

超新星の爆発メカニズムと マルチメッセンジャー

Explosion Mechanisms of
Core-Collapse Supernovae and
the Multi-messenger Observables

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新学術「地下素核研究」
第4回超新星ニュートリノ研究会
2018年1/7, 箱根静雲荘

Many workshops to celebrate 30 years of SN1987A !



Workshop on Supernova at Hyper-Kamiokande



12-13 February 2017

Asia/Tokyo timezone

Overview

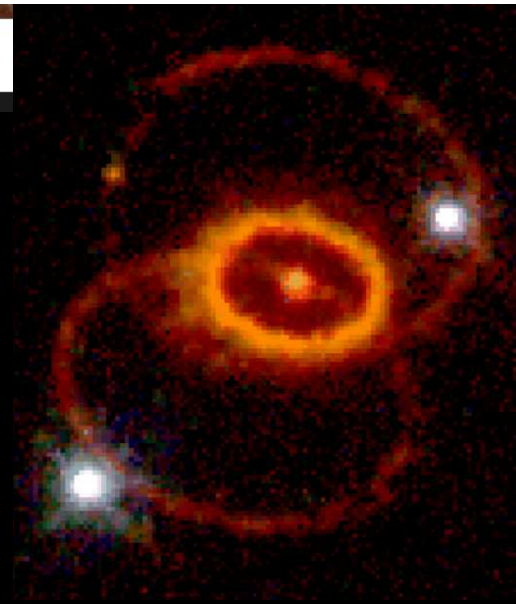
Scientific Programme

Timetable

Celebration of the 30th Anniversary of Supernova SN1987A



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prospekt" - al





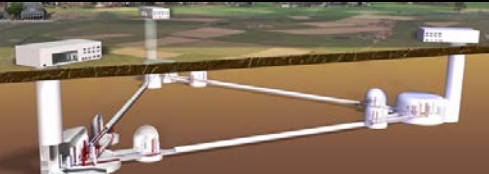
Looking back ~30 years, significant progress made in GW observation !

10^{-18}

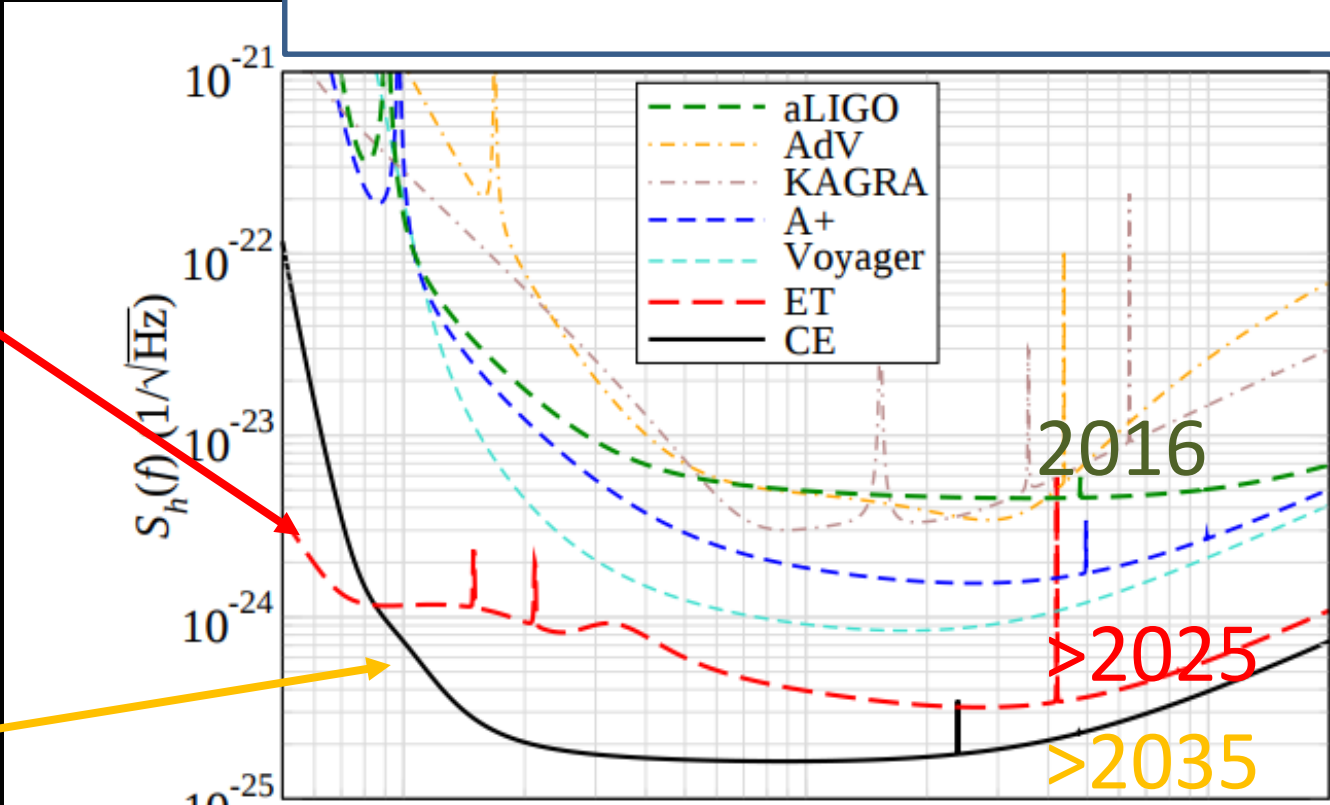
Typical thresholds of proto-types in 1989 (MIT, Garching, Caltech, Glasgow and Tokyo)

Sensitivity curves of laser interferometers

10 km long: Einstein Telescope (ET) could start ~2025.



40 km long: Cosmic Explore (CE) could operate ~2035.



GW astronomy is no more a dream !

The base-line and final goal (s)

What is the physics for exploding massive stars?

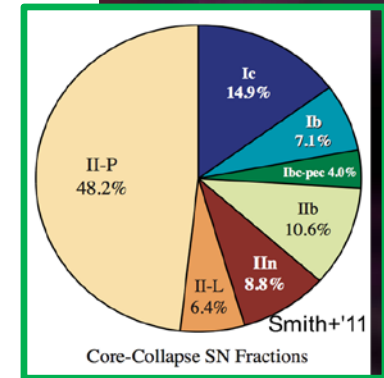
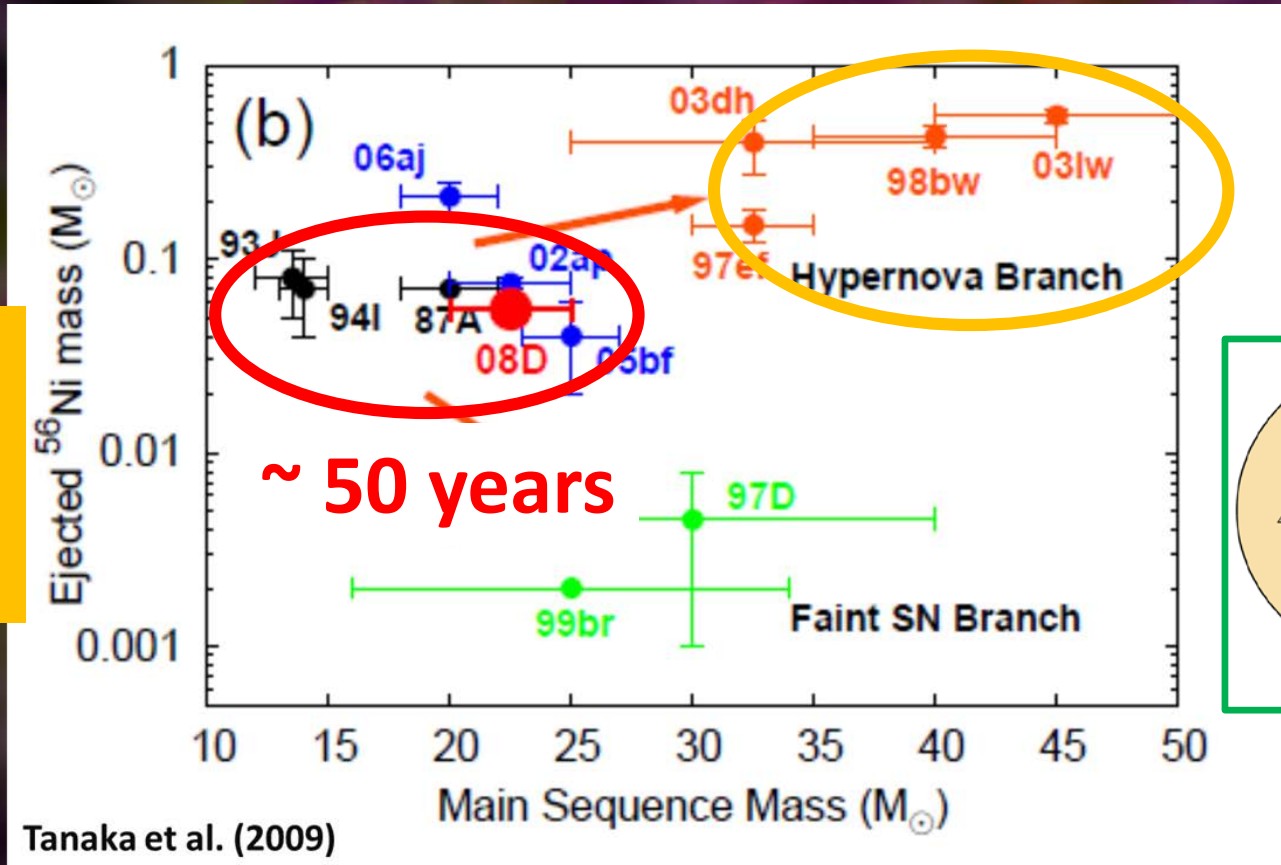
~ 10FOE

~ 1 FOE

FOE: Fifty-one-erg
 10^{51} erg

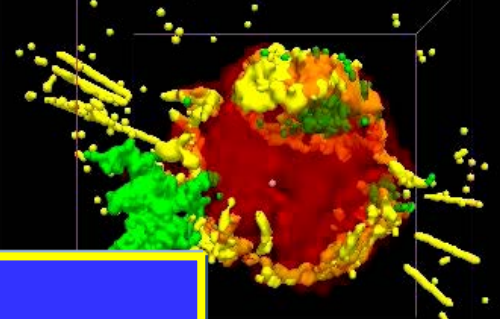
= 1 Bethe

Numerical study:
 Colgate & White
 (1966)

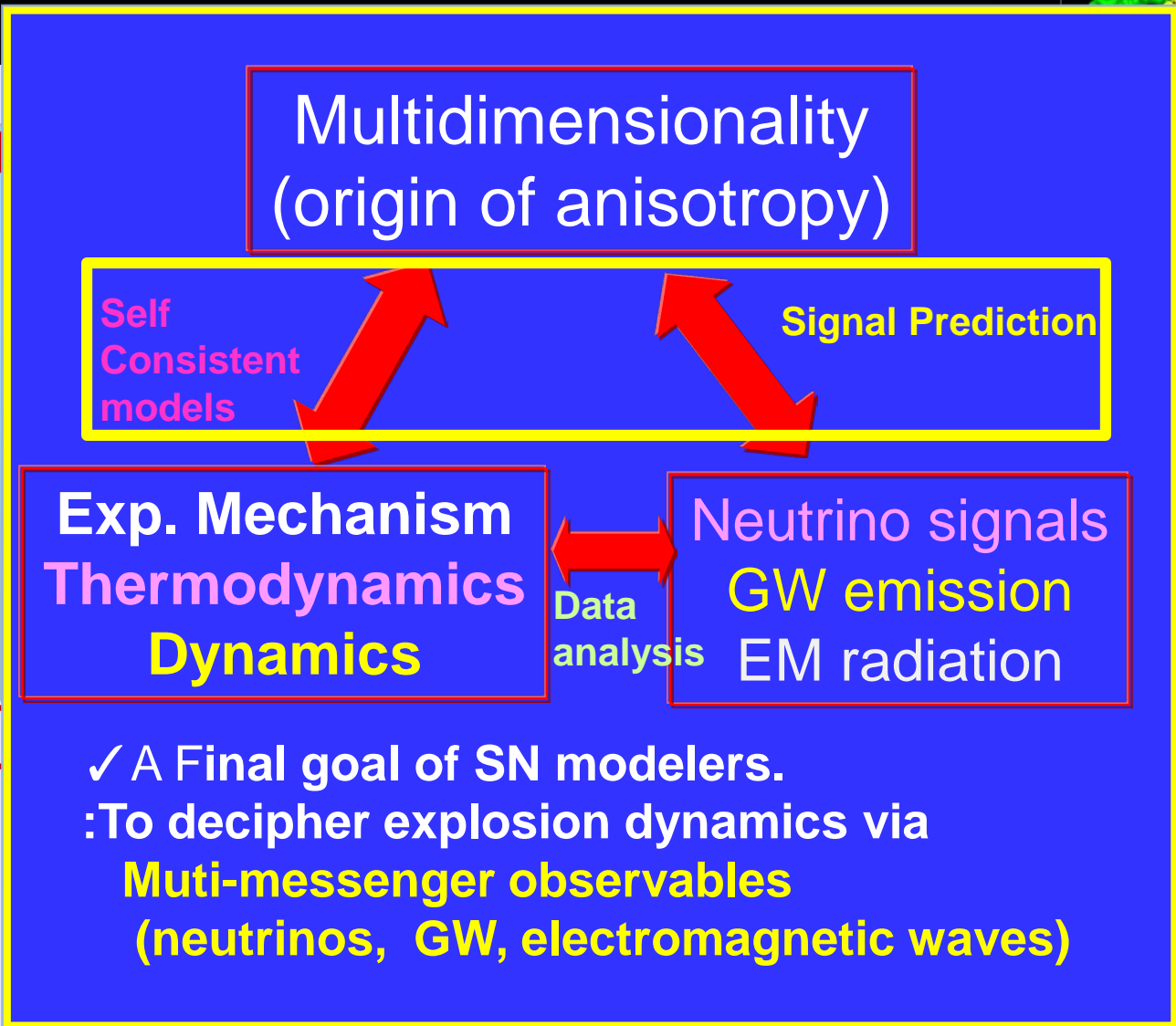
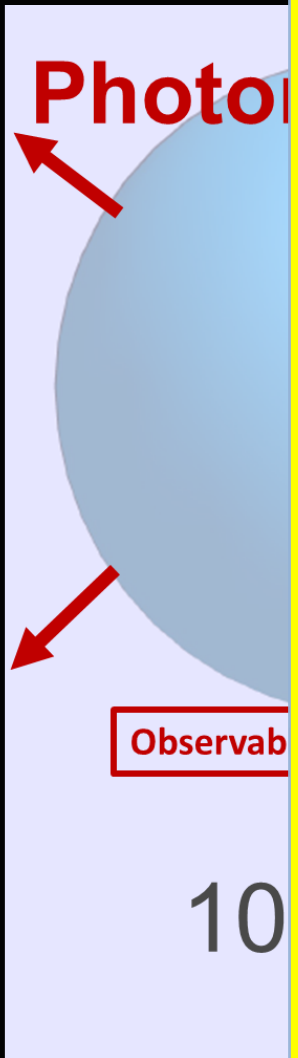


- 1). For which types of the progenitors (IIp, Ib/Ic, IIc) is rotation/B field most important ?
- 2). and 3). If important, why and how ?
- 4). Collapsar, Magnetar scenarios: Which one successful (or other) ? why ?
- 5). How long will it take before first-principles doable ? Strategies ?

Typical Scales of CCSN multimessenger



~350 years,
Type IIb



10

Outline

- ✓ **Brief introduction** (5 min)
what we can learn from SN multi-messengers ?

- ✓ **Recent progresses in “Supernova Theory”** (30 min)

- ☆ **The Core-Collapse Supernova Theory**

- :what is the essence to blow up massive stars?

- ☆ **Candidate mechanisms:** based on first-principle multi-D radiation-hydrodynamic simulations

- ✓ **Observational Signatures** (30 min)

- ☆ **Detectability of neutrino and gravitational-wave signals**

- ☆ **Perspectives toward “MM” astronomy**

- (correlation analysis between GWs/neutrinos,
electromagnetic messengers)

- What can we learn from the central engine ?**

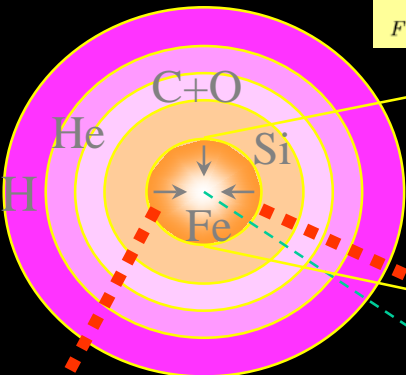
Standard scenario of core collapse

(e.g., Kotake+06, Foglizzo+14, M

$\rho_{\text{center}} \approx 10^{12} \text{g/cm}^3$
 $\rho_{\text{trap}} \approx 1.4 \times 10^{11} \text{g/cm}^3$ (view)

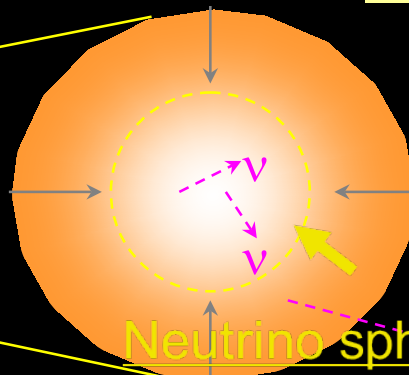
core collapse

$M \gtrsim 8M_{\odot}$
 $e^{-} + p \rightarrow \nu_e + n$
 $Fe + \gamma \rightarrow p + n$



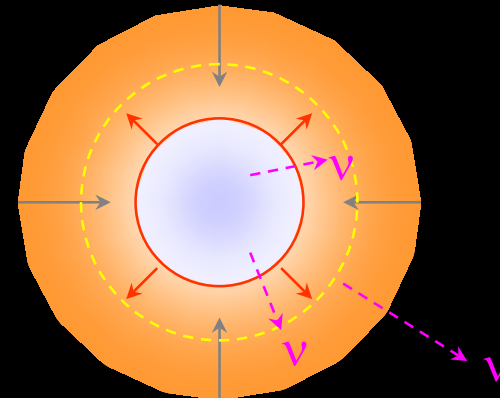
ν trapping

$\rho_c \sim 10^{12} \text{g/cm}^3$



core bounce

$\rho_c \sim 3 \times 10^{14} \text{g/cm}^3$

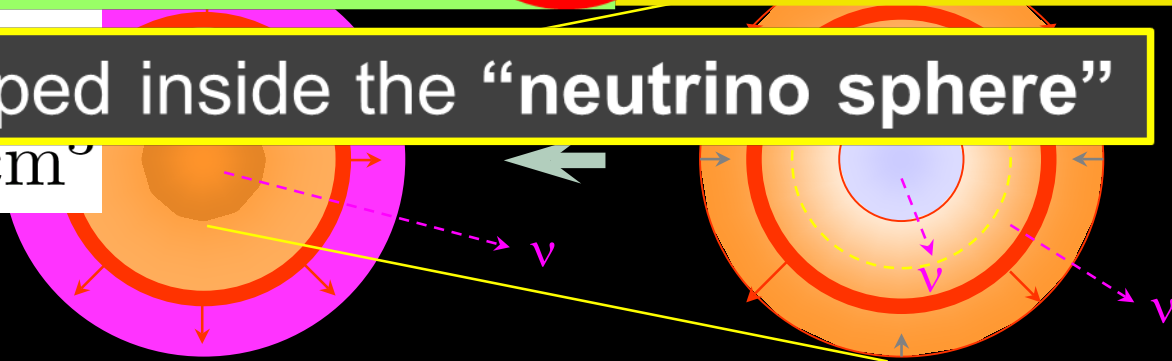


✓ Gravitational $e^{-} + p \rightarrow n + \nu_e$ in the iron core.

