

超新星ニュートリノ加熱物質中の元素合成

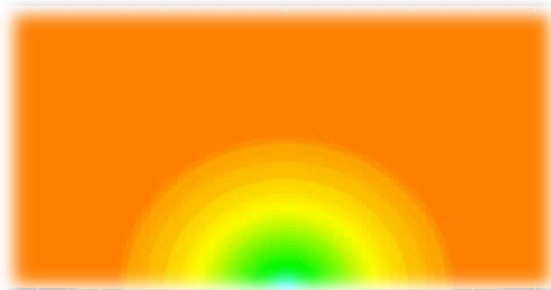
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新学術領域研究会
2015年5月15-17日, 神戸大学



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1. nucleosynthesis during the first 1 second;
what do we see in the innermost SN ejecta?
(Wanajo, Janka, Müller 2011, 2013, and in prep.)
2. nucleosynthesis during the next 10 seconds;
do we see an r-process? (Wanajo 2013)
3. how close to successful supernova explosions?
sensitivity to neutrino opacity (Melson+2015)

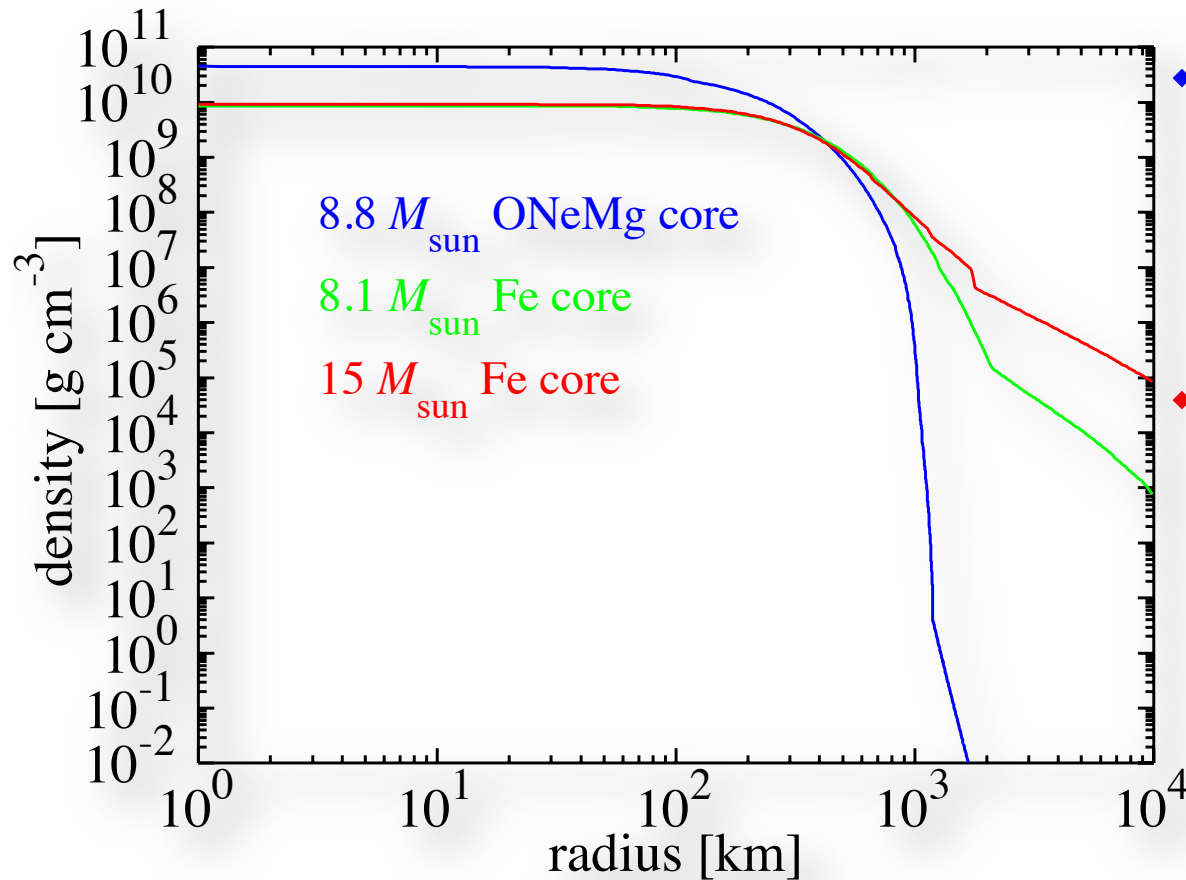


1. nucleosynthesis during the first 1 second;
what do we see in the innermost SN ejecta?
(Wanajo, Janka, Müller 2011, 2013, in prep.)

current status of SN simulations

	explosion ($\sim 10 M_{\odot}$)	explosion ($> 10 M_{\odot}$)
1D	yes	no
2D	yes	yes
3D	yes	no

pre-SN core density profiles

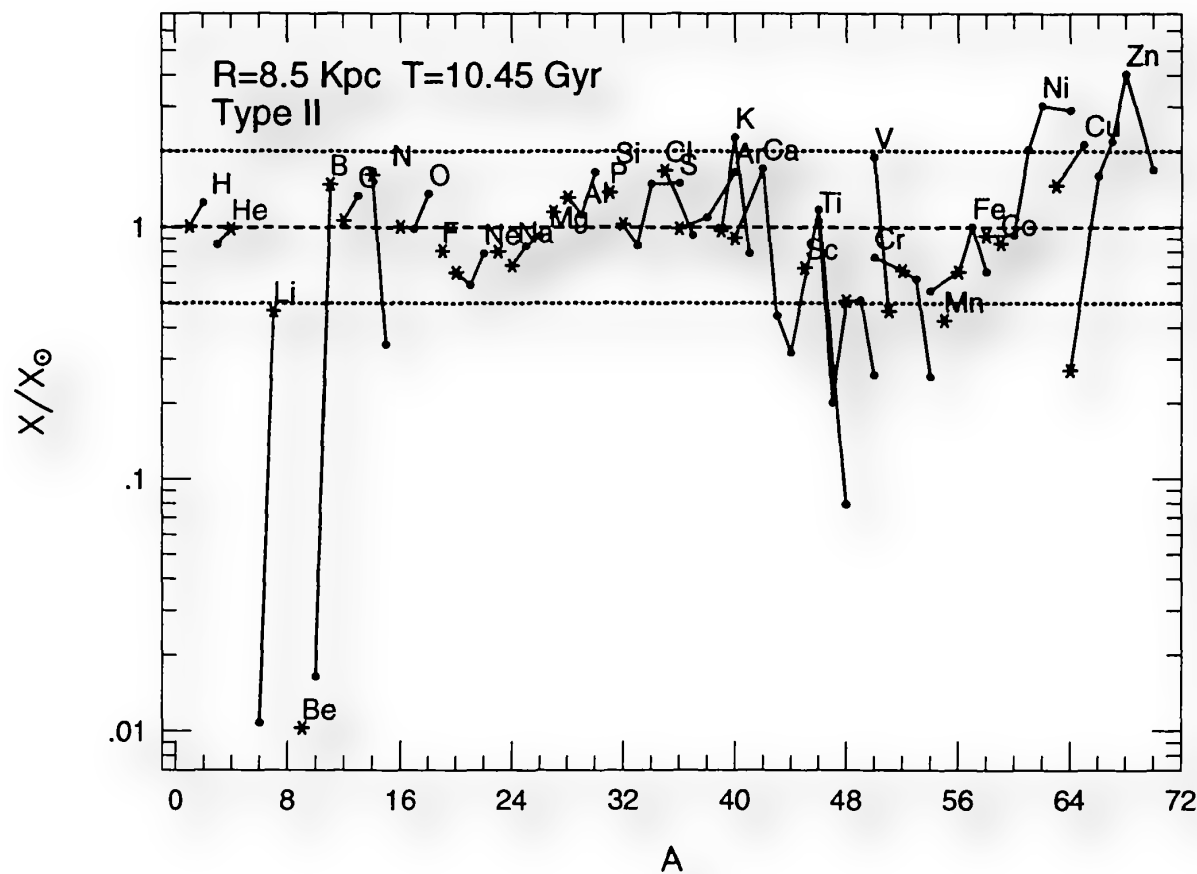


❖ low-mass SNe ($\sim 8-10 M_{\odot}$)
diluted outer envelope
makes explosion easy

❖ massive SNe ($> 10 M_{\odot}$)
dense outer envelope
makes explosion hard

known problems in SN nucleosynthesis

production factors from all SNe; Timmes+1995



previous SN nucleosynthesis based on parameterized 1D models

- ❖ no production of light trans-iron (Zn to Zr; $Z=30-40$, $A = 64-90$) nor r-process elements
- ❖ little production of ^{48}Ca (double magic), ^{64}Zn (main isotope), ^{92}Mo (p-isotope)

what crucial for nucleosynthesis are...