Neutrinos are trapped inside the "neutrino sphere"

$\rho_{\text{center}} \sim 3 \times 10^{14} \text{g/cm}^3$

$T \sim 1 \text{ MeV}$





(Weinberg)



(Salam)

Step 2 Neutrino Trapping

Neutrino



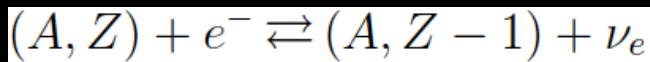
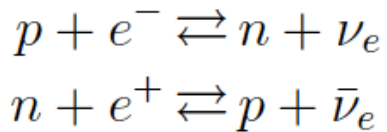
Weak interacting particle

(Neutral-rino: neutral-particle, light mass (<eV))

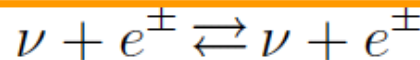
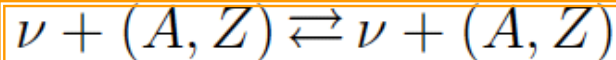
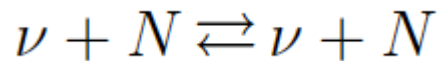
$$\sigma \sim 10^{-38} \text{cm}^2 \text{ (at 1GeV)} \ll \sigma_T \sim 10^{-25} \text{cm}^2$$

Representative Neutrino reactions in the SN core

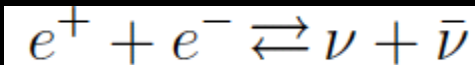
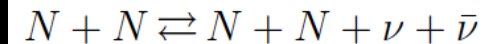
Neutrino emission/absorption



Scattering with (N: nucleon (A, Z): Nuclei)



Pair reaction
Bremsstrahlung

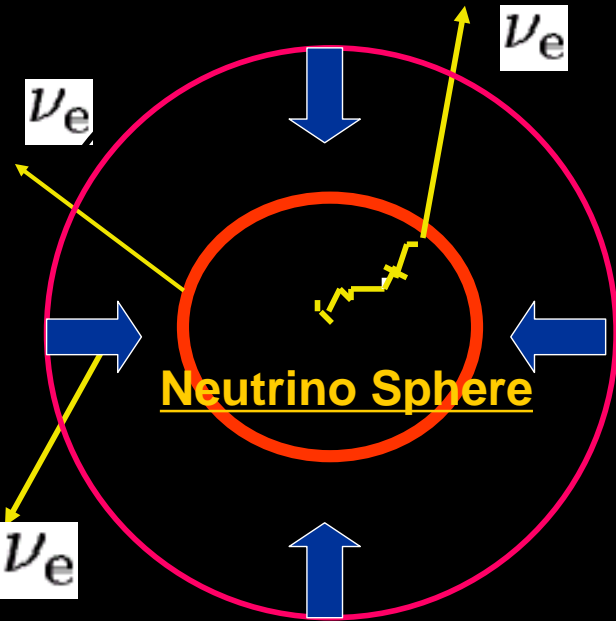


Mass energy				
0.000511 GeV	e	●	Electron	Thes
0.1066 GeV	μ	●	Muon	The t
1.777 GeV	τ	●	Tau	the e
				show
				be a

(see, however, Bollig et al (2017))



The condition of "Neutrino trapping"



To judge whether neutrinos can be trapped or not in the iron core, Compare the two timescales !

Free-fall timescale τ_{ff}

Diffusion timescale due to coherent neutrino-A scattering τ_{diff}

$$\left\{ \begin{aligned} \rho \frac{\partial^2 r}{\partial t^2} &= - \frac{GM_r \rho}{r^2} \\ \rho &\sim \frac{M}{4\pi/3R^3} \end{aligned} \right.$$

$$N_R = \left(\frac{R}{\lambda} \right)^2$$

Number of scattering (random walk)

$$N_R \frac{\lambda}{c} = \frac{R^2}{c\lambda}$$

Diffusion timescale

$$t_{diff} = \frac{3R^2}{c\lambda}$$

$$R_{core} = \left(\frac{3M^{1/3}}{4\pi\rho} \right)$$

$$t_{diff} \leq t_{dyn}$$

Mean free-path by the coherent scattering

Average neutrino energy

$$\lambda_\nu = \frac{1}{\sigma n_A} \approx 10^7 \text{cm} \left(\frac{\rho}{3 \times 10^{10} \text{gcm}^{-3}} \right)^{-1} \left(\frac{E_\nu}{12.6 \text{MeV}} \right)^{-2} \left(\frac{A}{56} \right)^{-1} \left(\frac{Y_e}{26/56} \right)^{-2/3}$$

$$E_\nu \approx \frac{5}{6} \mu_e = \frac{5}{6} \left(3\pi^2 \frac{\rho Y_e}{m_p} \right)^{1/3} \hbar c \approx 12.6 \text{MeV} \left(\frac{\rho}{3 \times 10^{10} \text{g cm}^{-3}} \right)^{1/3} \left(\frac{Y_e}{26/56} \right)^{1/3}$$

✓ The position of the neutrino sphere is energy-dependent !

$$R_\nu \approx 1.0 \times 10^7 \text{cm} \left(\frac{E_\nu}{10 \text{MeV}} \right)$$

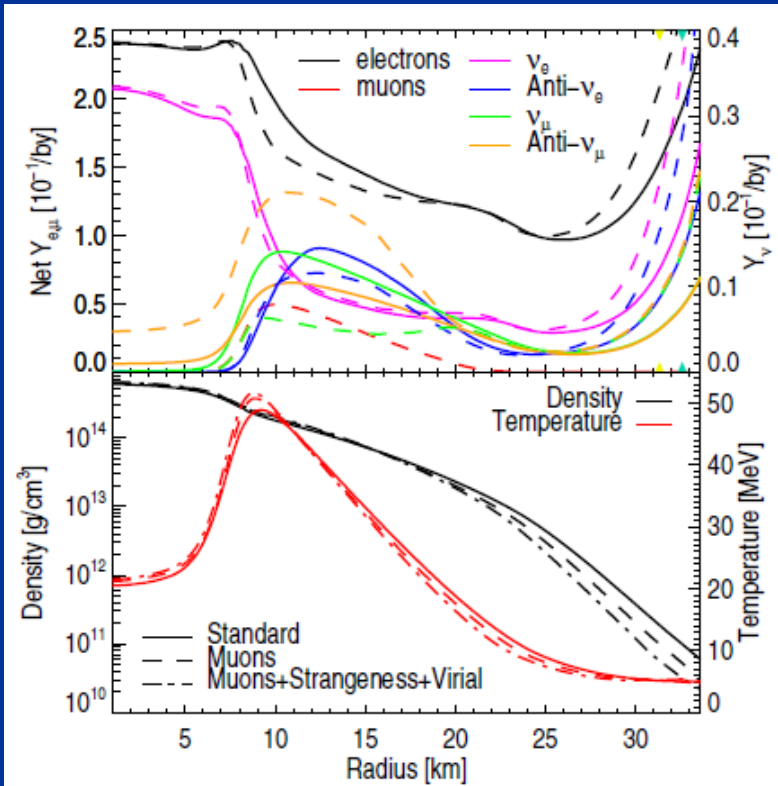
For lower-energy neutrinos, the neutrino sphere forms deeper inside, because they need a denser environment to be opaque!

Muon creation in supernova matter facilitates neutrino-driven explosions

2017,
PRL

R. Bollig,^{1,2} H.-T. Janka,¹ A. Lohs,³ G. Martínez-Pinedo,^{3,4} C.J. Horowitz,⁵ and T. Melson¹

¹Max-Planck-Institut für Astrophysik, Karl-Schwarzschild-Str. 1, 85748 Garching, Germany



- ✓ 20 M_{sun} star (Woosley & Heger 2007)
- ✓ VERTEX-PROMETHIUS code

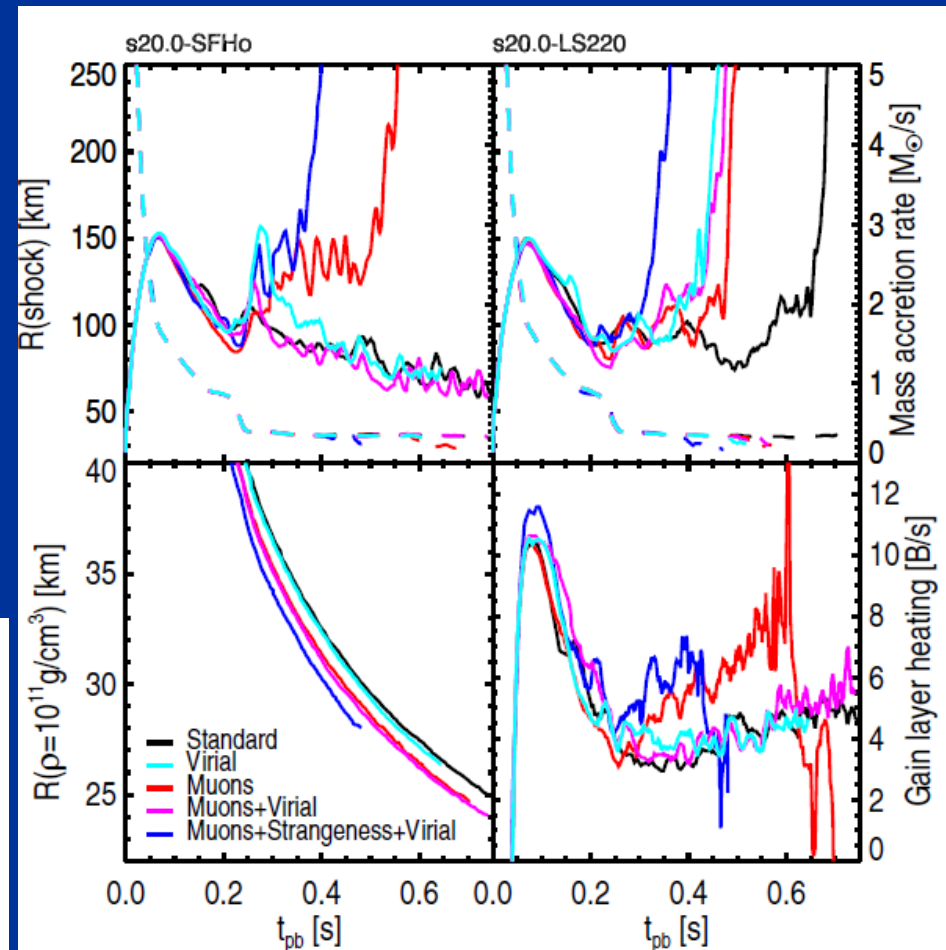


TABLE I. Neutrino reactions with muons.

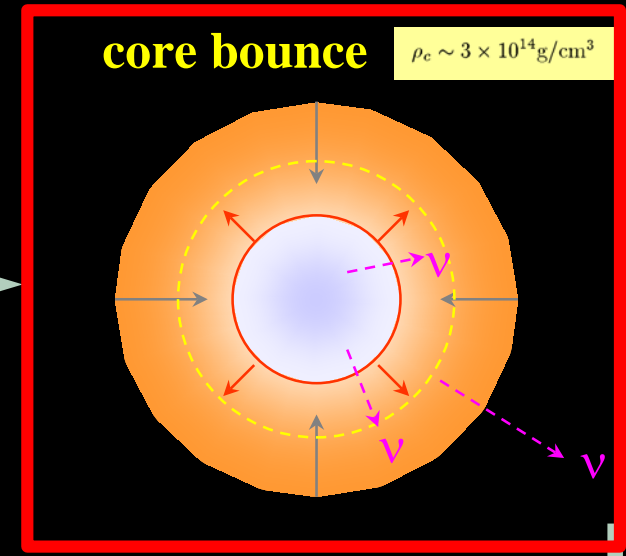
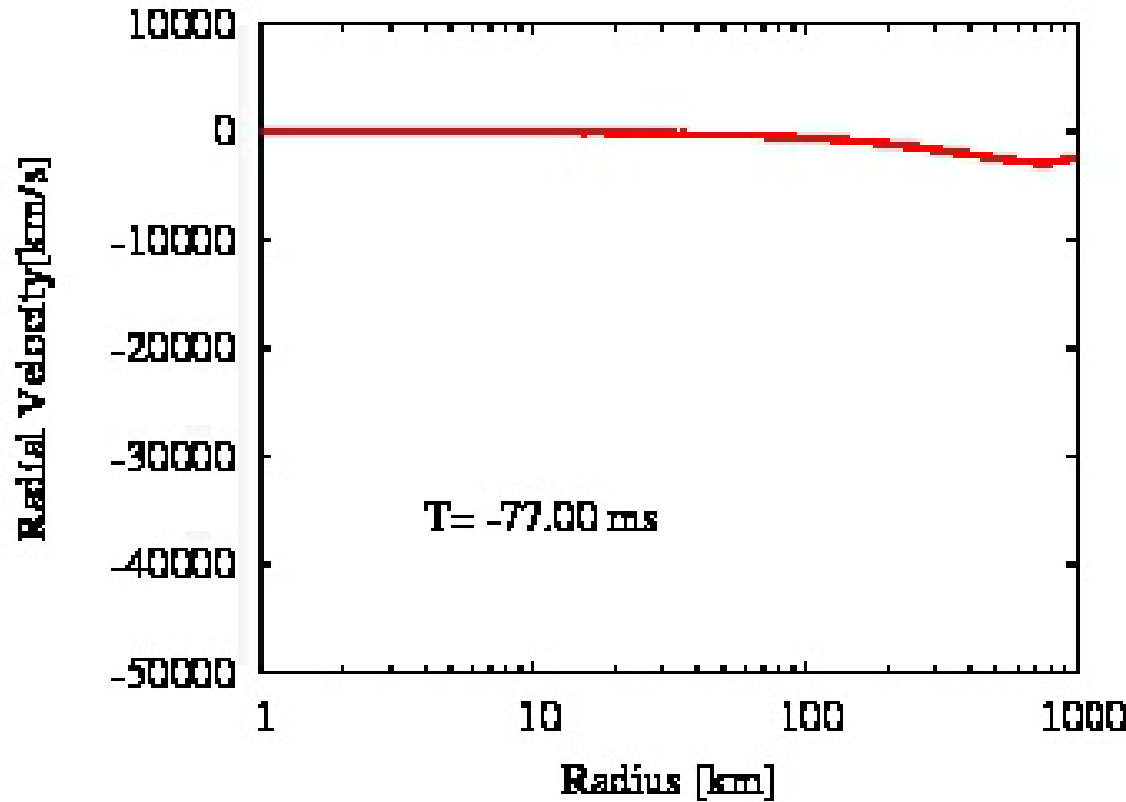
$\nu + \mu^- \rightleftharpoons \nu' + \mu^{-'}$	$\nu + \mu^+ \rightleftharpoons \nu' + \mu^{+'}$
$\nu_\mu + e^- \rightleftharpoons \nu_e + \mu^-$	$\bar{\nu}_\mu + e^+ \rightleftharpoons \bar{\nu}_e + \mu^+$
$\nu_\mu + \bar{\nu}_e + e^- \rightleftharpoons \mu^-$	$\bar{\nu}_\mu + \nu_e + e^+ \rightleftharpoons \mu^+$
$\bar{\nu}_e + e^- \rightleftharpoons \bar{\nu}_\mu + \mu^-$	$\nu_e + e^+ \rightleftharpoons \nu_\mu + \mu^+$
$\nu_\mu + n \rightleftharpoons p + \mu^-$	$\bar{\nu}_\mu + p \rightleftharpoons n + \mu^+$

Evolution of Radial velocity profiles

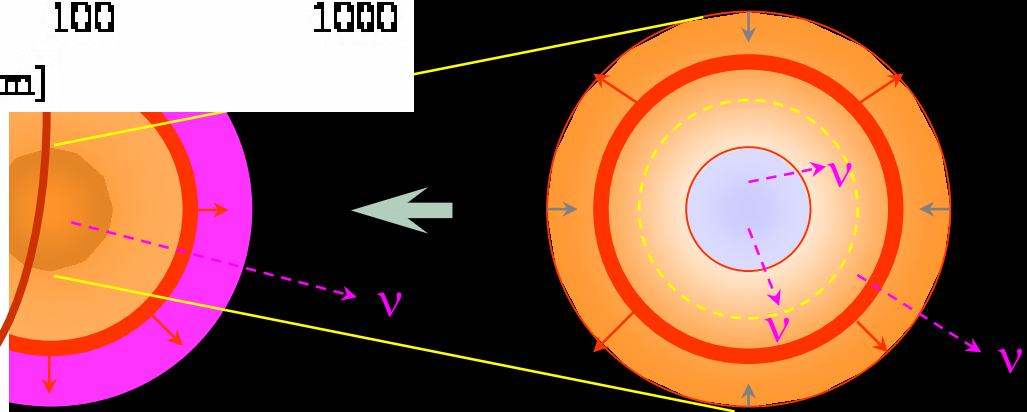
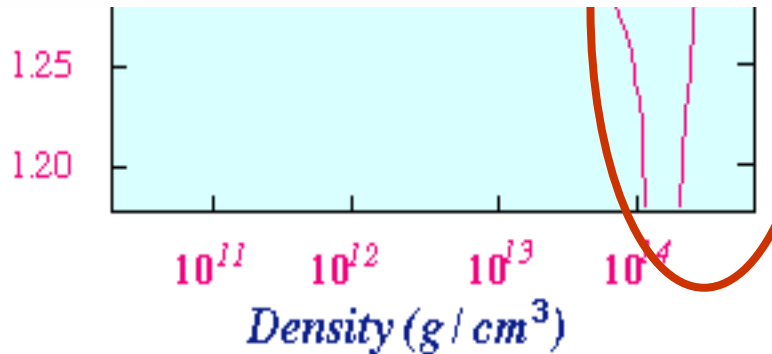
based on radiation-transport simulation (MGFLD) KK+06