Au 79

Gold

Atomic Radius 174pm
Crystal Structure Face Centered Cubic
Electron Shells [Xe]6s¹4f¹⁴5d¹⁰

Atomic Weight 196.96655
Density 19.3 g/cc
Melting Point 1064.18°C
1947.52°F
Boiling Point 2856°C
5173°F
Electronegativity 2.54

% in Universe 0.000000060%
% in Sun 0.00000010%
% in Crust 0.00000031%
% in Ocean 5.0 × 10⁻⁹%
% in Humans 0.000010%

computational knowledge from
WolframAlpha

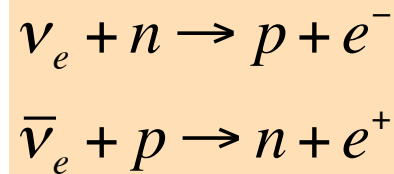
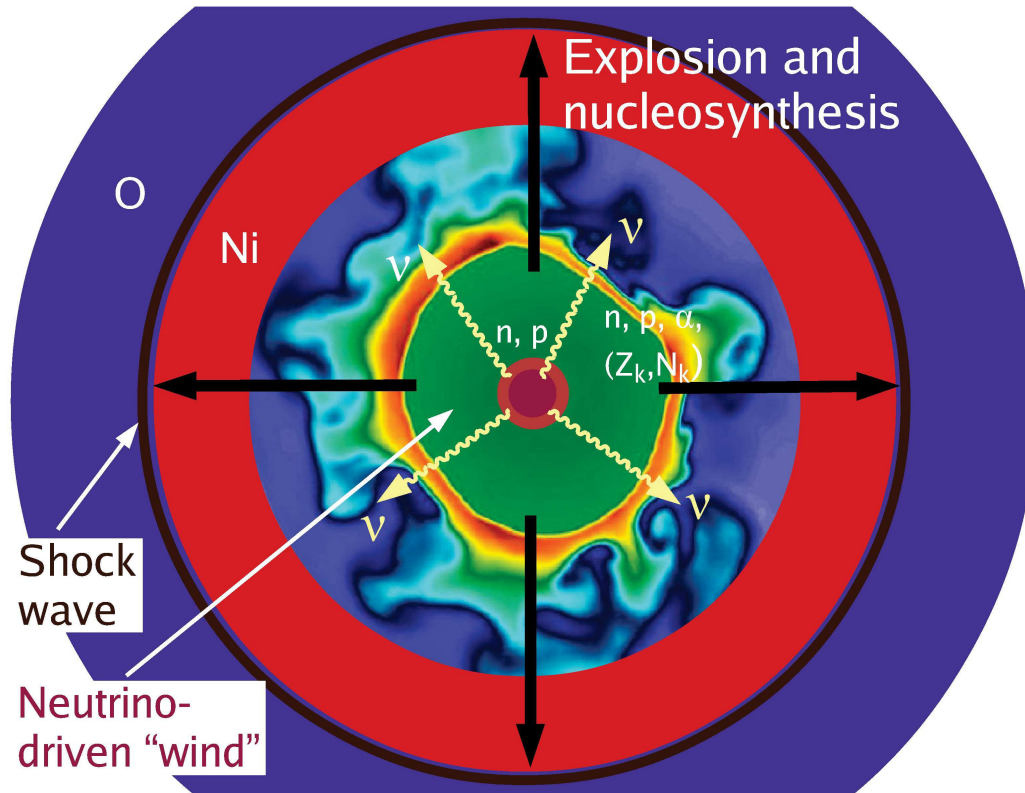
Pt 78 Home Au 79
Back

important physical parameters

- ❖ **entropy; S** ($\propto T^3/\rho$)
controls n , p , α amounts
- ❖ **expansion timescale; τ**
(e -folding time of T)
controls n , p , α amounts
- ❖ **electron fraction; Y_e**
(protons per nucleon)
controls n/p ratios

innermost ejecta of SNe

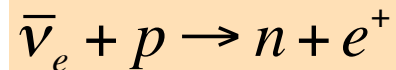
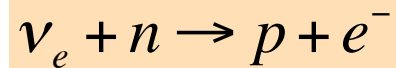
Janka+2012



- ❖ elements up to iron are formed in the outer layers
- ❖ (light) trans-iron elements likely to be formed in the innermost layers
- ❖ difficult to constrain Y_e from simulations because of its sensitivity to neutrino spectra

what determines Y_e ?

❖ Y_e is determined by



❖ equilibrium value is

$$Y_e \sim \left[1 + \frac{L_{\bar{\nu}_e} \varepsilon_{\bar{\nu}_e} - 2\Delta}{L_{\nu_e} \varepsilon_{\nu_e} + 2\Delta} \right]^{-1},$$

$$\Delta = M_n - M_p \approx 1.29 \text{ MeV}$$

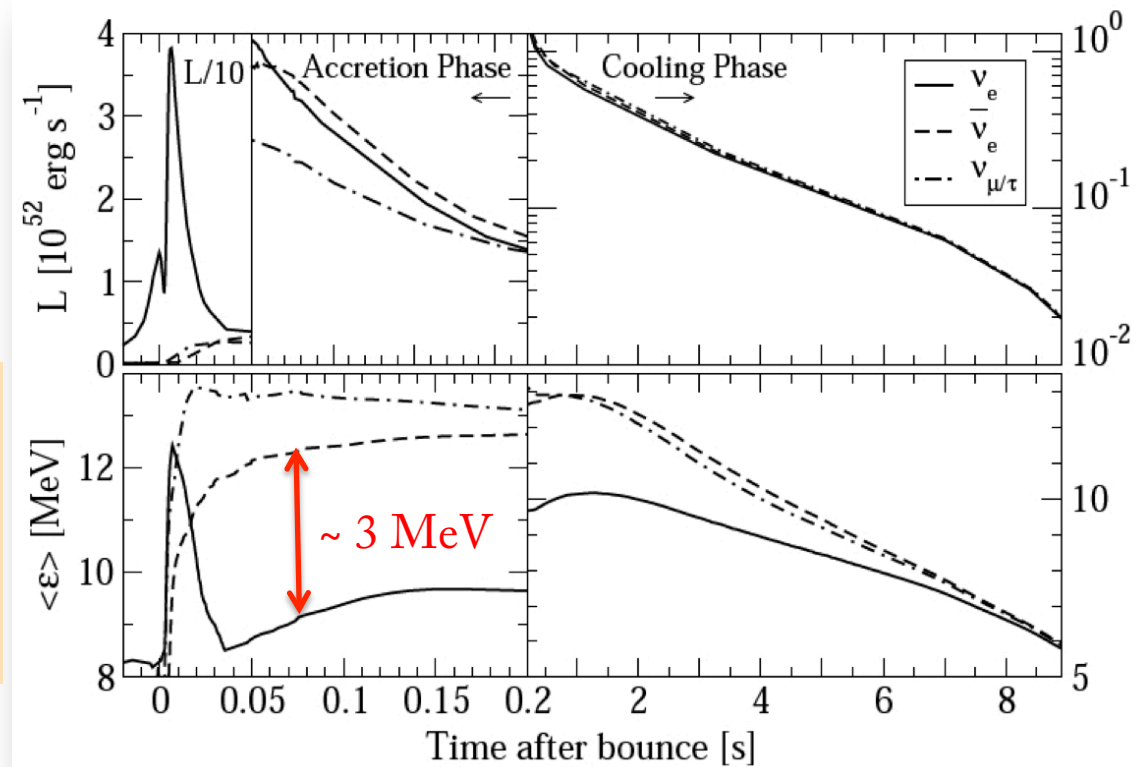
❖ for $Y_e < 0.5$ (i.e., n-rich)

$$\varepsilon_{\bar{\nu}_e} - \varepsilon_{\nu_e} > 4\Delta \sim 5 \text{ MeV}$$

$$\text{for } L_{\bar{\nu}_e} \approx L_{\nu_e}$$

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1D self-consistent explosion of a $9 M_\odot$ star
Hüdepohl+2009

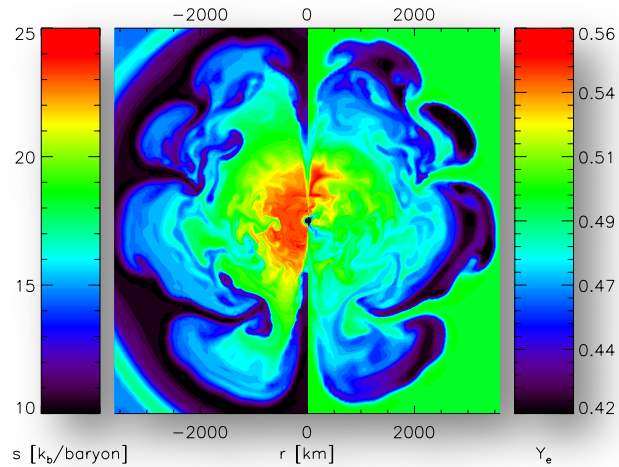


neutrino spectra in the early times are crucial for nucleosynthesis !!!

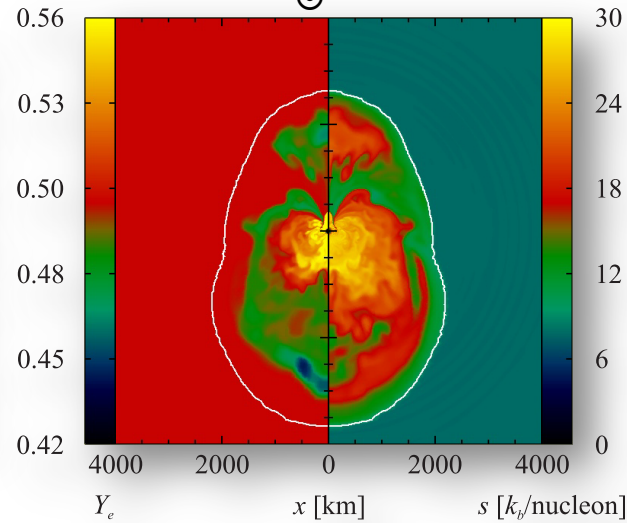
和南城伸也

2D SN simulations with ν -transport

8.8 M_{\odot} ECSN

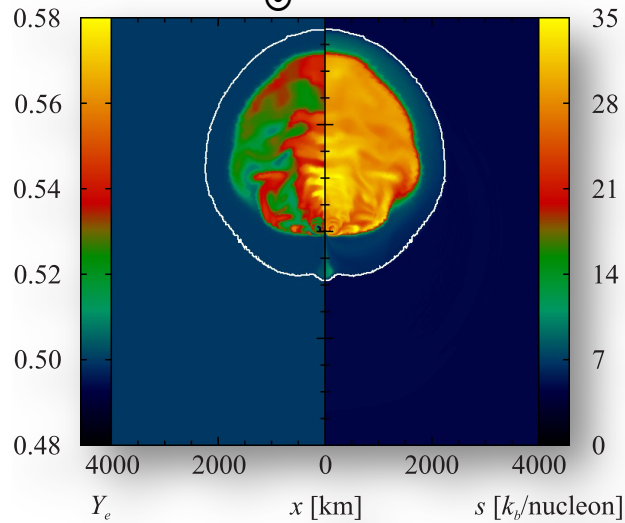


8.1 M_{\odot} CCSN

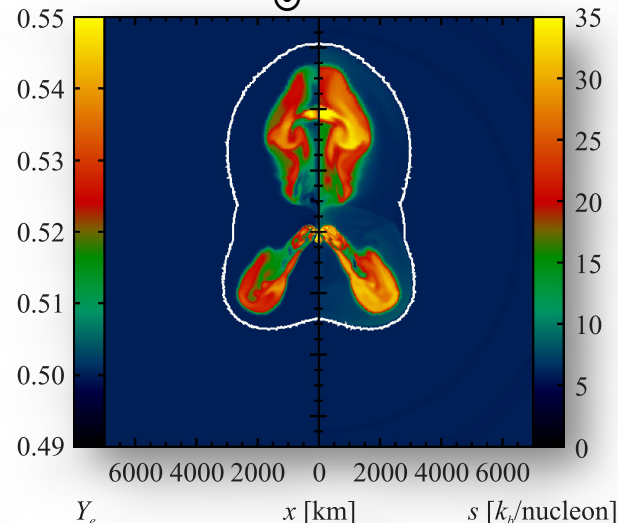


❖ a number of self-consistent SN models with neutrino transport are now available

15 M_{\odot} CCSN



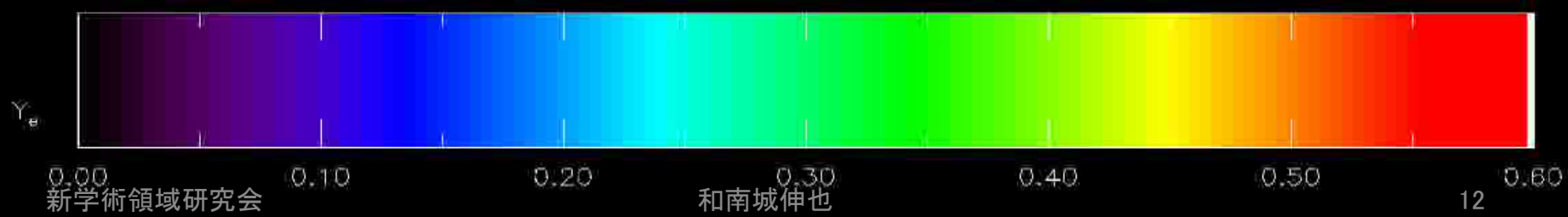
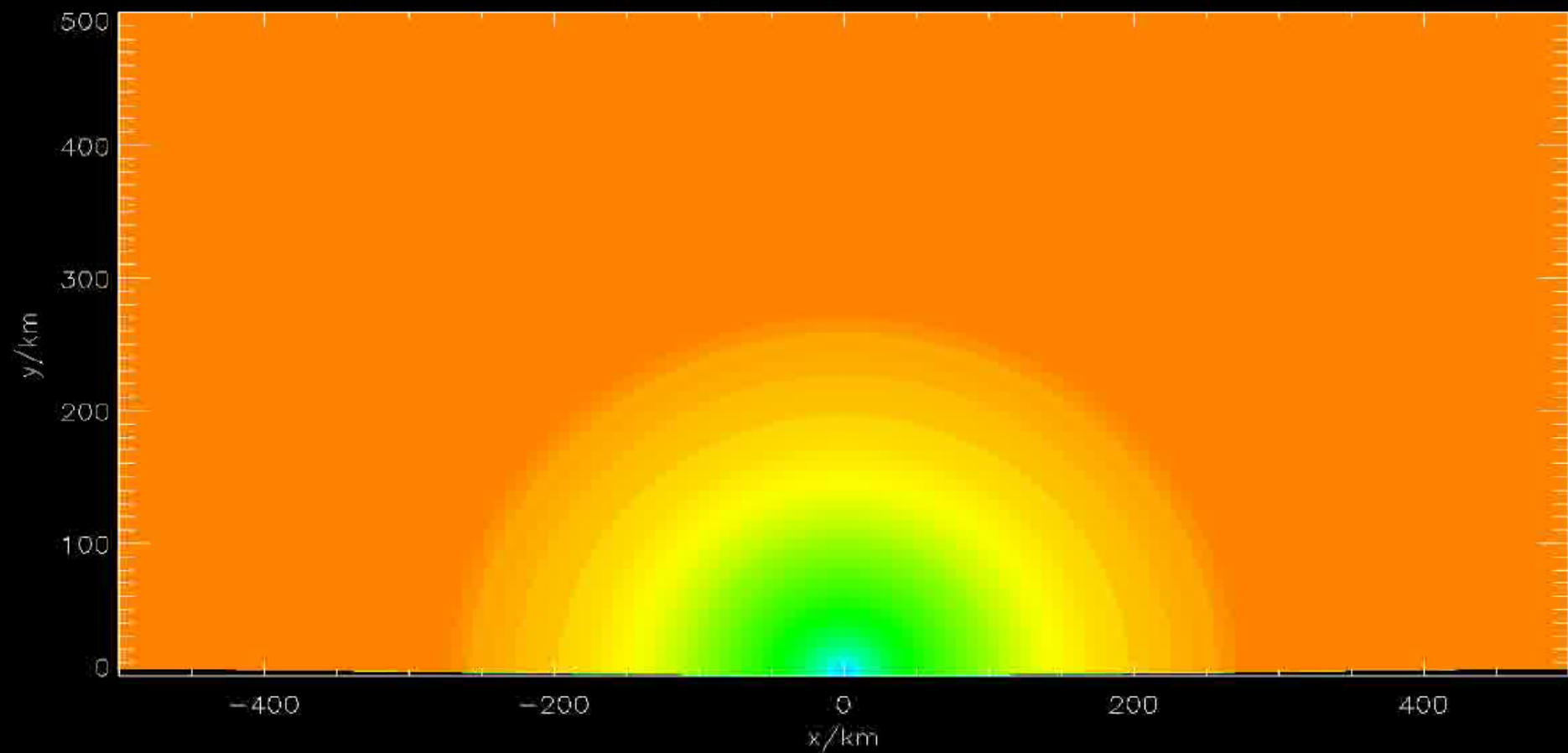
27 M_{\odot} CCSN