(e.g. Kotake+06, Fogliizzo+14, Mezzacappa+15, Janka17 for a review)

pre-collapse SNe

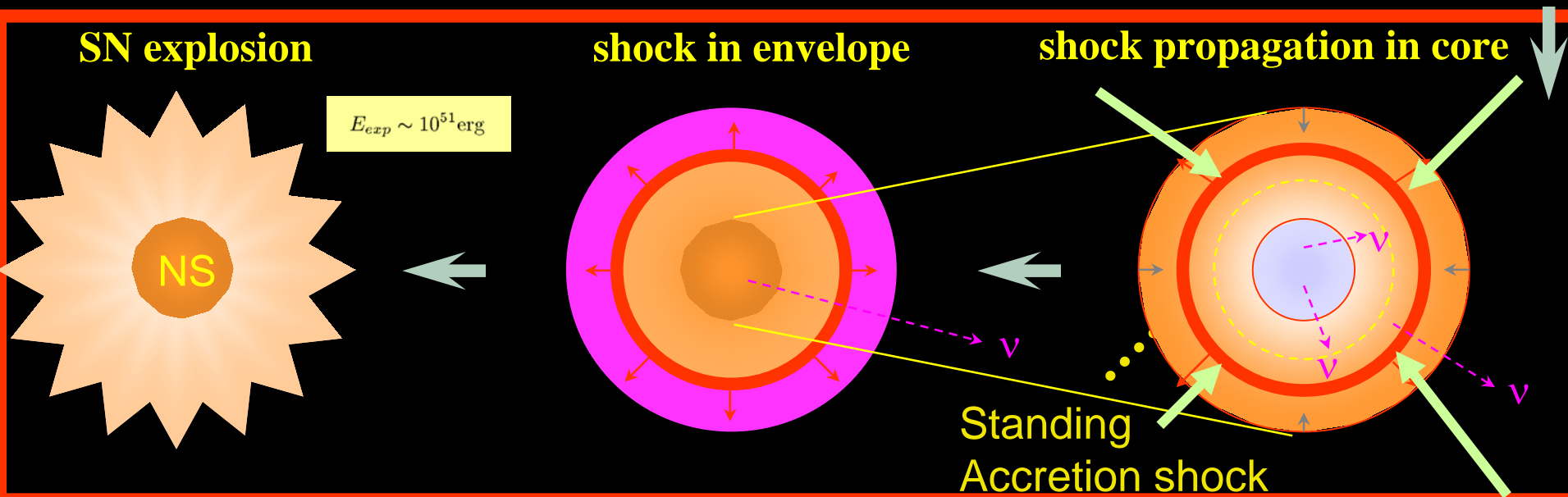


shock propagation in core

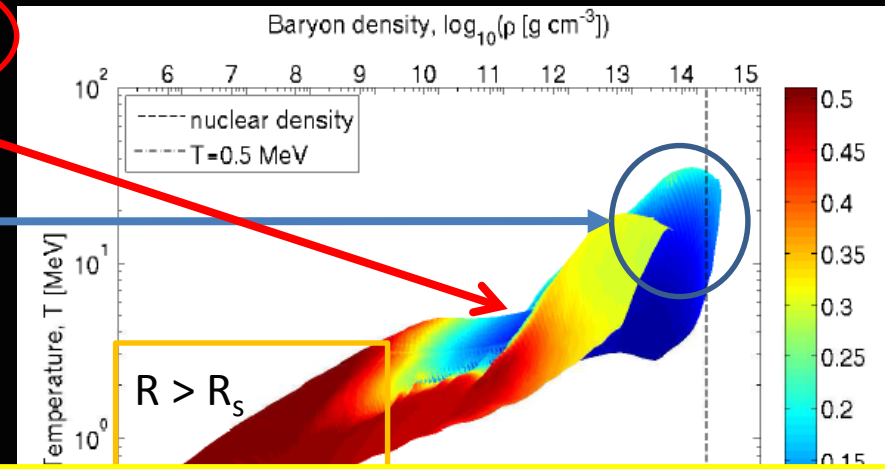
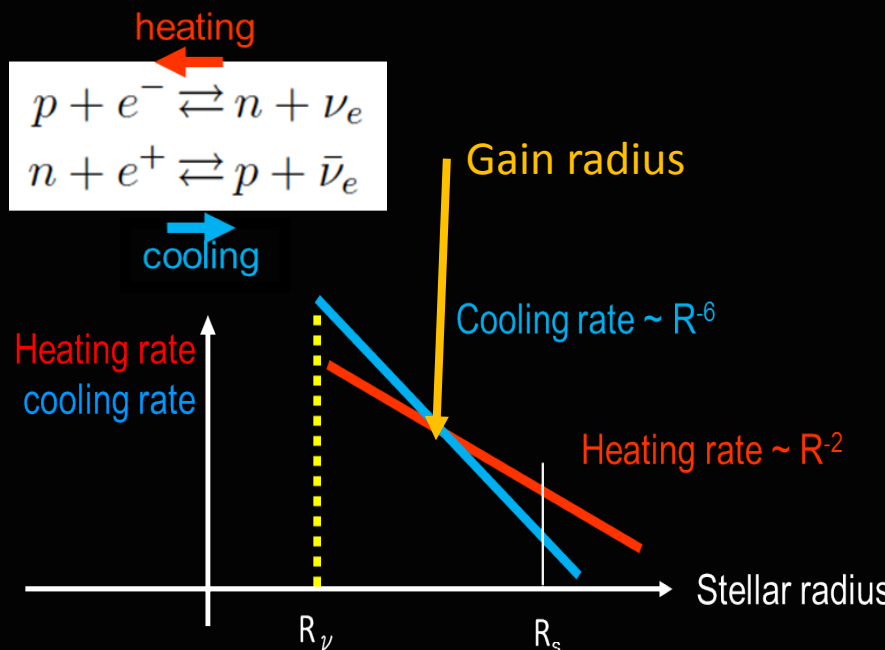
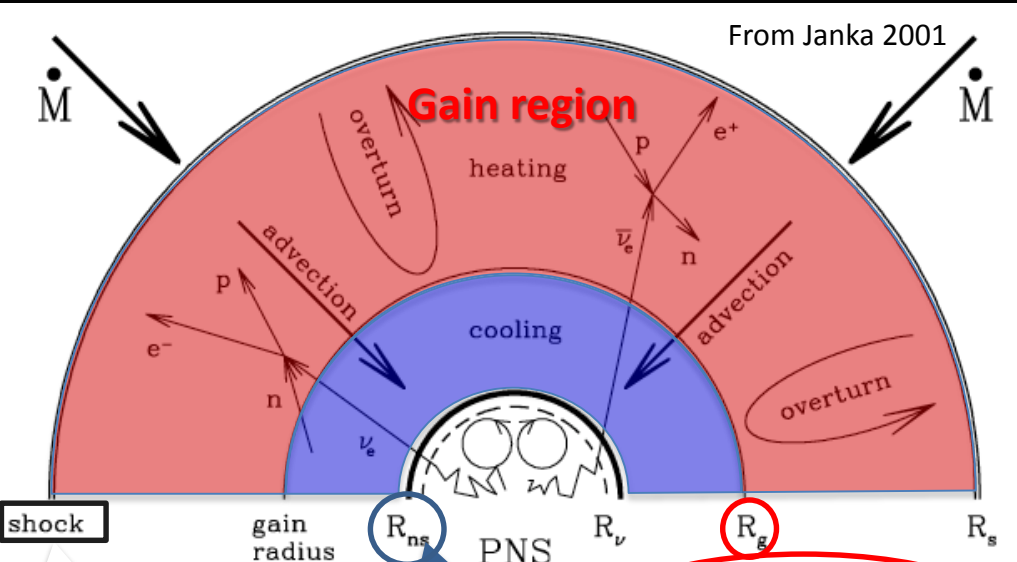


Short Summary (till shortly after bounce)

- ✓ SN simulations over these 20 years show that the bounce shock always stall because the kinetic energy of the shock is lost by the photo-dissociation of iron nuclei.
→ Direct “prompt” hydrodynamic explosion fails.
- ✓ The bounce shock turns into the standing accretion shock.
- ✓ The supernova problem is how to revive the stalled shock into explosion!



Typical scales after bounce and Density-Temperature relation



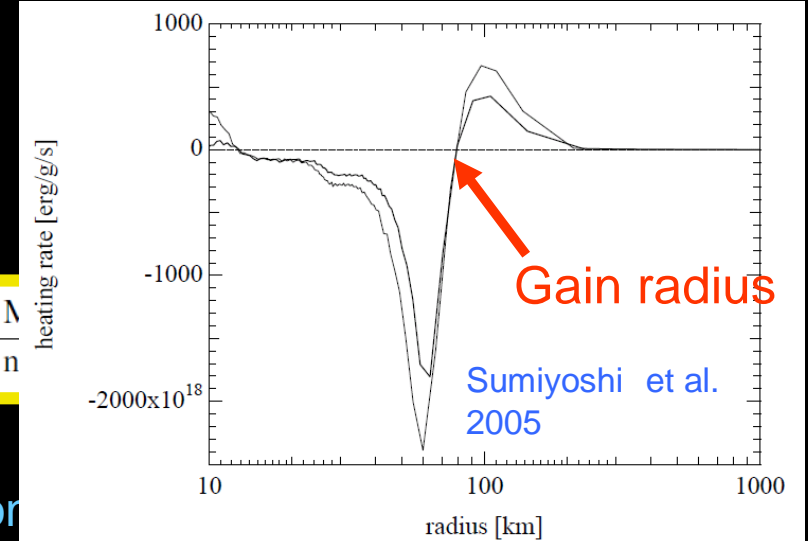
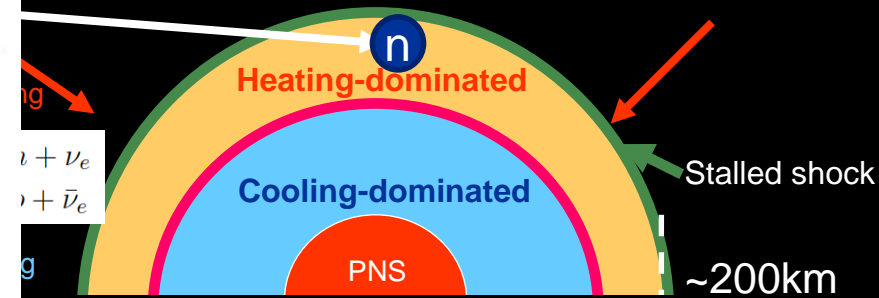
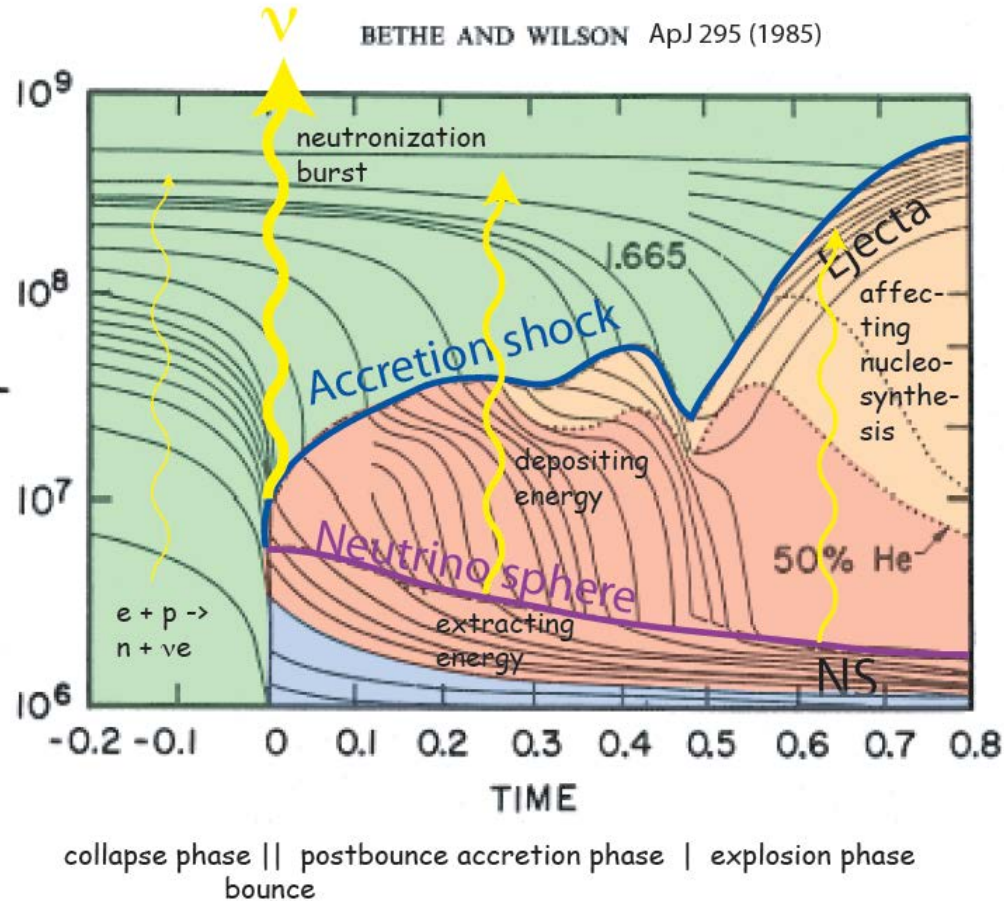
Gain Radius
 $R_g \sim 80\text{km}$

$R_{PNS}, R_\nu \sim 50\text{km}$
 $\rho_c \sim 10^{14} \text{ g/cc}, T_c \sim 10 \text{ MeV}$

$R_{\text{stalled_shock}} \sim 200\text{km}$
 ($R_{\text{core}} \sim 1500 \text{ km}$)
 $\rho < 10^9 \text{ g/cc}, T \sim 1 \text{ MeV}$

✓ **Travel time** (τ_{gain}) in the “gain region” **longer**
Gain mass (M_{gain} : mass in the “gain region”) **bigger,**
more favorable for explosions !

How the neutrino mechanism works ?



The gravitational binding energy of single neutron

$$-\frac{GM_{\text{NS}}m_u}{R} = -13.0 \left(\frac{M_{\text{NS}}}{1.4M_{\odot}} \right) \left(\frac{R}{150 \text{ km}} \right)^{-1} \text{ [MeV/nucleon]}$$

- ✓ If the neutrino heating could last $> \sim 0.25$ sec, the absorbed energy exceeds the local grav. binding energy \rightarrow inflows turns into outflows !
- ✓ More correctly neutrino cooling occurs, which delays the onset of explosion.

After +50 years of CCSN modeling : “Multi-D” neutrino mechanism

(pioneered by Colgate & White (1966), review by Kotake & Kuroda (2016), Janka (2012), Burrows (2013))

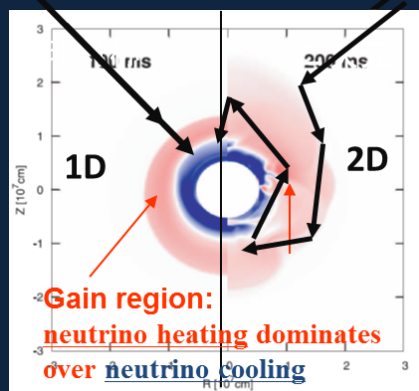
“Four steps” in neutrino-driven explosions

(see, e.g., Suwa et al. 2010,2011,2013, ApJ)

1st : After bounce, the bounce shock stalls.

2nd: Neutrino-driven convection and the **SASI**.
(**S**tanding-**A**ccretion-**S**hock-**I**nstability)

3rd: In the **heating region**, dwell-time of material **gets longer** due to non-radial motions in multi-D environments.
(Turbulence helps explosion).



4th: At around $O(100)$ s ms after bounce, neutrino-driven explosions set in.

2D radiation-hydro simulation of a $15 M_{\text{sun}}$ star

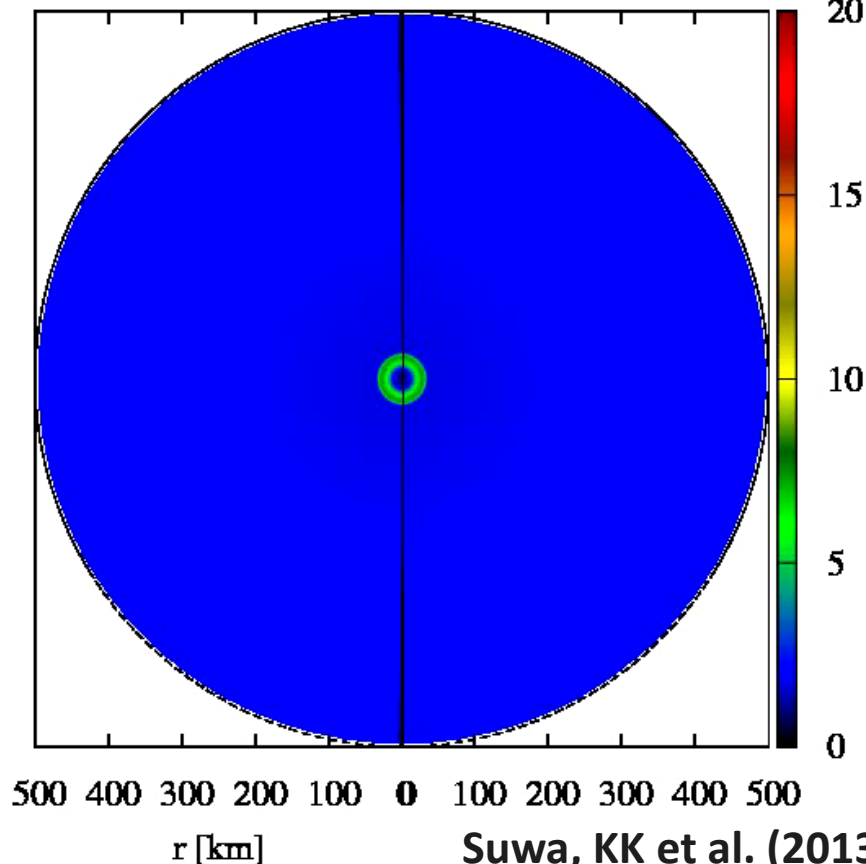
✓ IDSA scheme for spectral neutrino transport

✓ Lattimer-Swesty EOS ($K=220 \text{ MeV}$)

:compatible with $2 M_{\text{sun}}$ NS observation

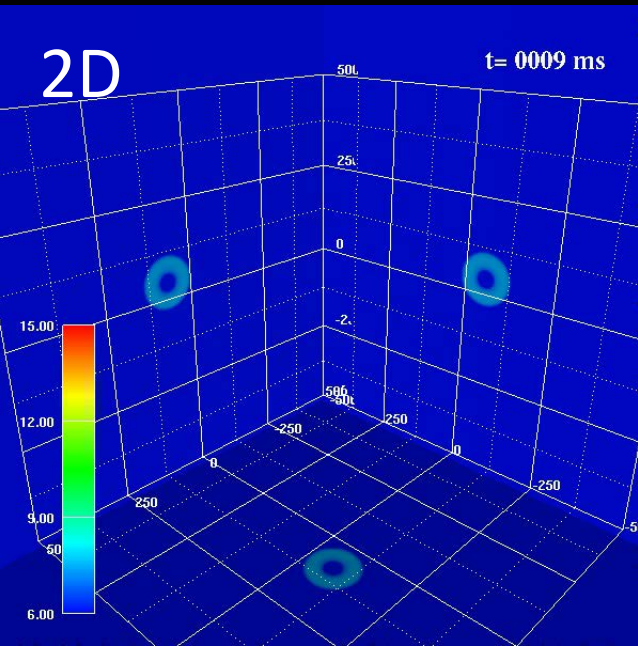
Color scale: entropy

$T = 188 \text{ ms}$

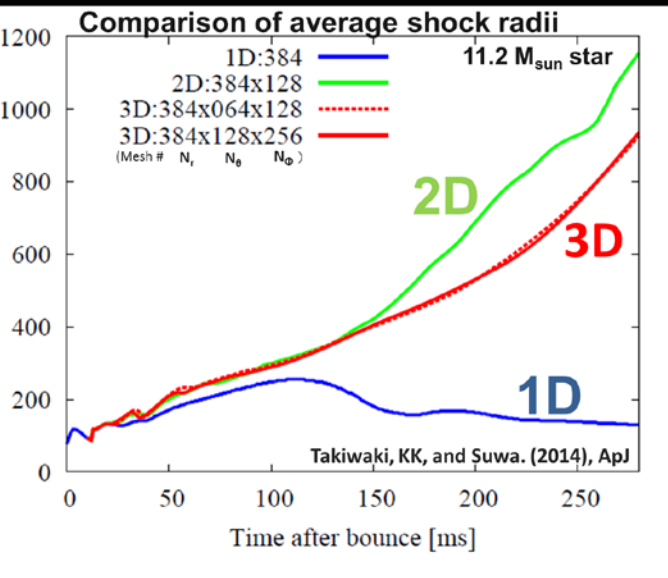


Suwa, KK et al. (2013)

3D vs. 2D



(e.g., Takiwaki, KK, Suwa (2012,2014), ApJ)



- ✓ 3D explosions are generally weaker than 2D.
(11.2, 27 M_{sun} : Hanke et al. (2014), however, not for 9.6 M_{sun} Melson et al. (2015))
- ⇒ The “3D vs. 2D problem” is progenitor dependent.
- ✓ No “Bethe” models obtained in 3D....
- ⇒ **Need to find ingredients to foster 3D explosions !**
Candidates: **Rotation** (Takiwaki+16, Summa+17),
General Relativity (Kuroda+14, Ott+17),
Microphysics (Melson+16, KK+17)

Multi-messengers from CCSNe:

Blue Giant (Red Giant: $\times 100$)

Phot

Explosion multi-dimensionality
(origin of anisotropy)

Simulations

Signal
Prediction

Data Analysis

Exp. Mechanism
Dynamics
Thermodynamics

GW emission
Neutrino signals
photon

Obse

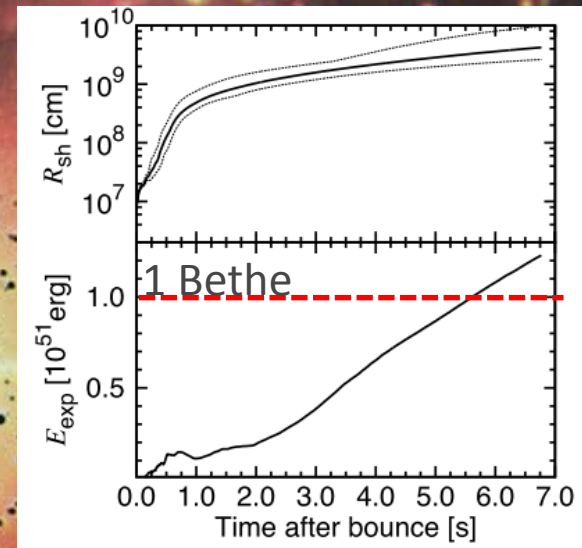
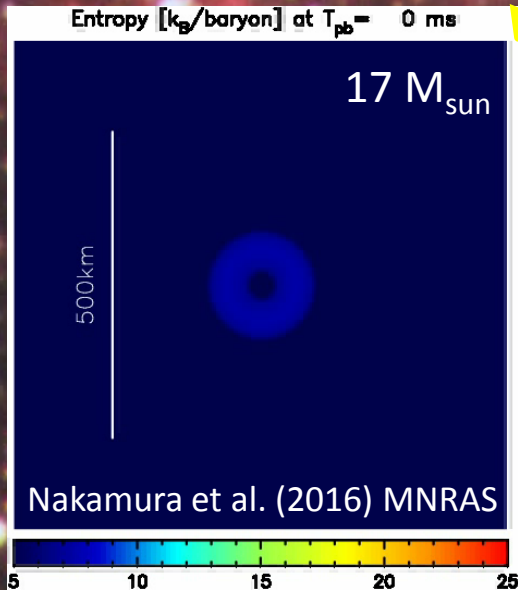
✓ A final goal of SN modellers is.
“To identify the supernova mechanism from
CCSN multi-messengers
(GW/neutrino/EM observations) ! “

15 km

neutron star

engine !

Drill for SN 2018xx !

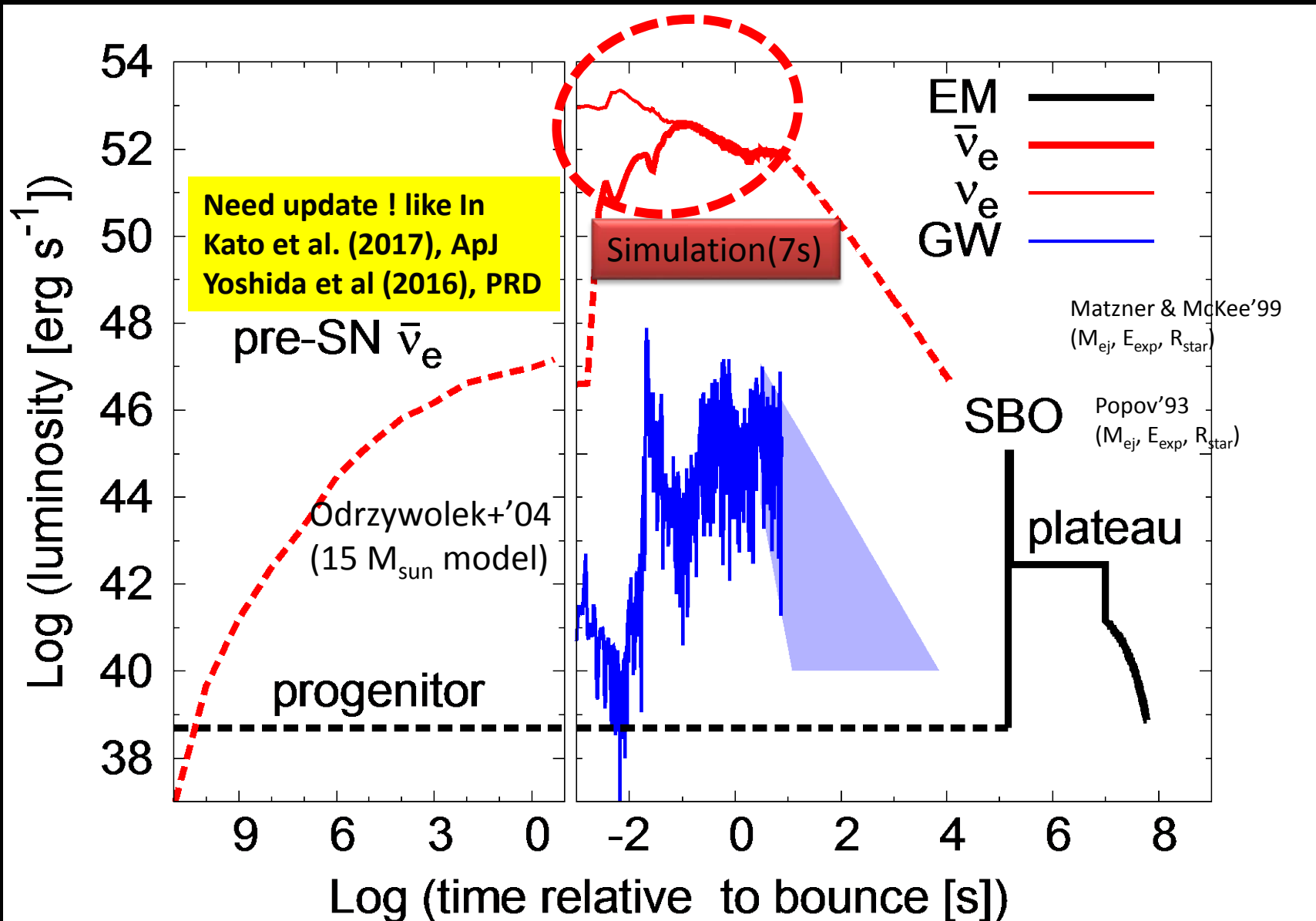


For the next Galactic event (several/century..),
how we observe multi-messengers and
what we can learn about the
supernova physics ?

Overview of “Multi-messenger signals” from exploding $17 M_{\text{sun}}$ star

Nakamura, Horiuchi, Tanaka, Hayama, Takiwaki, KK (MNRAS) 2016

Energetics: $E_{\text{neutrino}} \sim 10^{53}$ erg, $E_{\text{kinetic}} \sim 10^{51}$ erg, $E_{\text{photon}} \sim 10^{49}$ erg, $E_{\text{GW}} \sim 10^{46}$ erg

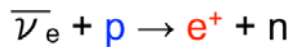


First Alert: Neutrinos ! (here for a Galactic event @ 8.5 kpc)

Nakamura, Horiuchi et al. (2016) MNRAS

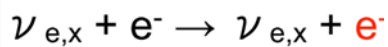
Super Kamiokande (SK):

Inverse- β decay

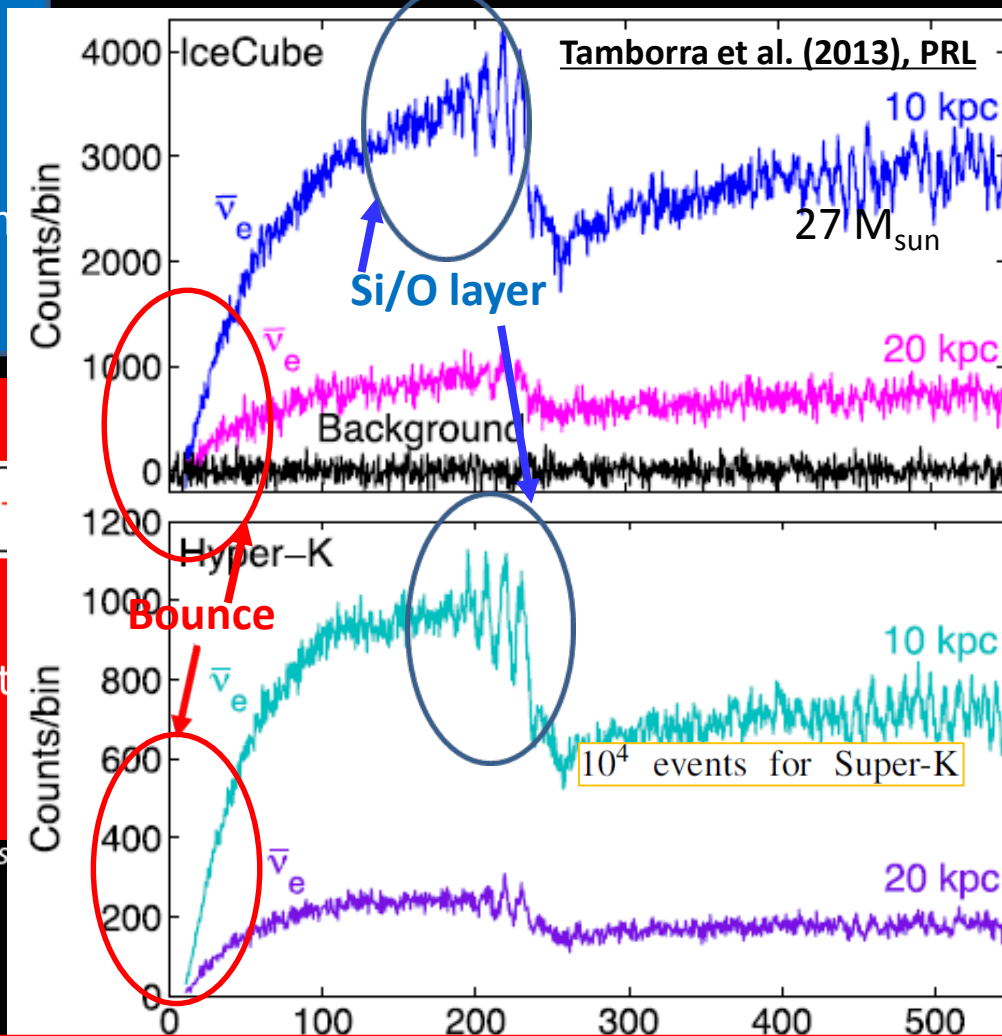
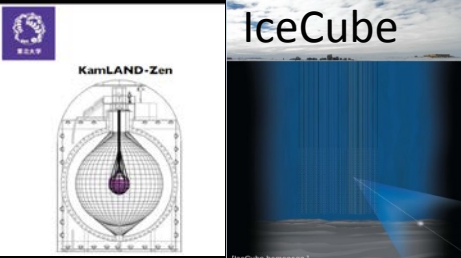
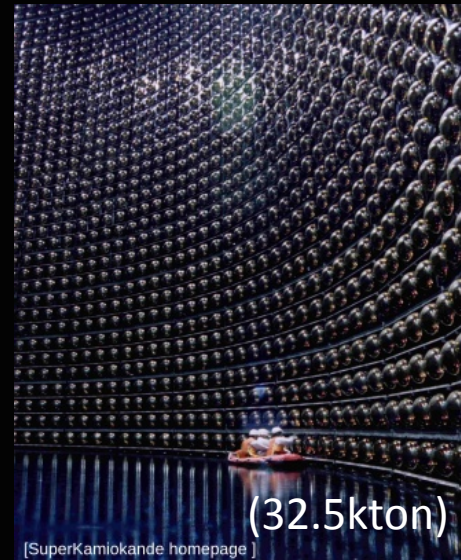


~15000 events
Timing information
~bounce ± 3 ms
(e.g., IceCube)

Electron scattering



~514 events
 \Rightarrow Position information
~6° (SK),
~3° (Gd-SK)
Beacom & Vagins



✓ For a galactic source, Can learn about SN physics.
(Bounce time, progenitor structure, compactness
(see Nakamura-san's talk !)
SASI modulation \Rightarrow Direct messenger of explosion.

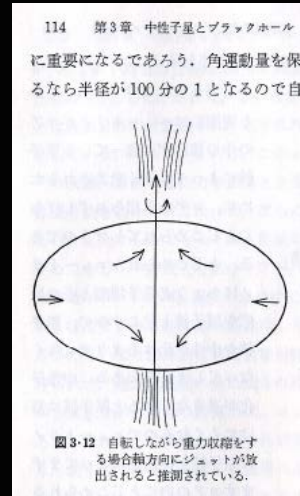
From an old textbook....

(1) Importance of rotation, B-field

第3章 中性子星とブラックホール

佐藤 勝彦
K. Sato

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(2) Multi-messenger astronomy

d) 超新星爆発とニュートリノ・重力波の観測

超新星爆発によりブラックホールや中性子星が作られる時、大量されるだけでなく、大量のニュートリノや強い重力波も放出される

(2) や (3) の場合であっても我々の銀河内で起こる場合その検出は可能であり、ニュートリノ検出との同時観測を通じて爆発機構の解明にきわめて重要である。

速く自転していたのでは、そもそも重力収縮が起こらないので弱い

恒星社

Stellar evolutions
and their fate:
edited by D. Sugimoto
published in 1979

1979年7月

杉本大一郎

Gravitational Waves (GWs) from Stellar Collapse

(see reviews in Ott (2009), Fryer & New (2011), Kotake (2013), Kotake and Kuroda (2016) in "Handbook of Supernovae")

GW amplitude from the quadrupole formula

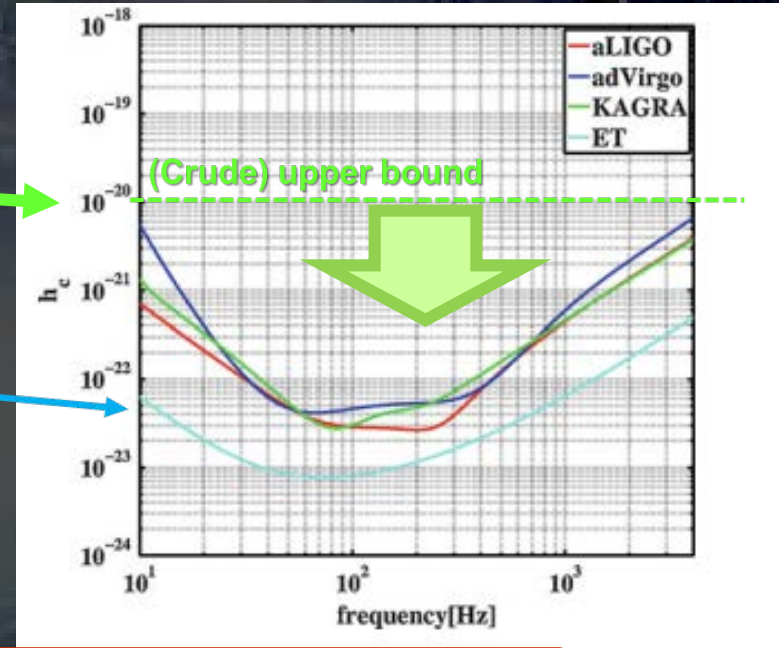
$$h_{ij} = \frac{2G}{c^4 R} \frac{\partial^2}{\partial t^2} Q_{ij} \sim \frac{R_s}{R} \left(\frac{v}{c}\right)^2$$

Quadrupole moment

Typical values at the formation of Neutron Star (NS)

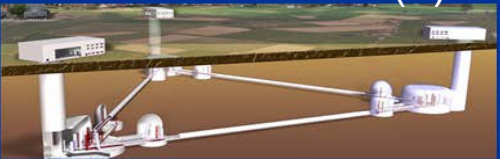
$$R_s = 3 \text{ km} \left(\frac{M}{M_\odot}\right) \quad v/c = 0.1 \quad R = 10 \text{ kpc}$$

$$h \sim 10^{-20}$$



Good news ! (Future)

10 km long: Einstein Telescope (ET) could start ~2025 (?)



40 km long: Cosmic Explore (CE) could operate ~2035 (?)

✓ CCSN event in our galaxy (several/century) is primary target !

$h_{ij} = \epsilon \frac{R_s}{R} \left(\frac{v}{c}\right)^2$ ϵ : the degree of anisotropy.
 If collapse proceeds spherically, $\epsilon = 0$ no GWs !

What makes the SN-dynamics deviate from spherical symmetry is essential for the GW emission mechanism !