❖ very first result of SN nucleosynthesis with such models

8.8 M_{\odot} self-consistently exploding ONeMg core supernova

simulation by Bernhard Müller

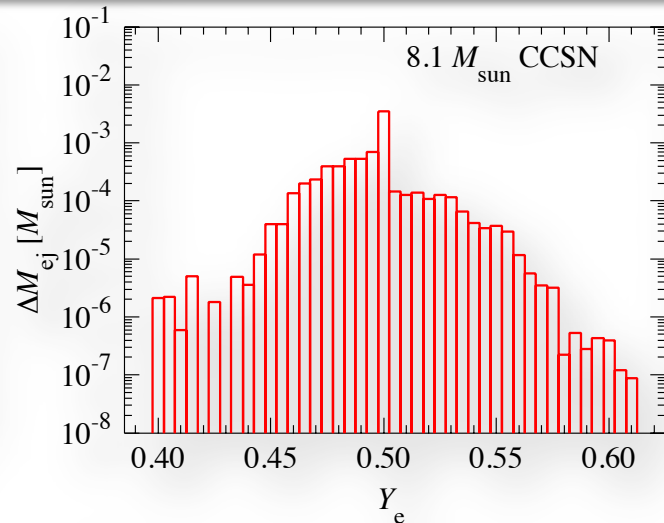
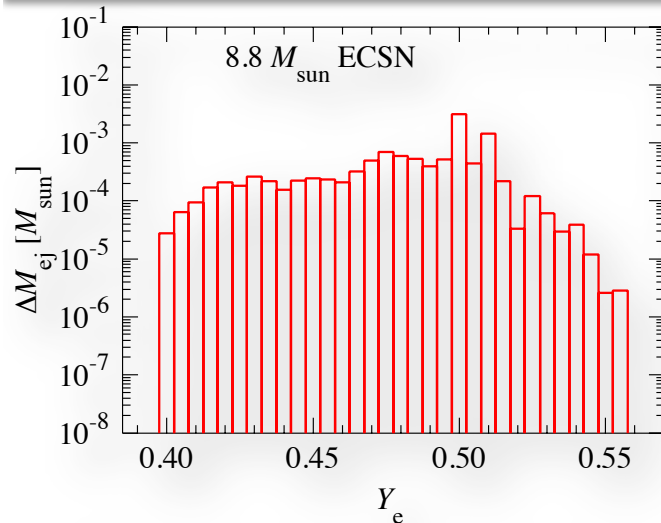


**27 M_{\odot} self-consistently exploding
Fe core**

simulation by Bernhard Müller

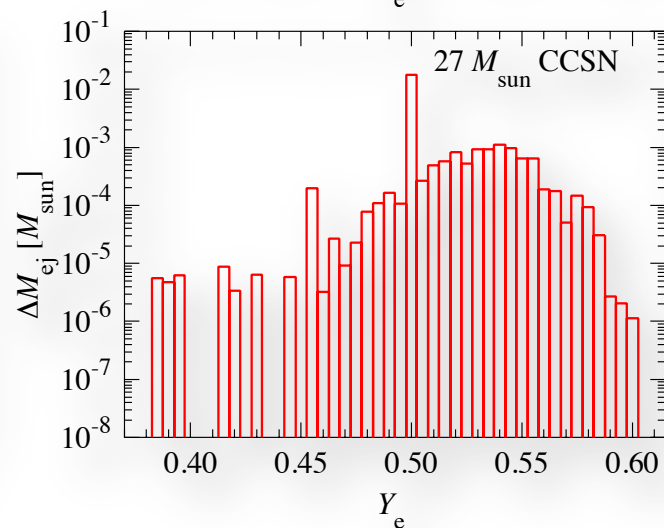
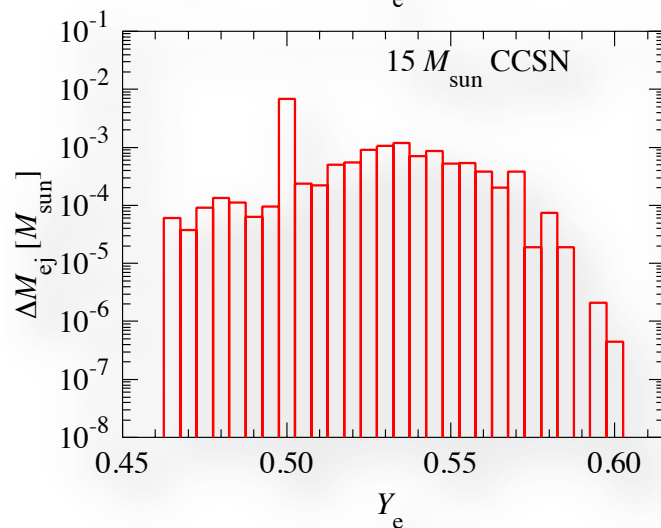


neutron-richness in the ejecta



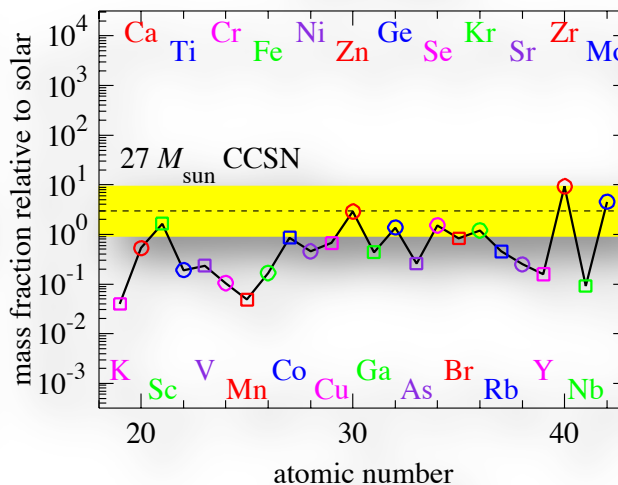
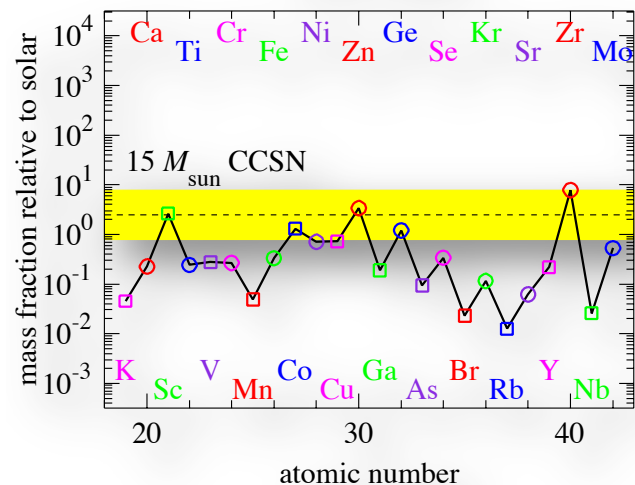
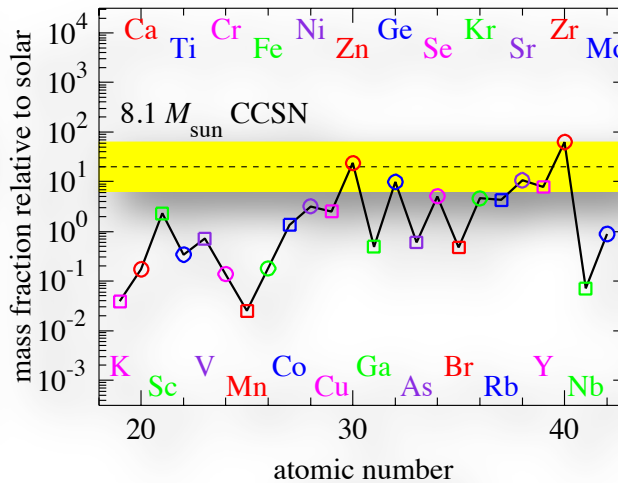
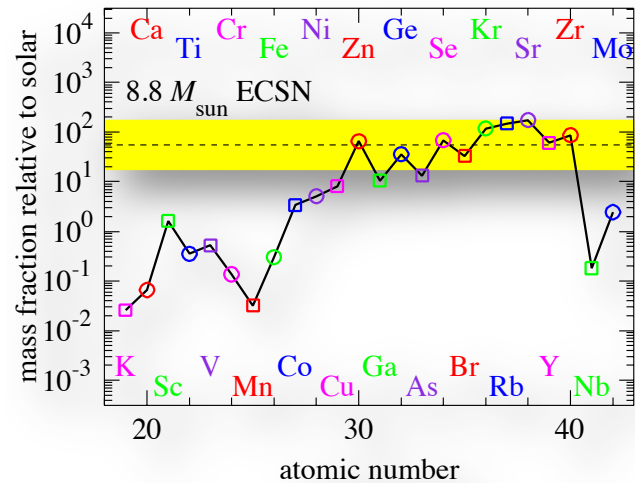
Y_e distribution in the innermost ejecta ($\sim 0.01 M_{\odot}$)