Two candidates : **The key** is “initial rotation rate” (Ω_0) of the iron core

(See reviews in Janka ('17), Mezzacappa et al. ('15), Foglizzo et al. ('15), Burrows ('13), Kotake et al. ('12))

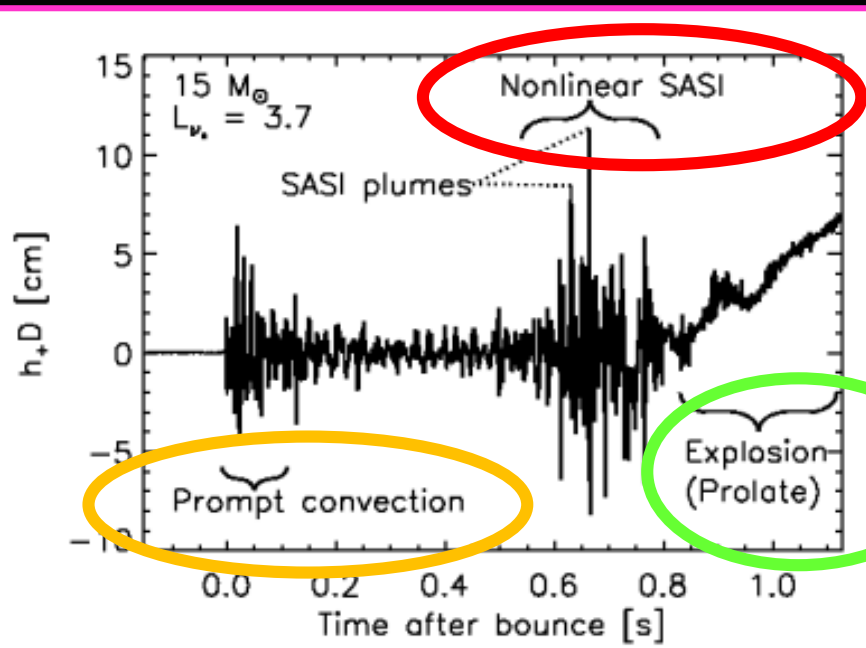
	Neutrino mechanism	MHD mechanism
Progenitor	Non- or slowing- rotating star ($\Omega_0 < \sim 0.1$ rad/s)	Rapidly rotating star with strong B fields ($\Omega_0 > \sim \pi$ rad/s, $B_0 > \sim 10^{11}$ G)
Main origin of GW emission	Turbulent Convection and SASI	Rotating bounce and Non-axisymmetric instabilities
Progenitor fraction	Main players	~ 1 (?) % (Woosley & Heger (07), ApJ): (hypothetical link to magnetar, collapsar)



(see also, Burrows et al. ('17), Melson et al. ('15), Lentz et al. ('15), Roberts et al. ('16), B. Mueller ('15), Takiwaki et al. ('16))

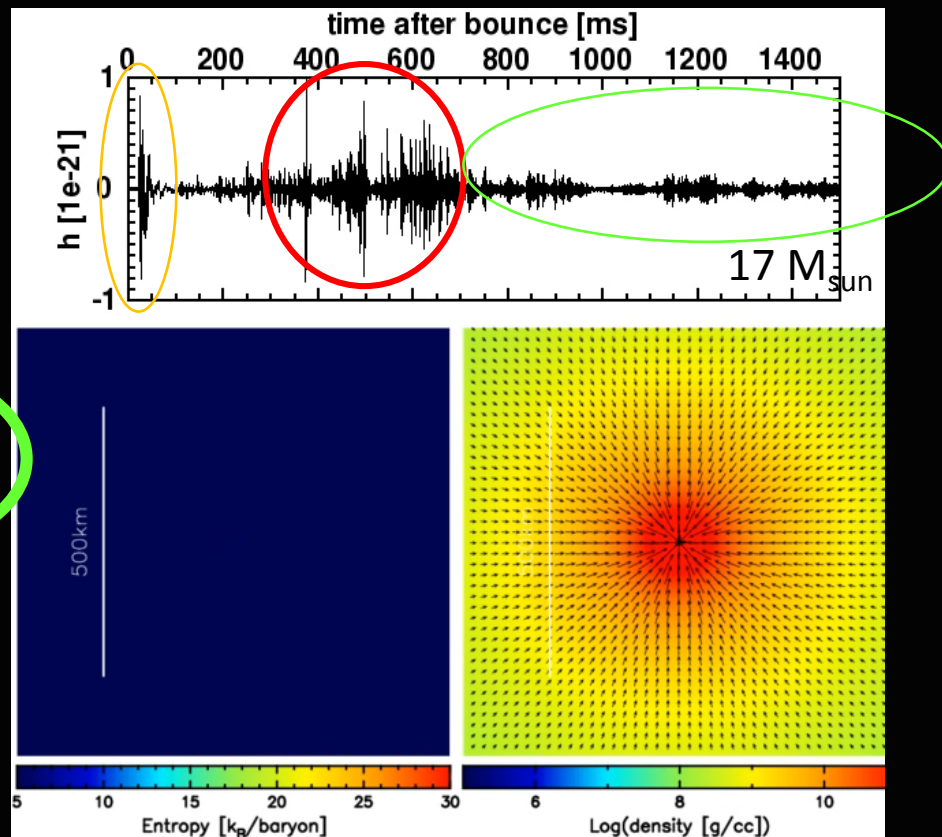
GW signatures from 2D neutrino-driven explosion

Waveform from Murphy et al. (2009) ApJ



(Later confirmed by B. Mueller et al. ('13), ApJ, Yakunin et al. (2015), PRD)

Waveform from Nakamura et al. ('16) MNRAS



✓ **Three generic phases** in neutrino-driven models:

- Prompt-convection phase** : within ~50 ms post-bounce
- Non-linear phase (Convection/SASI)** : Downflows hit the PNS surface
- Explosion phase** : Long-lasting signal but terminates if BH forms

(Müller et al. (2004, ApJ), Cerda-Duran et al. (2013, ApJ))

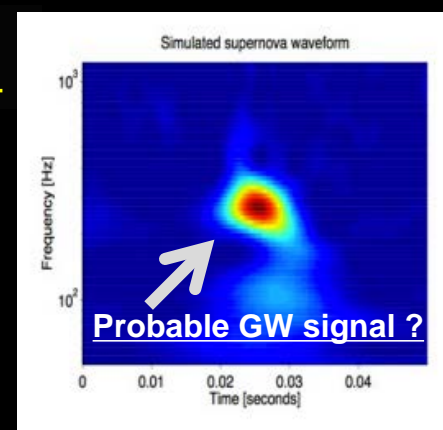
✓ Waveforms **have no template character**: stochastic explosion processes.

How to detect GWs with no-template features...

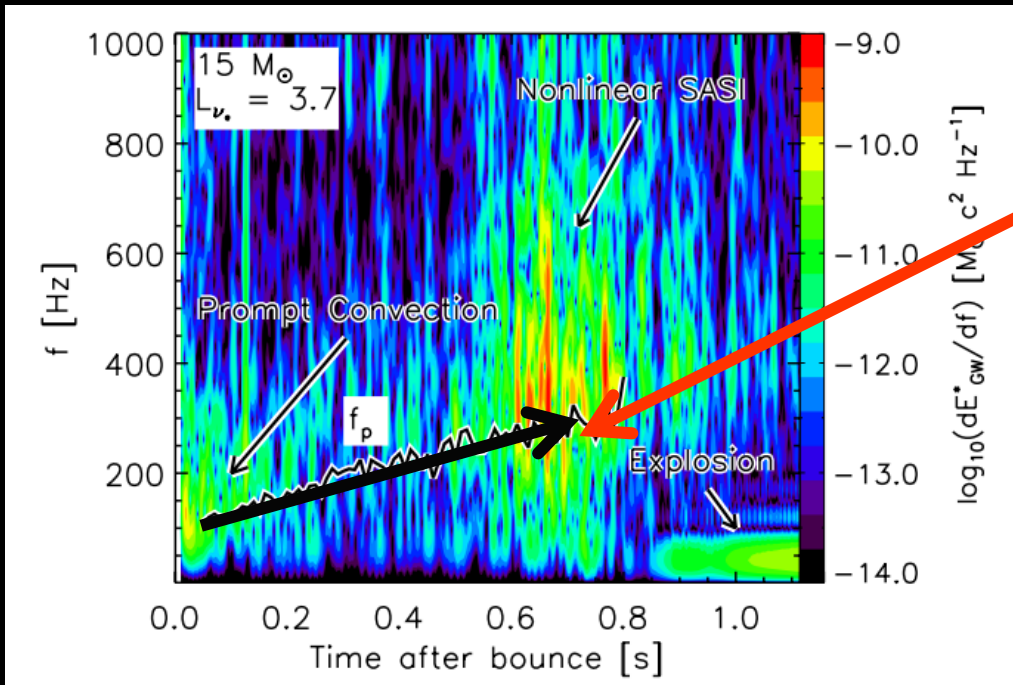
✓ **Excess power method:** Flanagan & Hugh (1998)

⇒ Decompose data-stream into time-frequency domains

⇒ Search for “hot” regions with excess power in the spectrogram !



✓ **GW spectrogram** from Murphy et al. ('09) ApJ.



★ Increase of typical frequency
⇒ the g-mode frequency of PNS

$$f_p = \frac{N}{2\pi} = \frac{1}{2\pi} \frac{GM}{R^2} \sqrt{\frac{(\Gamma - 1)m_n}{\Gamma k_b T}} \left(1 - \frac{GM}{Rc^2}\right)^{3/2}$$

M, R, T : mass, radius & temperature of PNS, Γ : stiffness of EOS

Due to mass accretion, $M \uparrow$, $R \downarrow$

⇒ $f_p \uparrow$

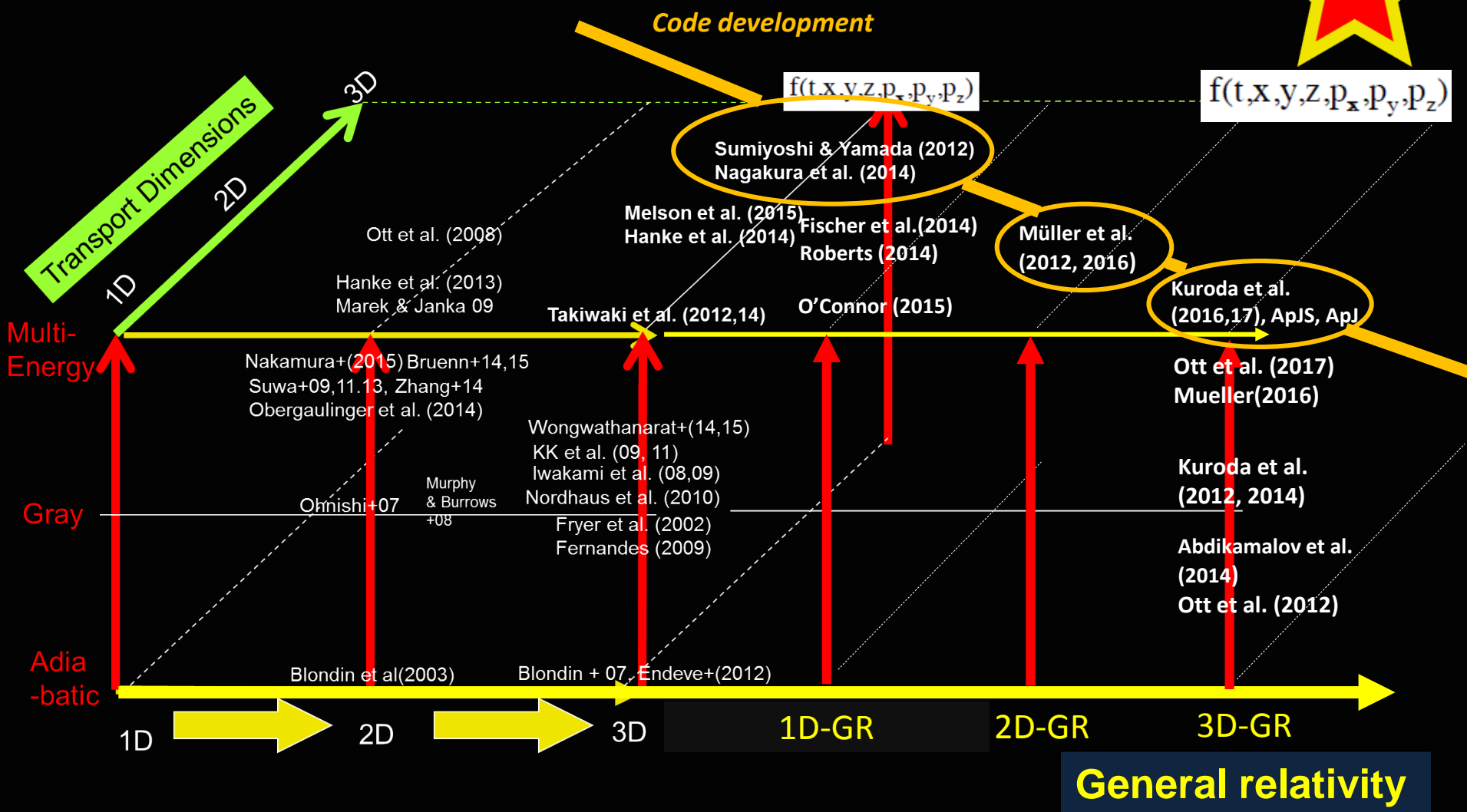
(see complete derivation in B. Mueller et al. ('13), ApJ)

- ✓ (With no template character...) **Three generic phases are in the spectrogram !**
- ✓ Secular increase of typical GW frequency (f_p) **reflects the PNS evolution.**
- ✓ On top of f_p , the high frequency component comes from strong downflows to PNS.
- ✓ **These qualitative features are common to more recent 2D and 3D models.**
- ✓ The GW amplitudes $\sim 1/10$ than in 2D (KK et al. 2009, Yakunin et al. (2017))

Recent Status of CCSN simulations

Disclaimer: only CCSNs

Ultimate goal:
7D Boltzmann transport in full GR MHD hydrodynamics
with increasing microphysical inputs !



First full-3D-GR simulations with multi-energy neutrino transport (M1)

Kuroda, KK, Takiwaki, Thielemann submitted

(see also, GR models using the CoCoNuT code (CFC(+)) by Cerda-Duran+2011, Obergaulinger and Aloy (2017): 2D by Dimmelmeier et al. (2007), B. Mueller (2015), B. Mueller et al. (2017):3D)

✓ **“FUGRA”**: Fully General Relativistic code with multi-energy neutrino transport

Kuroda, Takiwaki, and KK, ApJS. (2016)

The marriage of **BSSNOK formalism** (3D GR code, Kuroda & Umeda (2010, ApJS) $\mathcal{G} = \{\tilde{\gamma}_{ij}, \tilde{A}_{ij}, \phi, K, \tilde{\Gamma}^l, \alpha, \beta^i\}$) + **M1 scheme**; Shibata+2011, Thorne 1981, (see also, Just et al. (2015), O'Connor (2015) for recent work)

✓ Evolution equation of neutrino radiation energy

$$\begin{aligned} \partial_t \sqrt{\gamma} E_{(\varepsilon)} + \partial_i \sqrt{\gamma} (\alpha F_{(\varepsilon)}^i - \beta^i E_{(\varepsilon)}) + \sqrt{\gamma} \alpha \partial_\varepsilon (\varepsilon \tilde{M}_{(\varepsilon)}^\mu n_\mu) \\ = \sqrt{\gamma} (\alpha P_{(\varepsilon)}^{ij} K_{ij} - F_{(\varepsilon)}^i \partial_i \alpha - \alpha S_{(\varepsilon)}^\mu n_\mu), \end{aligned}$$

✓ Evolution equation of radiation flux

$$\begin{aligned} \partial_t \sqrt{\gamma} F_{(\varepsilon)i} + \partial_j \sqrt{\gamma} (\alpha P_{(\varepsilon)i}^j - \beta^j F_{(\varepsilon)i}) - \sqrt{\gamma} \alpha \partial_\varepsilon (\varepsilon \tilde{M}_{(\varepsilon)}^\mu \gamma_{i\mu}) \\ = \sqrt{\gamma} [-E_{(\varepsilon)} \partial_i \alpha + F_{(\varepsilon)j} \partial_i \beta^j + (\alpha/2) P_{(\varepsilon)}^{jk} \partial_i \gamma_{jk} + \alpha S_{(\varepsilon)}^\mu \gamma_{i\mu}] \end{aligned}$$

✓ Analytic Closure with the use of **Minerbo-type Eddington factor** (Murchikova, Abdikamalov + (2017))

$$P_{(\varepsilon)}^{ij} = \frac{3\chi_{(\varepsilon)} - 1}{2} P_{\text{thin}(\varepsilon)}^{ij} + \frac{3(1 - \chi_{(\varepsilon)})}{2} P_{\text{thick}(\varepsilon)}^{ij}$$

$$\chi_{(\varepsilon)} = \frac{5 + 6\bar{F}_{(\varepsilon)}^2 - 2\bar{F}_{(\varepsilon)}^3 + 6\bar{F}_{(\varepsilon)}^4}{15}$$

Closed set of rad-hydro equations

$$\begin{aligned} \partial_t \rho_* + \partial_i (\rho_* v^i) &= 0, \\ \partial_t \sqrt{\gamma} S_i + \partial_j \sqrt{\gamma} (S_i v^j + \alpha P \delta_j^i) \\ &= -\sqrt{\gamma} [S_0 \partial_i \alpha - S_k \partial_i \beta^k - 2\alpha S_k^k \partial_i \phi \\ &\quad + \alpha e^{-4\phi} (S_{jk} - P \gamma_{jk}) \partial_i \tilde{\gamma}^{jk} / 2 + \alpha \int d\varepsilon S_{(\varepsilon)}^\mu \gamma_{i\mu}], \\ \partial_t \sqrt{\gamma} \tau + \partial_i \sqrt{\gamma} (\tau v^i + P (v^i + \beta^i)) \\ &= \sqrt{\gamma} [\alpha K S_k^k / 3 + \alpha e^{-4\phi} (S_{ij} - P \gamma_{ij}) \tilde{A}^{ij} \\ &\quad - S_i D^i \alpha + \alpha \int d\varepsilon S_{(\varepsilon)}^\mu u_\mu], \\ \partial_t (\rho_* Y_e) + \partial_i (\rho_* Y_e v^i) &= \sqrt{\gamma} \alpha m_e \int \frac{d\varepsilon}{\varepsilon} (S_{(v_e, \varepsilon)}^\mu - S_{(v_e, \varepsilon)}^\mu) u_\mu \end{aligned}$$

Table 1
The Opacity Set Included in this Study and their References

Process	Reference
$\nu_e \leftrightarrow e^- p$	Bruenn (1985), Rampp & Janka (2002)
$p \bar{\nu}_e \leftrightarrow e^+ n$	Bruenn (1985), Rampp & Janka (2002)
$\nu_e A \leftrightarrow e^- A'$	Bruenn (1985), Rampp & Janka (2002)
$\nu p \leftrightarrow \nu p$	Bruenn (1985), Rampp & Janka (2002)
$\nu n \leftrightarrow \nu n$	Bruenn (1985), Rampp & Janka (2002)
$\nu A \leftrightarrow \nu A$	Bruenn (1985), Rampp & Janka (2002)
$\nu e^\pm \leftrightarrow \nu e^\pm$	Bruenn (1985)
$e^- e^+ \leftrightarrow \nu \bar{\nu}$	Bruenn (1985)
$NN \leftrightarrow \nu \bar{\nu} NN$	Hannestad & Raffelt (1998)

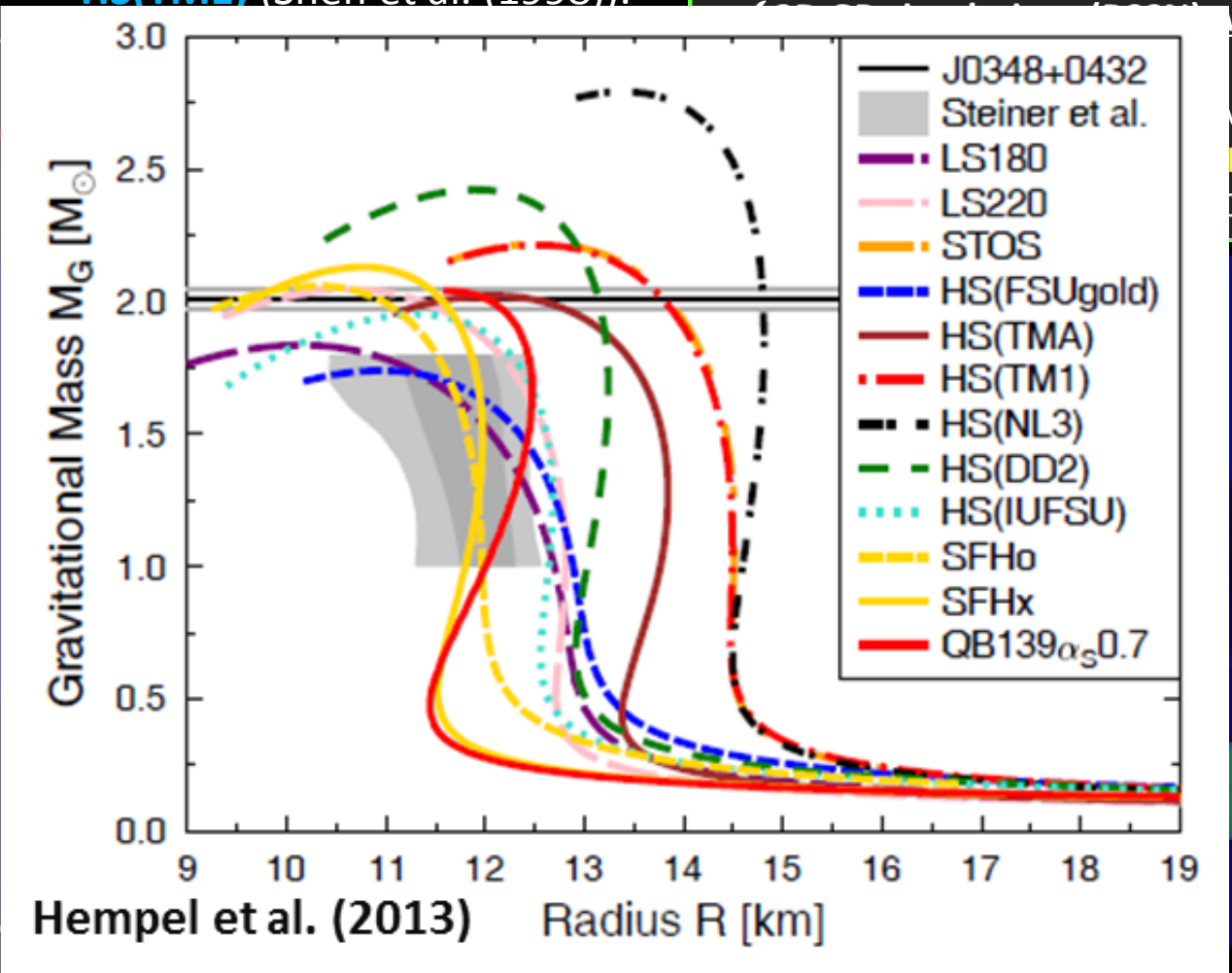
✓ 3 flavor neutrino transport
✓ Base-line opacity (t.b.updated)