❖ light SNe have more n-rich ejecta (less neutrino-processed)



❖ massive SNe have more p-rich ejecta (more neutrino-processed)

elemental abundances for each SN



nucleosynthesis in the innermost ejecta ($M_{\text{ej}} \sim 0.01 M_{\odot}$)

❖ light SNe make more trans-iron elements (Zn to Zr)

❖ massive SNe make few trans-iron elements (mainly Zn and Zr)

Wanajo+2015, in prep.

a little bit more of ECSNe: ^{48}Ca

www.kadonis.org

^{42}Ti 199.00 ms	^{43}Ti 509.00 ms	^{44}Ti 59.99 a	^{45}Ti 3.08 h	^{46}Ti 8.25	^{47}Ti 7.44	^{48}Ti 73.72	^{49}Ti 5.41	^{50}Ti 5.18
^{41}Sc 596.00 ms	^{42}Sc 681.00 ms	^{43}Sc 3.89 h	^{44}Sc 3.97 h	^{45}Sc 100	^{46}Sc 83.79 d	^{47}Sc 3.35 d	^{48}Sc 1.82 d	^{49}Sc 57.20 m
^{40}Ca 96.94	^{41}Ca 102.01 ka	^{42}Ca 0.647	^{43}Ca 0.135	^{44}Ca 2.09	^{45}Ca 162.62 d	^{46}Ca 0.004	^{47}Ca 4.54 d	^{48}Ca 0.187
^{39}K 93.2581	^{40}K 1.25×10^9 y	^{41}K 6.7302	^{42}K 12.32 h	^{43}K 22.30 h	^{44}K 22.13 m	^{45}K 17.30 m	^{46}K 1.75 m	^{47}K 17.50 s

^{48}Ca

- ❖ doubly magic ($Z = 20$, $N = 28$) n-rich isotope
- ❖ almost stable ($t_{1/2} = 1.9 \times 10^{19}$ yr by double β -decay)
- ❖ origin is unknown; only suggested are hypothetical SNe Ia

a little bit more of ECSNe: ^{60}Fe

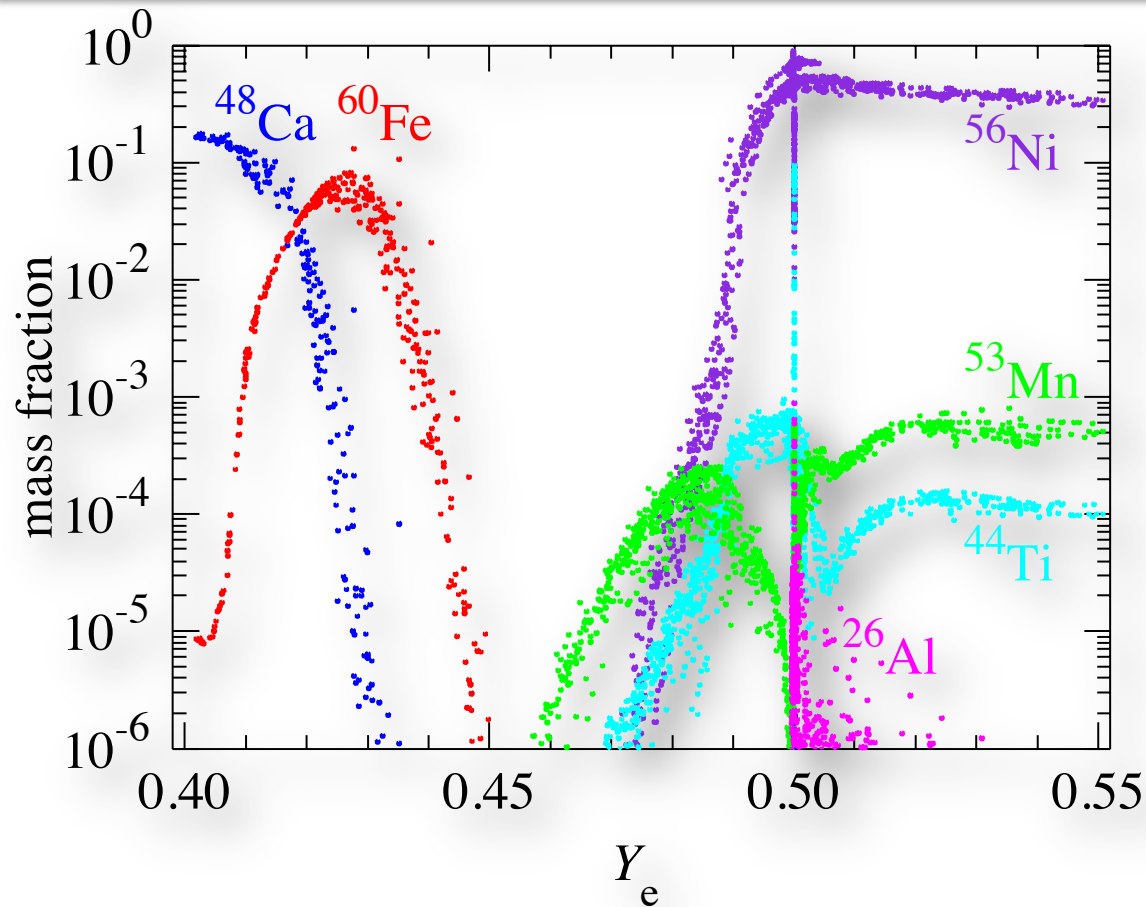
www.kadonis.org

^{56}Ni 6.08 d β^+	^{57}Ni 1.48 d β^+	^{58}Ni 68.077 38.7 mb	^{59}Ni 75.99 ka 87 mb, β^+	^{60}Ni 26.223 30 mb	^{61}Ni 1.14 82 mb	^{62}Ni 3.634 22.3 mb
^{55}Co 17.53 h β^+	^{56}Co 77.23 d β^+	^{57}Co 271.76 d β^+	^{58}Co 70.86 d β^+	^{59}Co 100 38 mb	^{60}Co 5.27 a β^-	^{61}Co 1.65 h β^-
^{54}Fe 5.845 27.6 mb	^{55}Fe 2.74 a 75 mb, β^+	^{56}Fe 91.754 11.7 mb	^{57}Fe 2.119 40 mb	^{58}Fe 0.282 12.1 mb	^{59}Fe 44.50 d β^-	^{60}Fe 1.50 Ma β^-

^{60}Fe

- ❖ n-rich radionuclide with $t_{1/2} = 2.62 \times 10^6$ yr (Rugel+2009)
- ❖ live ^{60}Fe in the early solar system (e.g., Tachibana+2003)
- ❖ live ^{60}Fe in the Milky Way Galaxy (e.g., Harris+2005)
- ❖ CCSNe (n-capture in He, O/Ne layers) can be the sources?