GW signatures from 3D-GR models with strong SASI vs. weak SASI activity

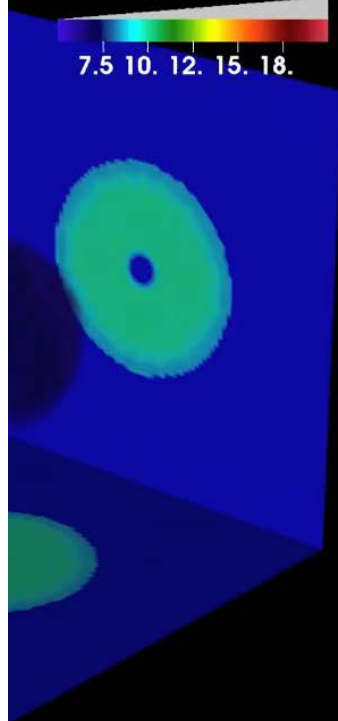
(from Kuroda, KK, & Takiwaki ApJL (2016), see also Andresen, B, E Müller and Janka (2017))

- ✓ Two EOSs → **SFHx** (Steiner et al. (2013), fits well with experiment/NS radius, Steiner+(2011)), **HS(TM1)** (Shen et al. (1998)).

✓ 15 M_{sun} S
 S
 Tpb(ms)=-0



with neutrino
 pJ, 2014, PRD)
 ble !
 016))



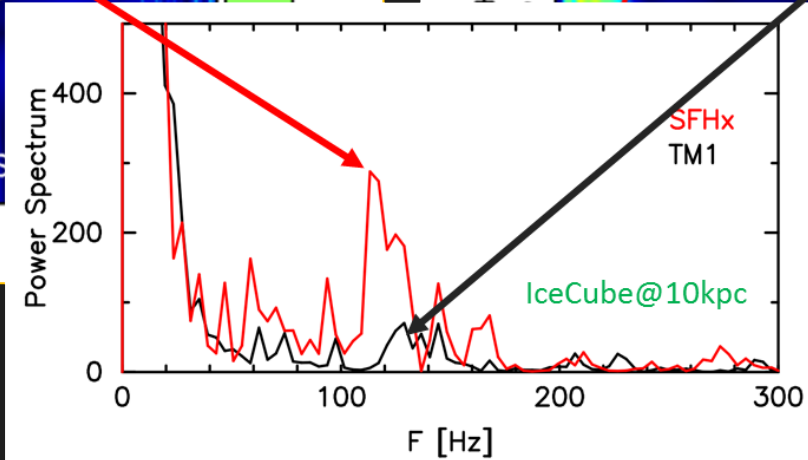
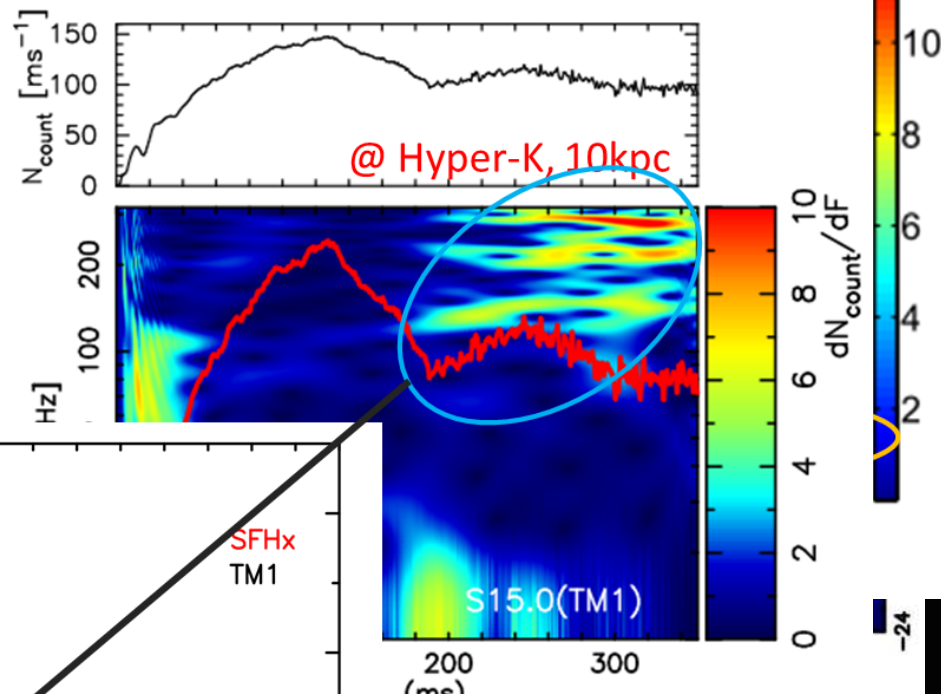
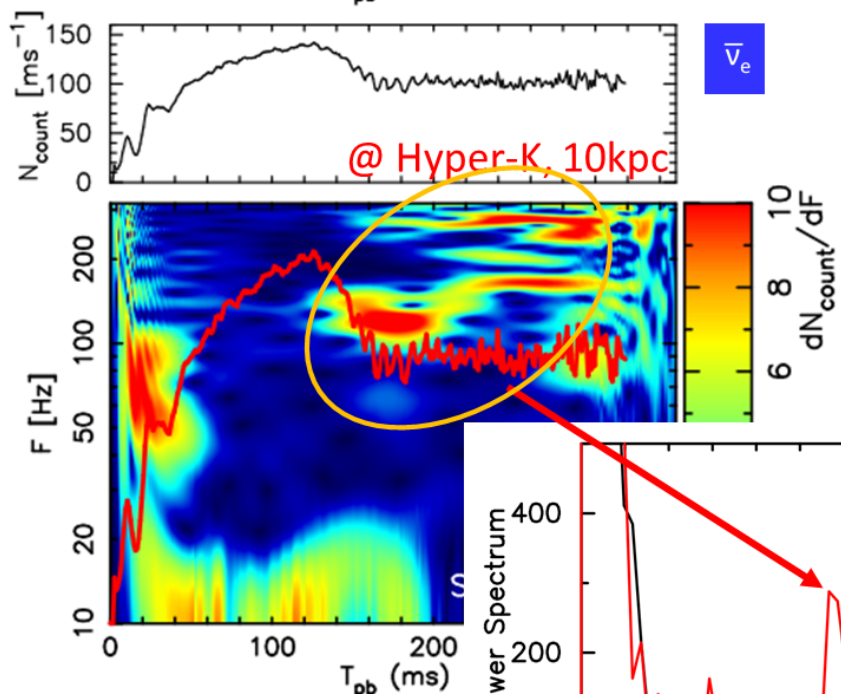
✓ **SASI activity higher for softer EOS** (due to shorter growth rate, e.g., Foglizzo et al. ('06)).

Reconstructed Spectrogram

7

SFHx :softer

TM1 :stiffer



$$T_{\text{SASI}} = \tau_{\text{adv}} + \tau_{\text{ac}} = \int_{r_V}^{r_{\text{sh}}} \frac{dr}{|v_r|} + \int_{r_V}^{r_{\text{sh}}} \frac{dr}{c_s - |v_r|}$$

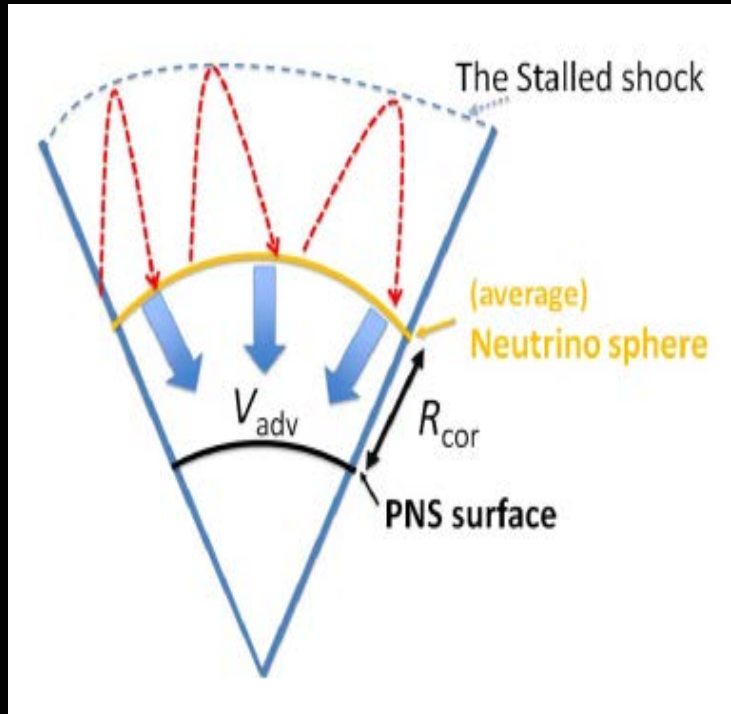
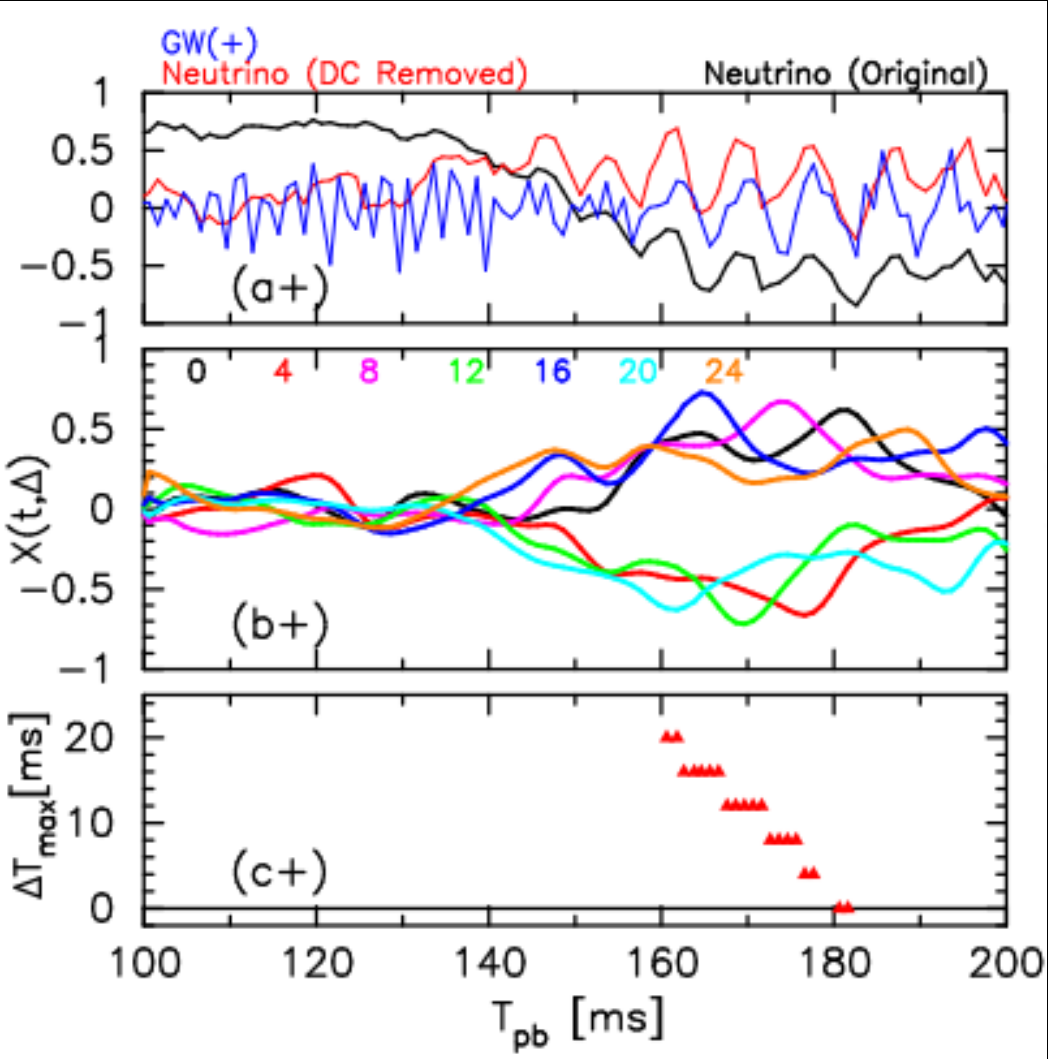
is only 2~3 kpc, but could extend out to 100 kpc when **FT and CE are on-line (>2035)**.

- ✓ **Detection of neutrinos (Super-K, IceCube) important to get timestamp of GW detection.**
- ✓ **The SASI activity, if very high, results in characteristic signatures in both GWs and neutrino signals (e.g., Tamborra et al. (2012) for SASI-induced neutrino signals).**

Correlation between GWs and neutrinos with strong SASI activity ($15 M_{\text{sun}} + \text{SFHx}$)

$$X(t, \Delta T) = \frac{\int d\tau H(t - \tau) A_\nu(\tau + \Delta T) A_{\text{GW}}(\tau)}{\sqrt{\int d\tau H(t - \tau) (A_\nu(\tau + \Delta T))^2} \sqrt{\int d\tau H(t - \tau) (A_{\text{GW}}(\tau))^2}}$$

Kuroda, KK, Hayama, Takiwaki
(2017, ApJ)



$$\begin{aligned} \Delta T_{\text{max}} &\sim R_{\text{cor}} / V_{\text{adv}} \\ &\sim \text{O}(10) \text{ km} / (10^{7\sim 8} \text{ cm/s}) \\ &\sim \text{O}(10) \text{ ms} \end{aligned}$$

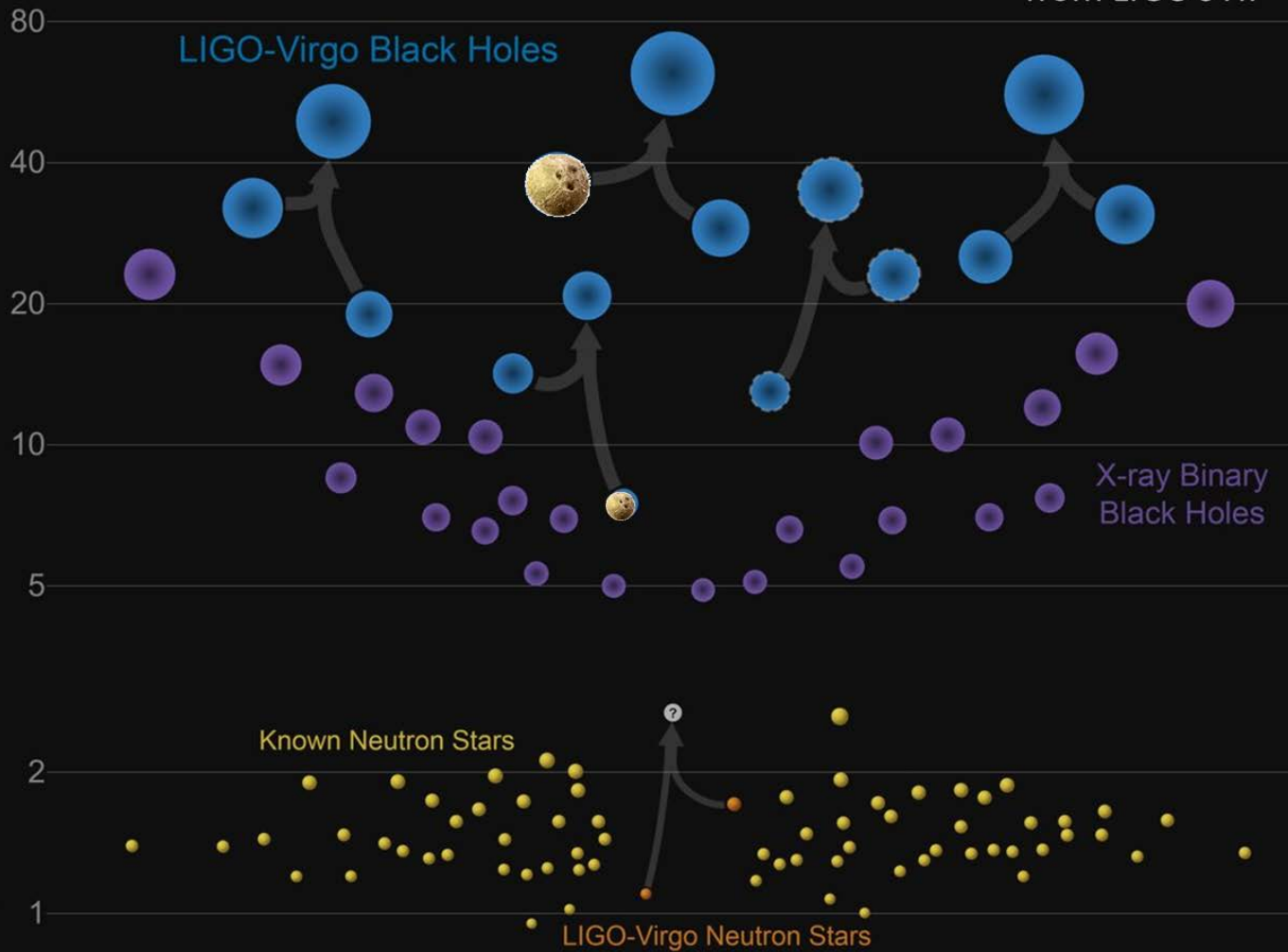
✓ The simultaneous detection potentially tells the distance between the neutrino sphere and PNS radius! (Need to follow long-term 3D evolution how long this continues..)

The Origin of the Nobel-Prize-awarded BHs ($7 \sim 40 M_{\text{sun}}$) ?

Masses in the Stellar Graveyard

in Solar Masses

from LIGO's HP



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201



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Elmehed
Rainer We
Prize share:

The Nob
Rainer V
Thorne
observa

✓ Lo
ne
BH
(e.g., Inagawa et al. (2017/2018))

56

57

51

51

41

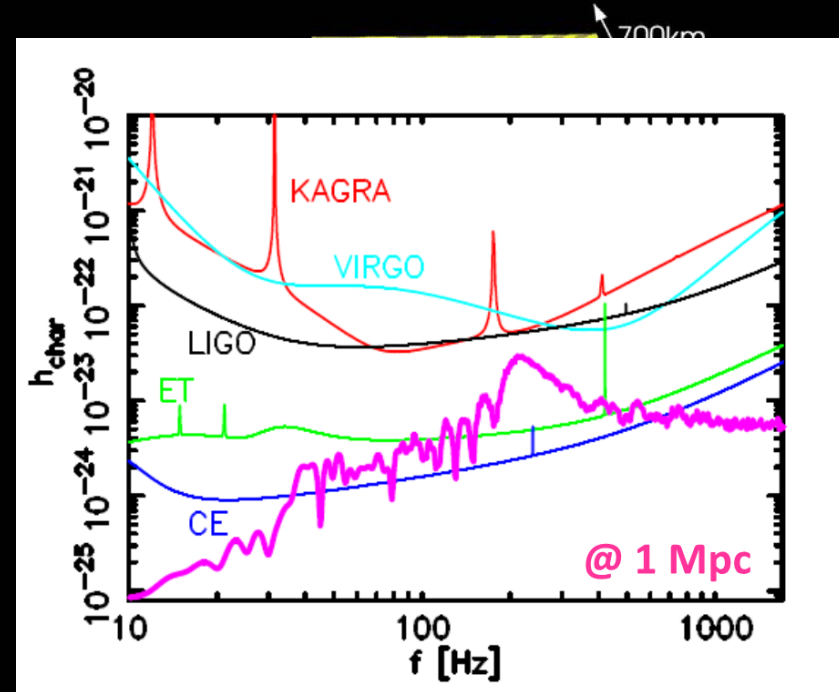
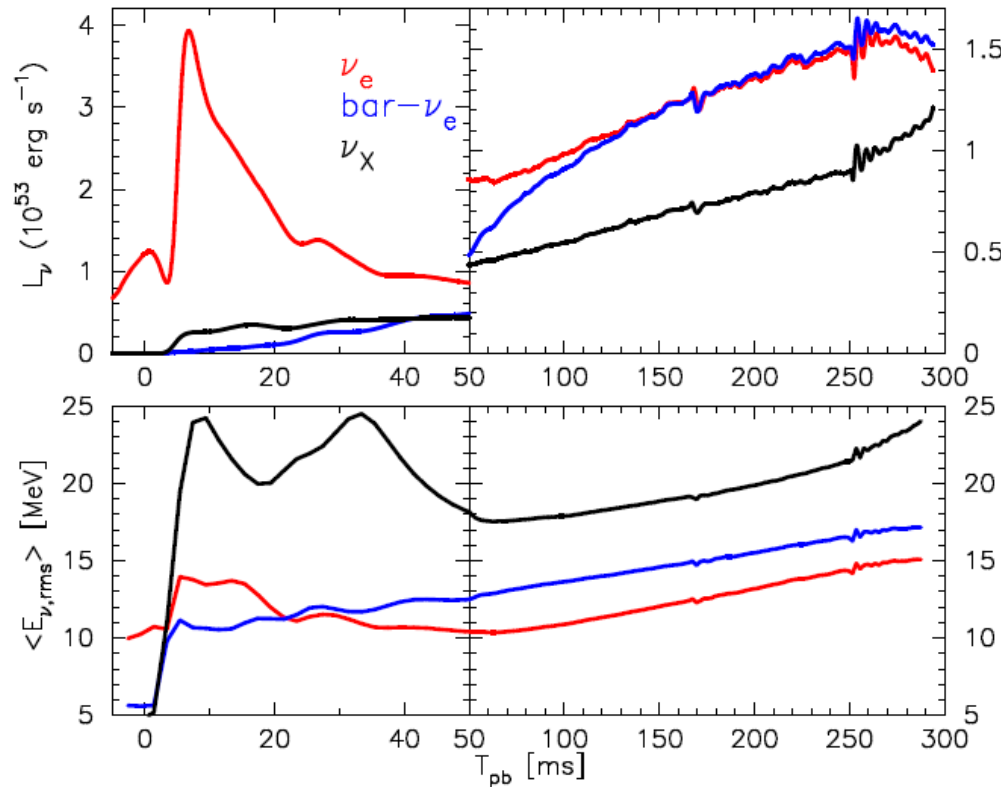
2600 Myr

al. (2006)

✓ FUGRA results of $70 M_{\text{sun}}$ ($M_{\text{CO}} \sim 28.5 M_{\text{sun}}$) (progenitor from Takahashi et al. (2014))

Z70.0(LS220)
Tpb(ms)=-2.59983

S11.2(LS220)
Tpb(ms)=-33.8579



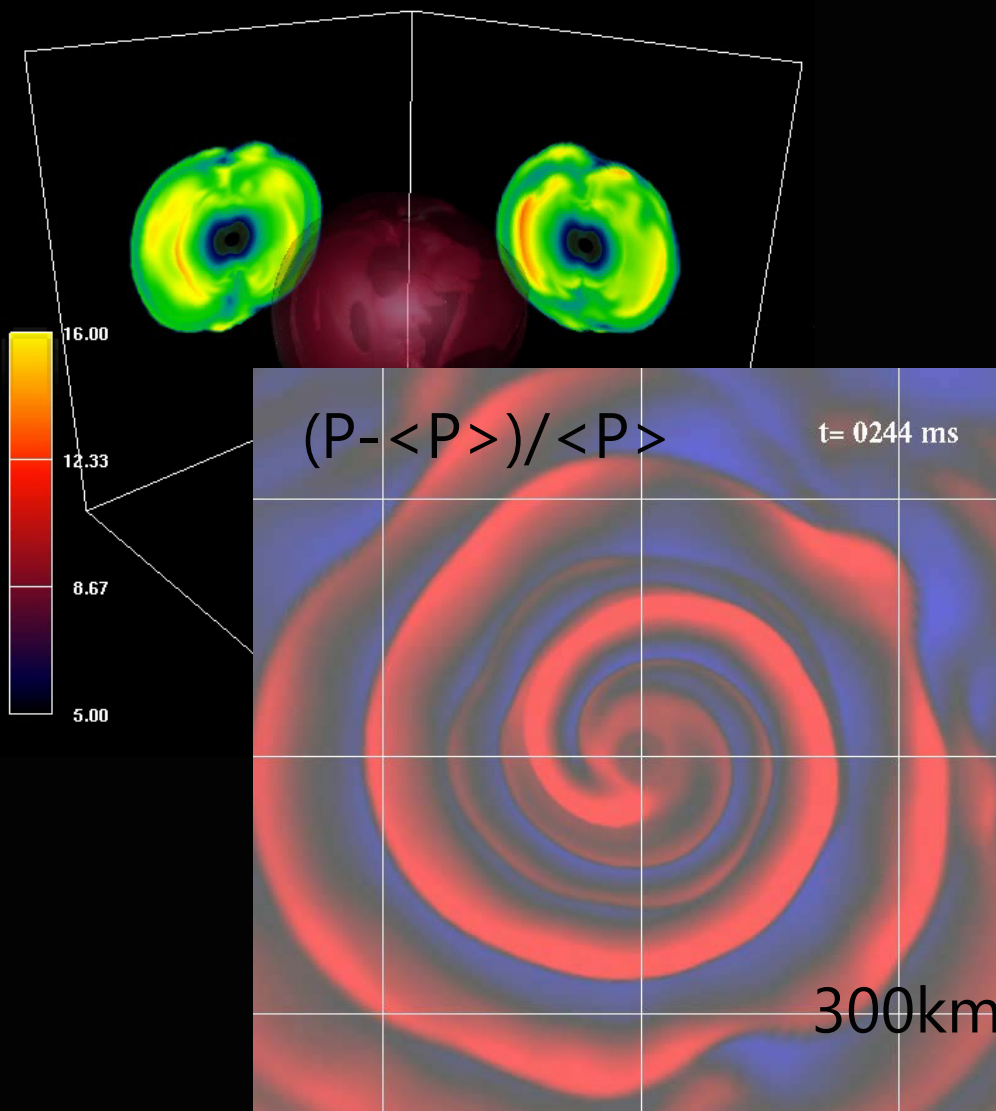
- ✓ **Earliest BH formation** after bounce (~ 300 ms postbounce) !
- ✓ Before the BH formation, **monotonic increase** of neutrino luminosity and rms energy. (**consistent with** 1D, e.g., Sumiyoshi+ (2006), Nakazato(+2008,2013), Fischer+ (2009), Huedepohl+(2016))
- ✓ Strong GW emission is visible to 1 Mpc, **but not** O(100) Mpc...
- ✓ Our code needs upgrade to follow long after BH formation...

Switching gears to MHD mechanism (rapid rotation required !!)

3D rotating explosion simulation of a $27 M_{\text{sun}}$ star ($\Omega_0 = 2 \text{ rad/s}$) with IDSA.
(Takiwaki, KK, and Suwa, MNRAS Letters, (2016), see also Summa et al. (2017)).

Entropy

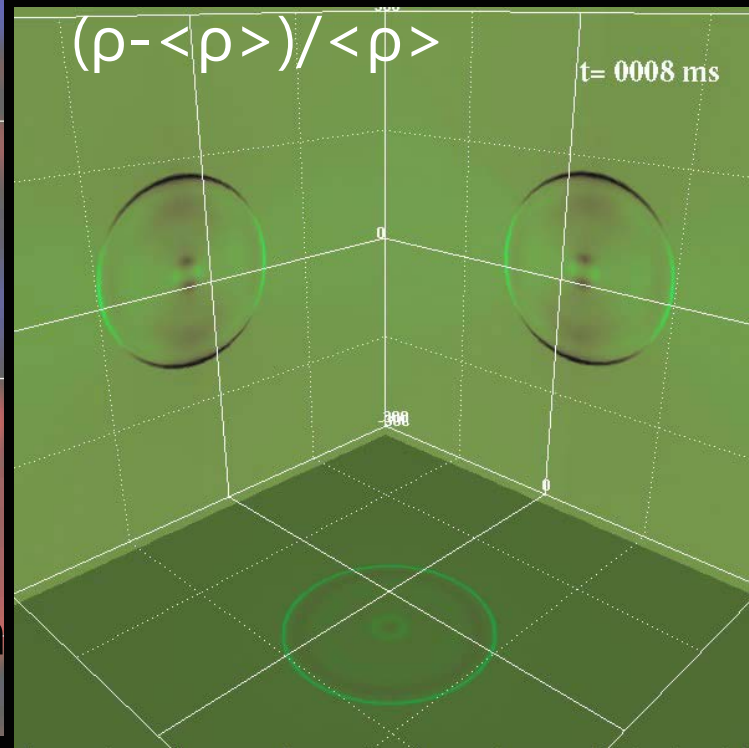
$t = 0102 \text{ ms}$



✓ One-armed (low $T/|W|$) instability



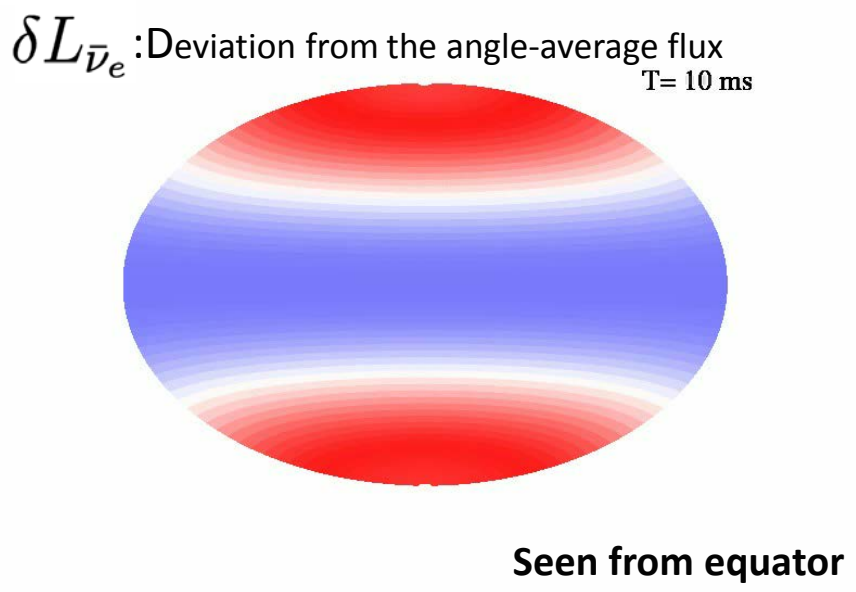
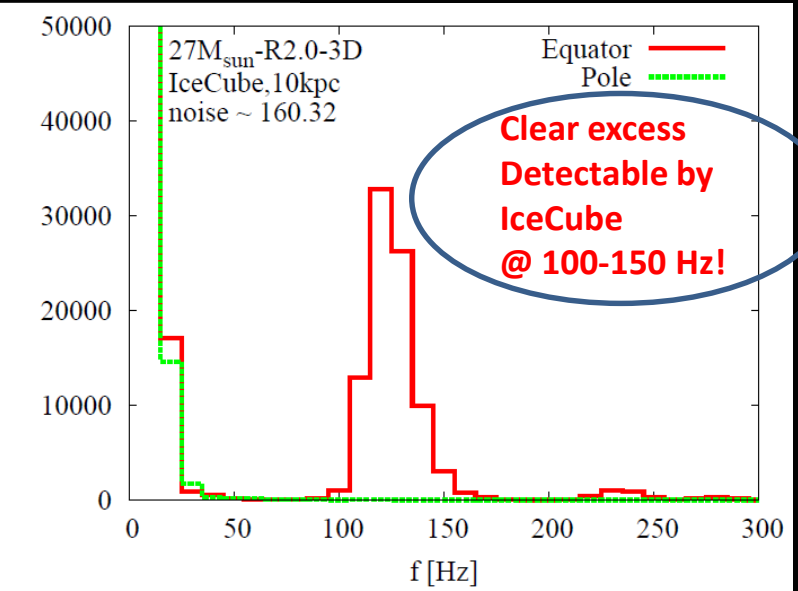
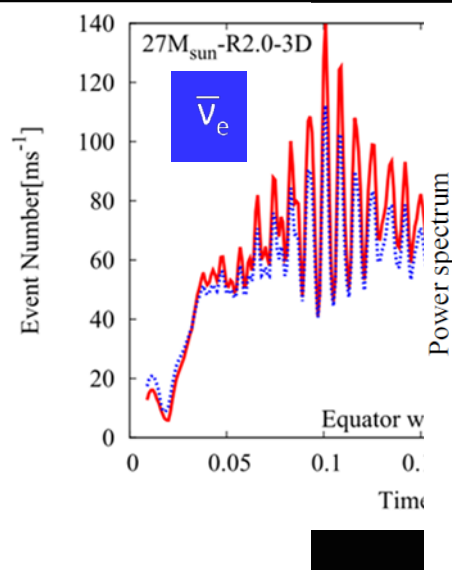
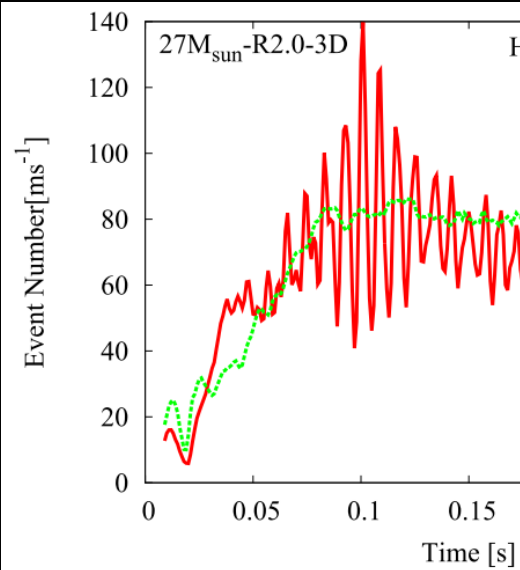
✓ Spiral waves enhance energy transport from PNS to gain region !



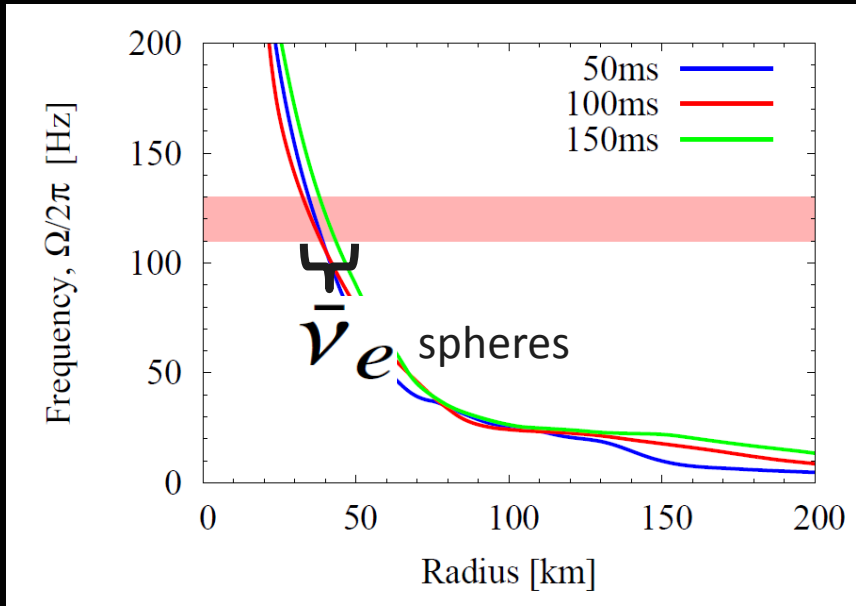
Neutrino signatures from rapidly rotating explosion of $27 M_{\text{sun}}$ star

Takiwaki and KK (2018), MNRAS Letters

Quasi-periodic variation ! May survive with coll. oscillation



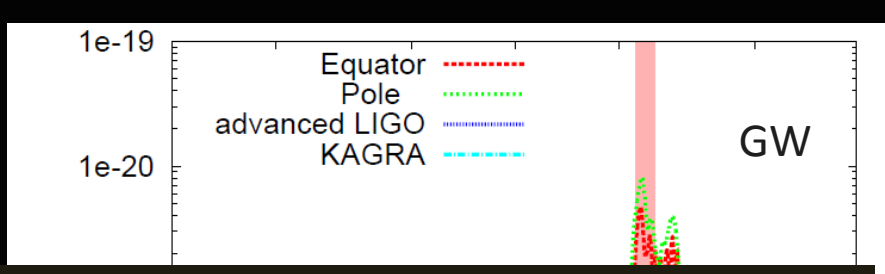
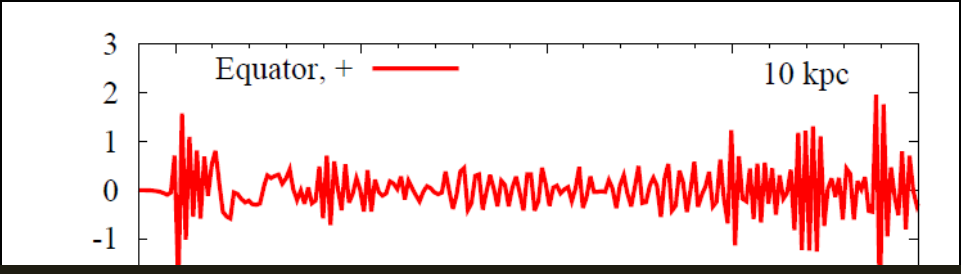
"Lighthouse effect"



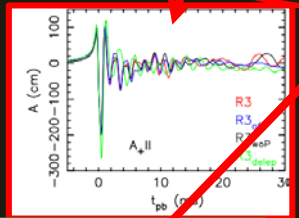
Correlation of GW and neutrino signatures from the 3D rotating model,

Gravitational waveform ($27 M_{\text{sun}}, \Omega_0 = 2\text{rad/s}$)

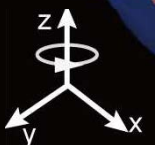
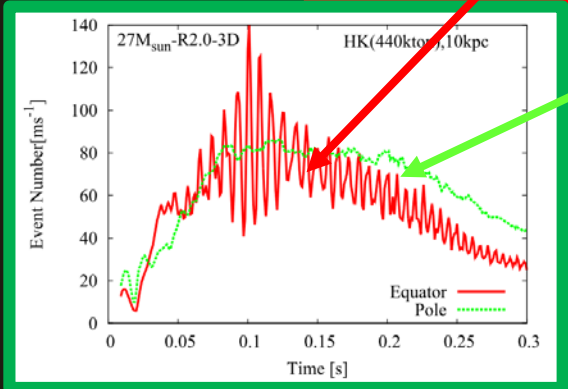
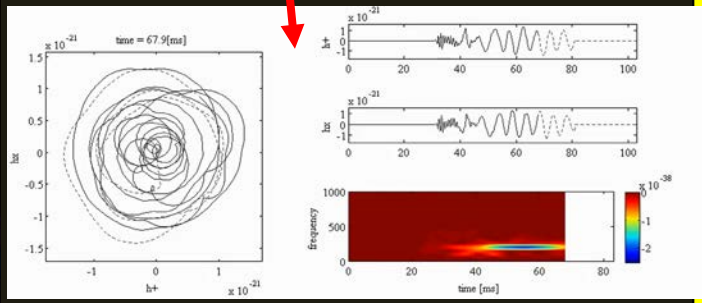
Takiwaki and KK (2018), MNRAS Letters



Directionality	Equator	Pole
Gravitational Wave	Type I signal	<ul style="list-style-type: none"> ✓ Quasi-periodic signals from non-axis. instability ✓ Circular polarization
Neutrinos	Light-house effect	No surprise ...