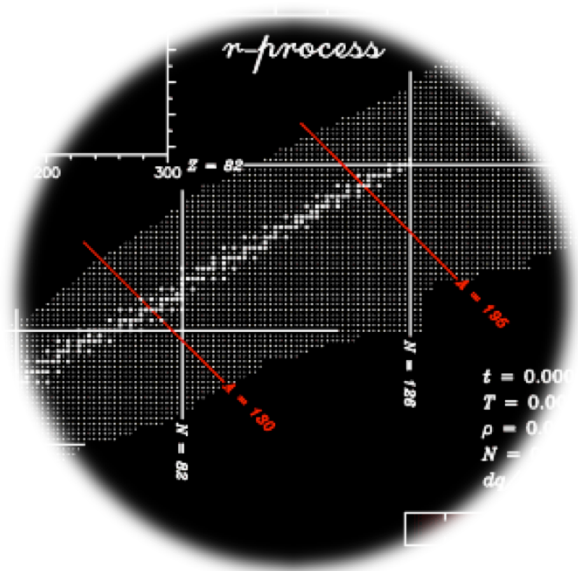
ECSNe make both ^{48}Ca and ^{60}Fe !



both are made in
NSE-like conditions with
 $Y_e = 0.40-0.45$

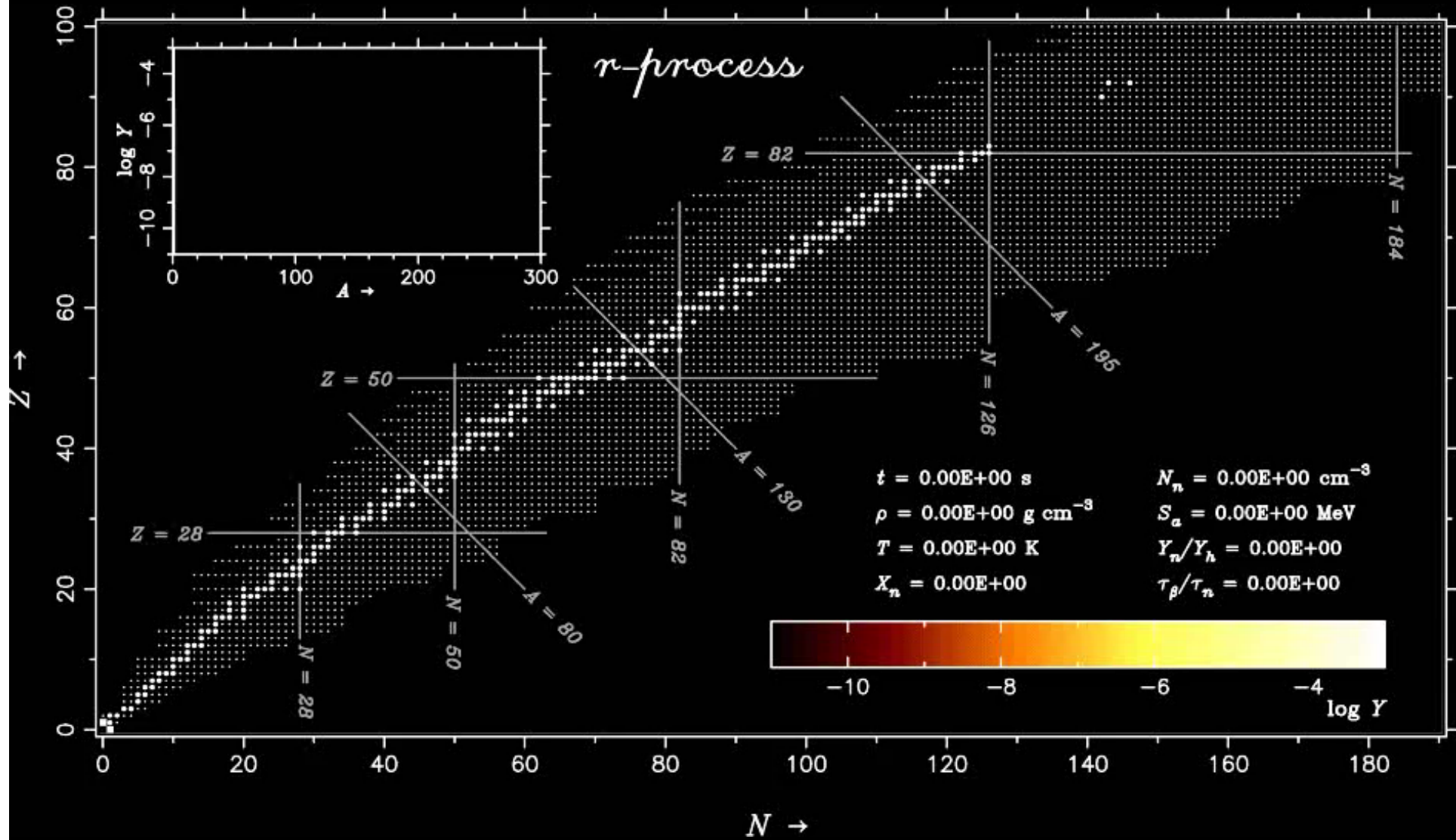
❖ nucleosynthesis is
extremely sensitive to Y_e ;
a few percent level of
accuracy is needed !!!

Wanajo+2013

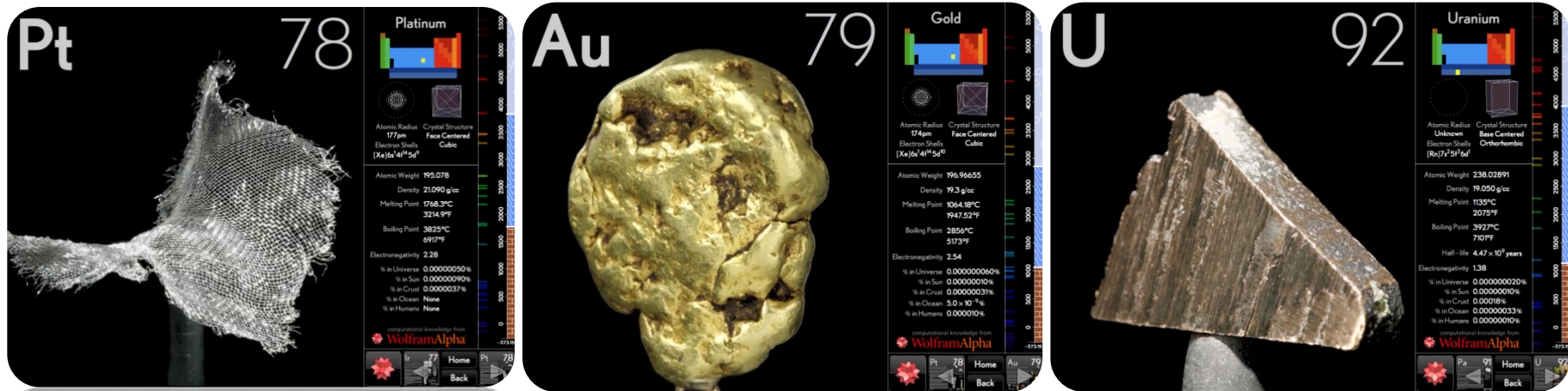


2. nucleosynthesis during the next 10 seconds;
is the answer blowing in the wind?
(Wanajo 2013)

do we find an r-process ?



how to make the 3rd peak and beyond



physical condition for making $A \geq 200$ (Hoffman+1997)