150ms



Need improvement in opacity of our 3D-GR code (with energy transport)!

Table 1

The Opacity Set Included in this Study and their References

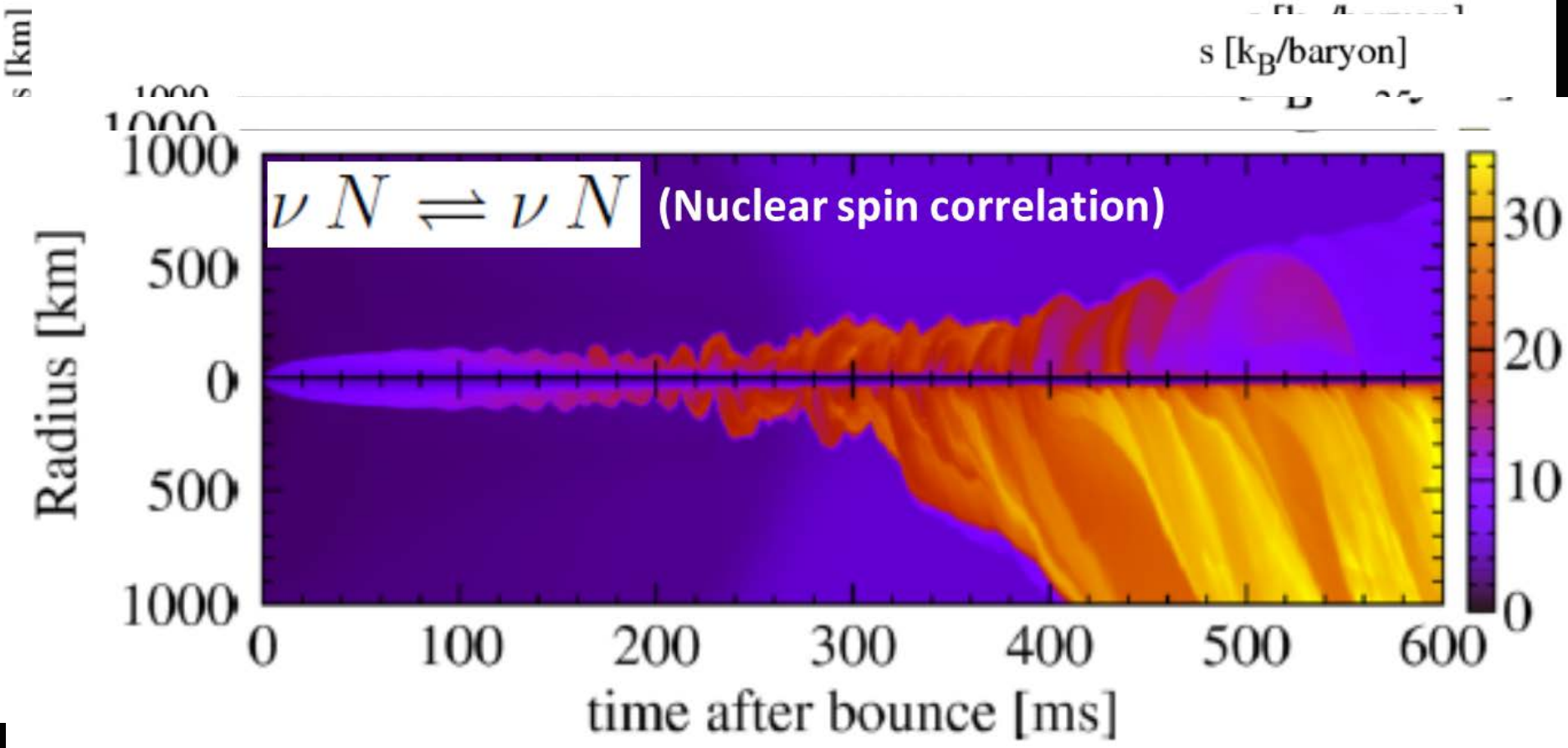
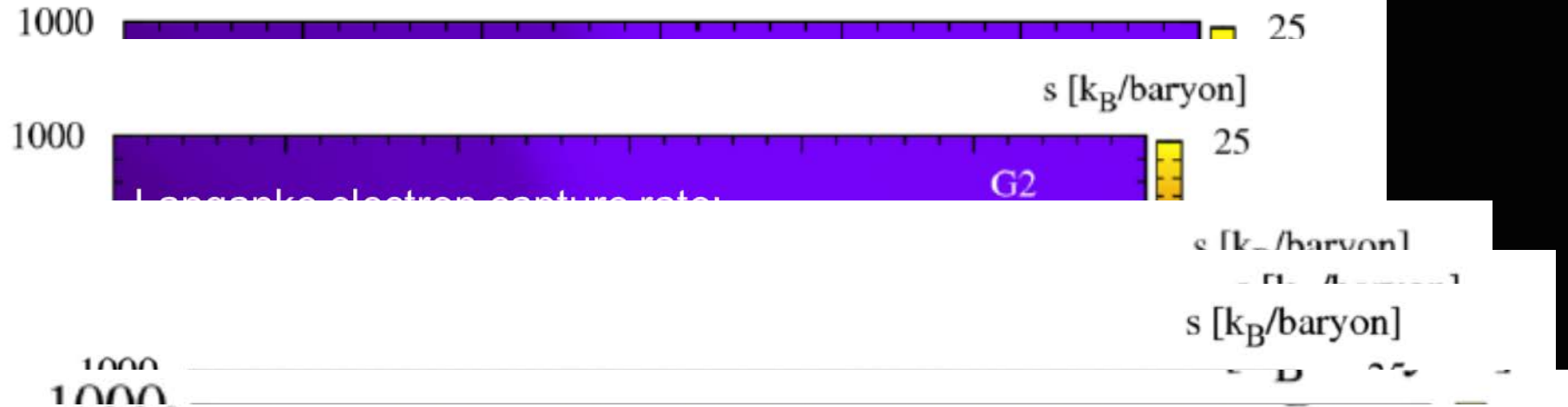
Process	Reference	Summarized In
$n\nu_e \leftrightarrow e^-p$	Bruenn (1985), Rampp & Janka (2002)	Appendix A.1
$p\bar{\nu}_e \leftrightarrow e^+n$	Bruenn (1985), Rampp & Janka (2002)	Appendix A.1
$\nu_e A \leftrightarrow e^-A'$	Bruenn (1985), Rampp & Janka (2002)	Appendix A.1
$\nu p \leftrightarrow \nu p$	Bruenn (1985), Rampp & Janka (2002)	Appendix A.2
$\nu n \leftrightarrow \nu n$	Bruenn (1985), Rampp & Janka (2002)	Appendix A.2
$\nu A \leftrightarrow \nu A$	Bruenn (1985), Rampp & Janka (2002)	Appendix A.2
$\nu e^\pm \leftrightarrow \nu e^\pm$	Bruenn (1985)	Appendix A.3
$e^-e^+ \leftrightarrow \nu\bar{\nu}$	Bruenn (1985)	Appendix A.4
$NN \leftrightarrow \nu\bar{\nu}NN$	Hannestad & Raffelt (1998)	Appendix A.5

KTK (2016), ApJS
(essentially,
Bruenn rates +
Bremsstrahlung)

Most advanced set
(e.g., Fischer(2016),
Bollig et al. (2017))

	Weak process	References
1	$e^- + p \rightleftharpoons n + \nu_e$	Reddy et al. (1998); Horowitz (2002)
2	$e^+ + n \rightleftharpoons p + \bar{\nu}_e$	Reddy et al. (1998); Horowitz (2002)
3	$n \rightleftharpoons p + e^- + \bar{\nu}_e$	Fischer et al. (2016b)
4	$e^- + (A, Z) \rightleftharpoons (A, Z - 1) + \nu_e$	Juodagalvis et al. (2010)
5	$\nu + N \rightleftharpoons N + \nu'$	Bruenn (1985); Mezzacappa & Bruenn (1993a); Horowitz (2002)
6	$\nu + (A, Z) \rightleftharpoons (A, Z) + \nu'$	Bruenn (1985); Mezzacappa & Bruenn (1993a)
7	$\nu + e^\pm \rightleftharpoons e^\pm + \nu'$	Bruenn (1985); Mezzacappa & Bruenn (1993b)
8	$e^- + e^+ \rightleftharpoons \nu + \bar{\nu}$	Bruenn (1985)
9	$N + N \rightleftharpoons N + N + \nu + \bar{\nu}$	Hannestad & Raffelt (1998)
10	$\nu_e + \bar{\nu}_e \rightleftharpoons \nu_{\mu/\tau} + \bar{\nu}_{\mu/\tau}$	Buras et al. (2003); Fischer et al. (2009)
11	$(A, Z)^* \rightleftharpoons (A, Z) + \nu + \bar{\nu}$	Fuller & Meyer (1991); Fischer et al. (2013)

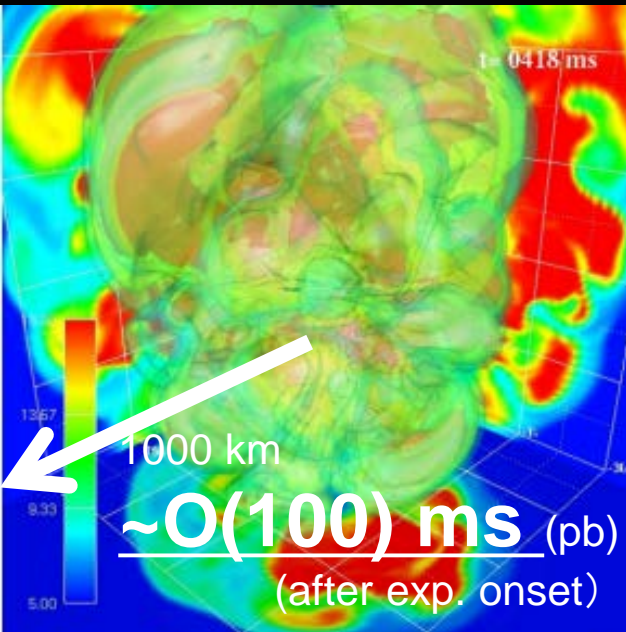
Note: unless stated otherwise, $\nu = \{\nu_e, \bar{\nu}_e, \nu_{\mu/\tau}, \bar{\nu}_{\mu/\tau}\}$ and $N = \{n, p\}$.



✓ Quantitative GW-neutrino signal prediction, the updates in opacities mandatory!

Next 10 years: Where are we and are we going?

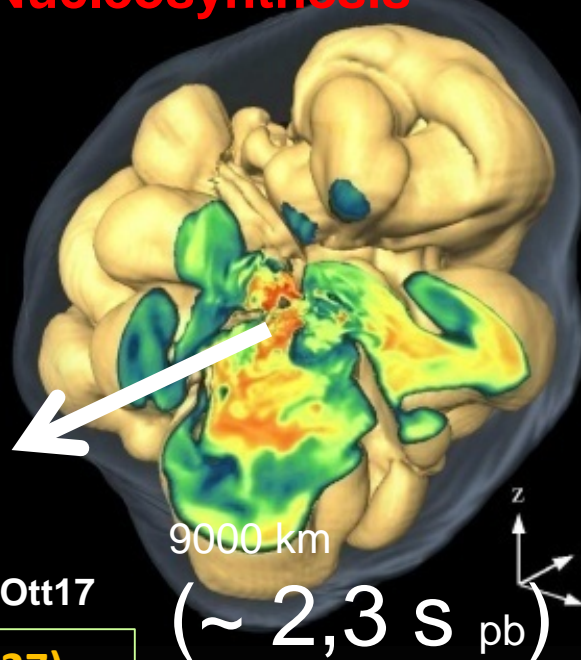
“A” self-consistent 3D model



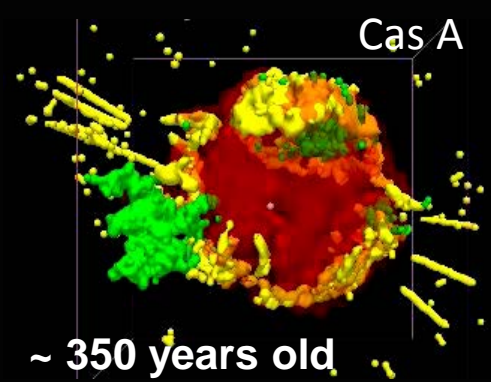
Takiwaki, KK, Suwa (2016), Melson+15, Ott17

**For some progenitors (11.2,15,20.27),
the stalled shock revived !
(5D/4D with approximate transport)**

**Gray-transport simulation
Nucleosynthesis**

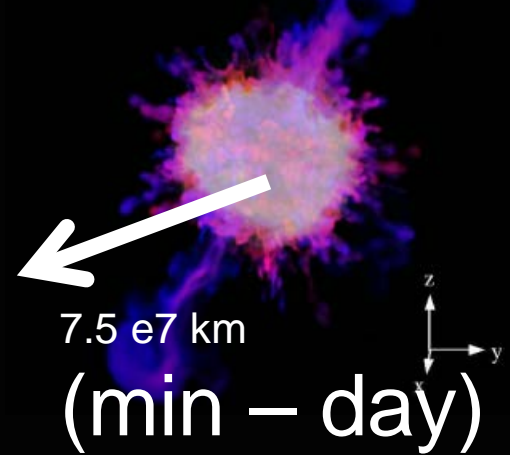


Wongwathanarat et al. (2015)



DeLaney et al. (2010)

**Hydrodynamic model:
Mixing, RT, RM instabilities**



Wongwathanarat et al. (2016)

**To-do-1: Long-term evolution in self-consistent 3D (GR) models
⇒ confront CCSN theory with observation ⇒ Pragmatism**

**To-do-2 : Full GR and Boltzmann project :
⇒ ultimately test whether the stalled shock would revive. ⇒ Perfectionism**

SN 20xx ! in the Galactic center: End-to-End Bridging Simulations

sec min hours day years

- 4 - 3 - 2 - 1 0 2 >3

SK detects ~ 10,000 neutrinos

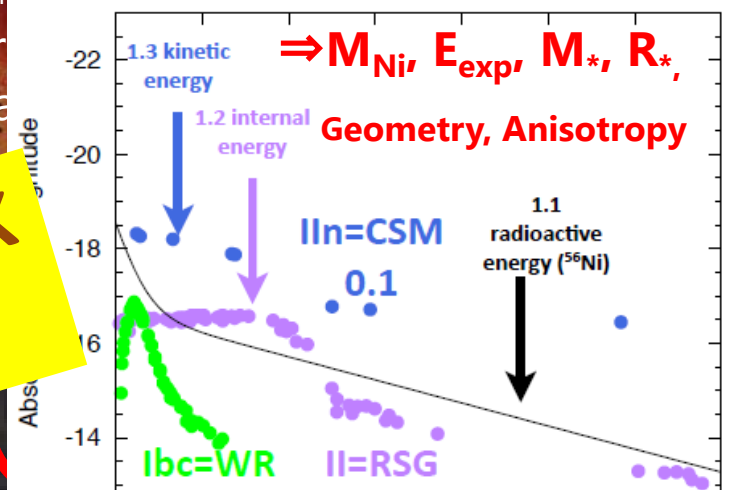
< 15min SURGE meeting (Supernova Urgent Response Group)

< 1 hour SK provide alert: Astronomer
(onset of neutrino burst, duration)

Log (day)

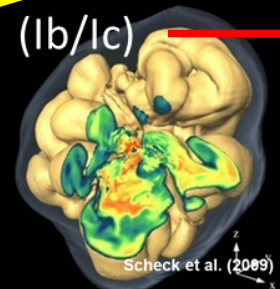
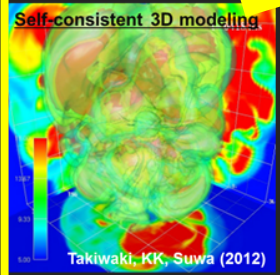


K, SK, KamLAND, LIGO, Virgo
KAGRA, Hyper-K, and Post-K
ET, CE, DUNE.....



Self-Consistent 3D modeling

Self-consistent shock-breakout



Multi-messenger research in steady progress !

Hammer et al. (2011)

~ 350 years old

Summary

1. In 2D, a number of explosion models (> 400) obtained by independent groups. Some are enough energetic to account for observations (E_{exp} , Ni).
2. 3D explosions generally under-energetic than 2D.
 - progenitor dependence yet unclear.
 - ✓ **Need to find some ingredients to foster 3D explosions.**
 - Need neutrino physics update ? (e.g., Melson et al. (2015), KK+(2017))
 - Impact of rotation/magnetic fields needs to be clarified in 3D self-consistent models.
(e.g., Takiwaki and Kotake (2018), Obergaulinger et al. (2016))
3. 3D GR modelling has just started with increasing microphysical inputs.
(e.g., It takes time ... next generation machines needed !)
4. **Multi-messenger analysis of neutrino and GWs are in steady progress.**
 - : provide information to measure “SASI” activity.
 - and to break the degeneracy (M_{PNS} , R_{PNS} , T_{PNS} , R_{shock} , EOS etc.)
 - ⇒ important probe to the explosion physics **for the SN20xx !**

Many thanks!