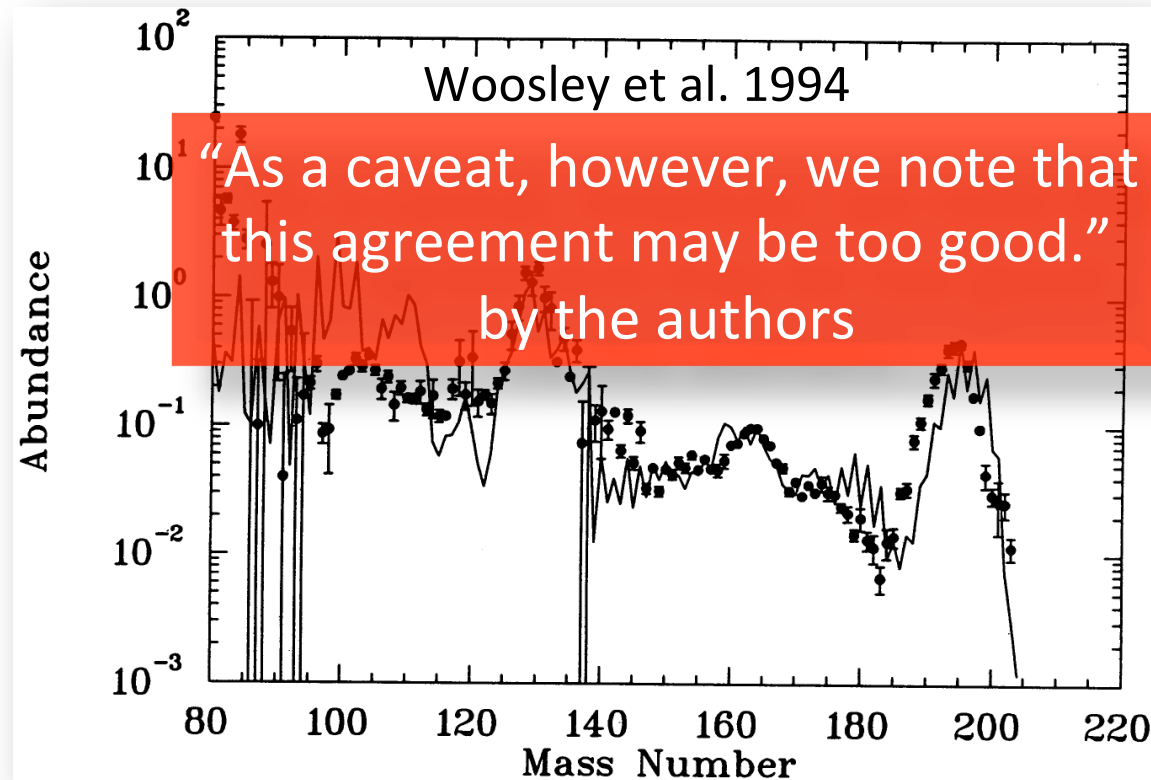
❖ high entropy scenario ($S > 100 k_B \text{ nuc}^{-1}$)

$$f_{200} = (S / 230 k_B \text{ nuc}^{-1})^3 / [(Y_e / 0.4)^3 (\tau / 20 \text{ ms})] \geq 1$$

❖ low entropy scenario ($S < 100 k_B \text{ nuc}^{-1}$)

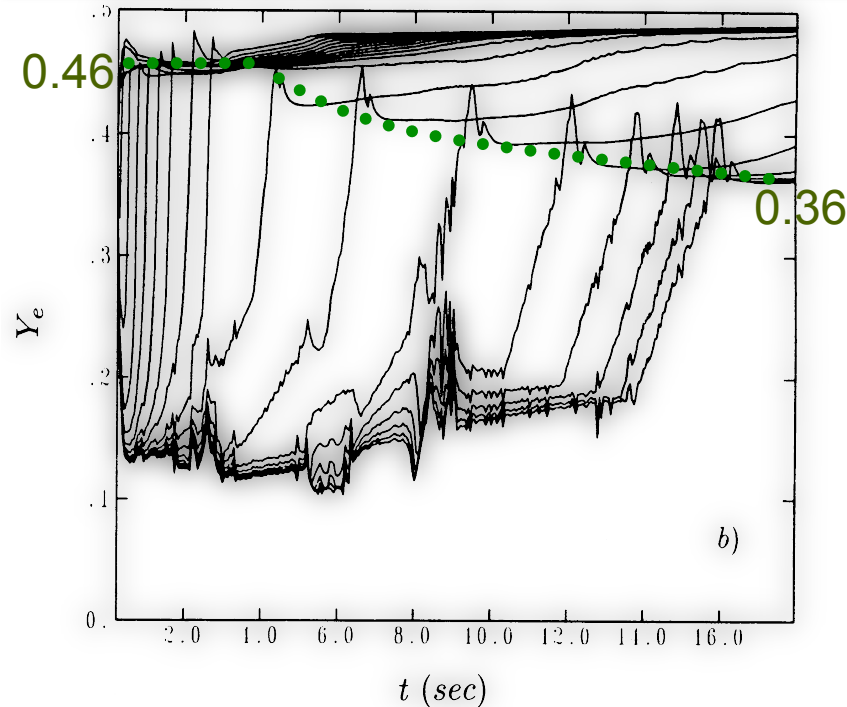
$$Y_e < 0.2 \quad \text{with any } S, \tau$$

high-entropy SN neutrino-driven wind

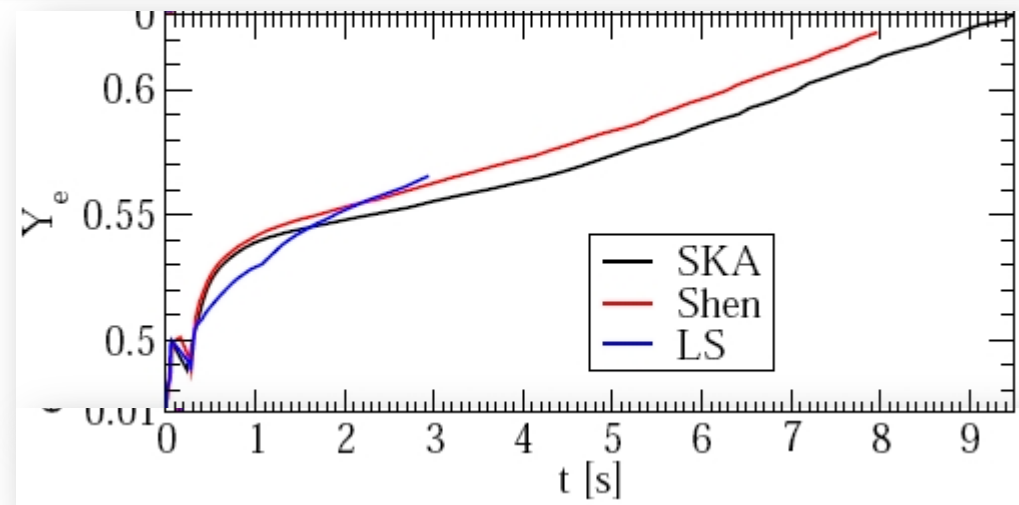


- ❖ successful r-process in the neutrino-driven winds of $S_{\text{rad}} \sim 400 k_B/\text{nuc}$ (1D hydro, $20 M_{\odot}$ star; Meyer+1992; Woosley+1994)
- ❖ but such high entropy is unlikely ($< 200 k_B/\text{nuc}$; Takahashi+1994; Qian+1996; Otsuki+2000)

neutrino-driven wind is “proton-rich” ?



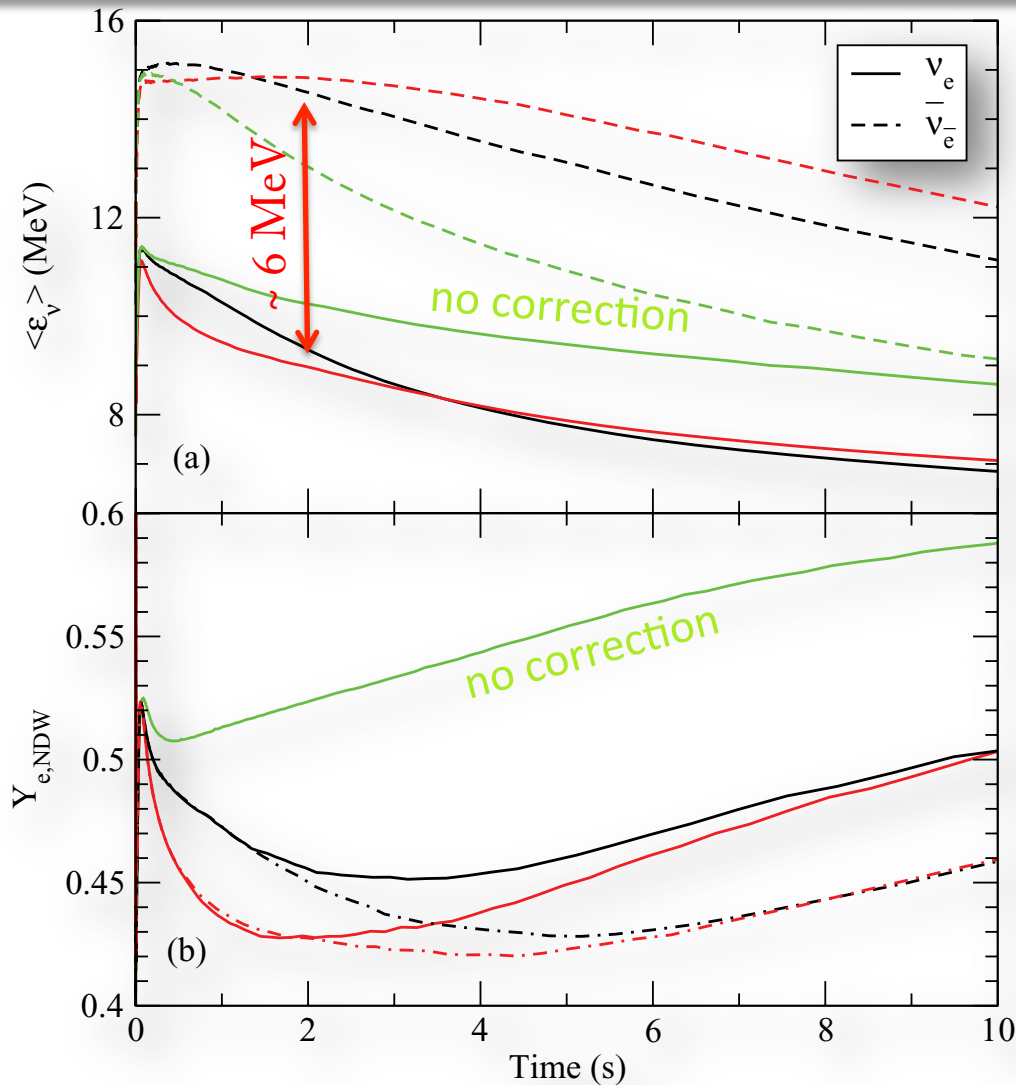
Woosley+1994



self-consistent explosion of a $9 M_{\odot}$ star
Hüdepohl+2009

- ❖ $Y_e > 0.5$ in all recent neutrino-transport simulations because of similar neutrino energies and luminosities for all flavors (i.e., protons are favored due to the p-n-mass difference)
- ❖ no r-process is expected regardless of S or τ ??

no, PNS wind is “slightly” n-rich!?

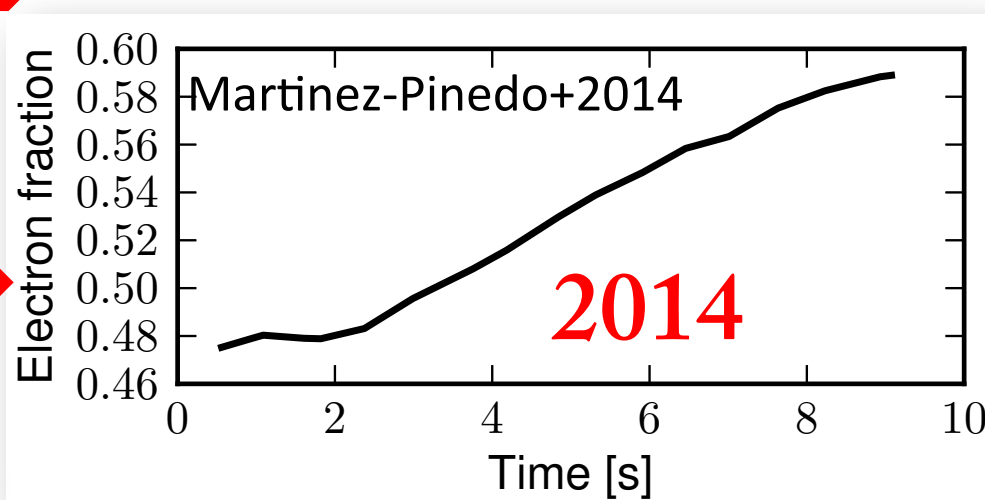
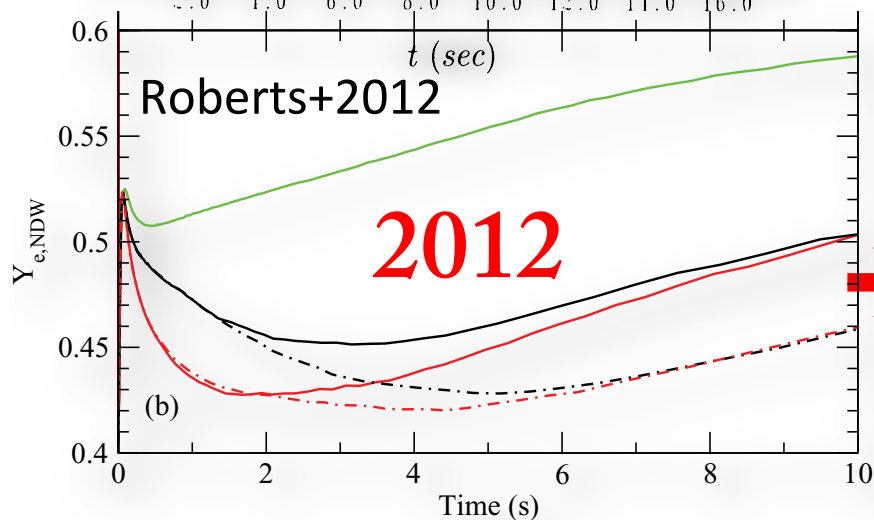
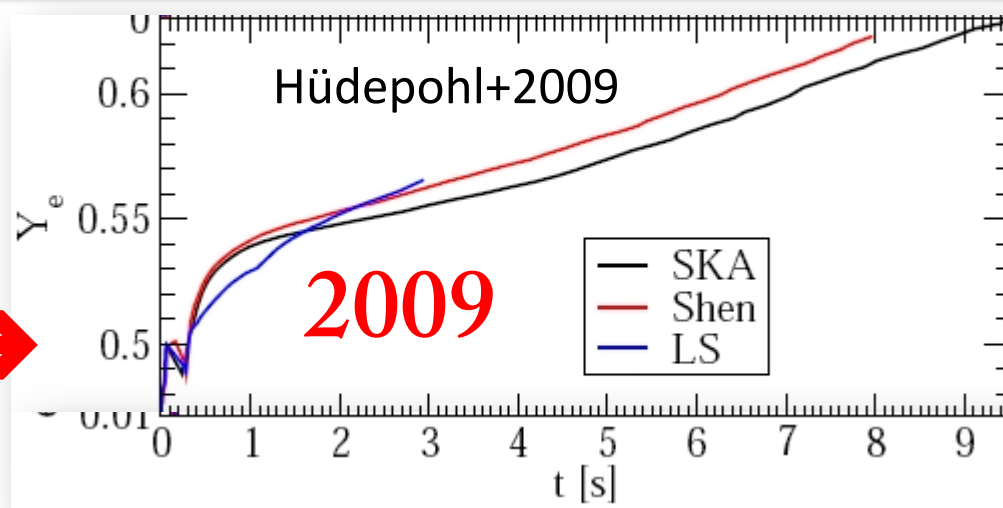
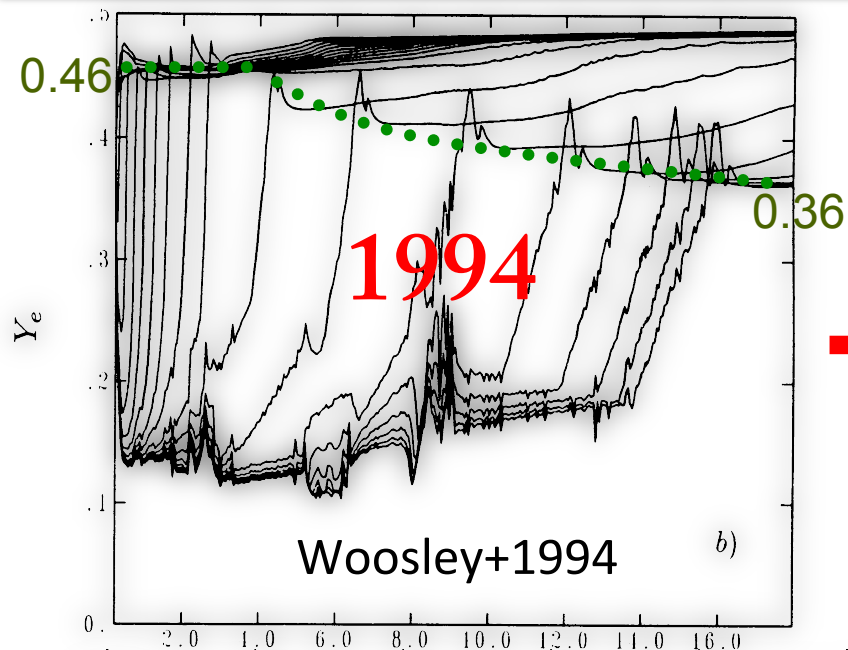


Roberts+2012

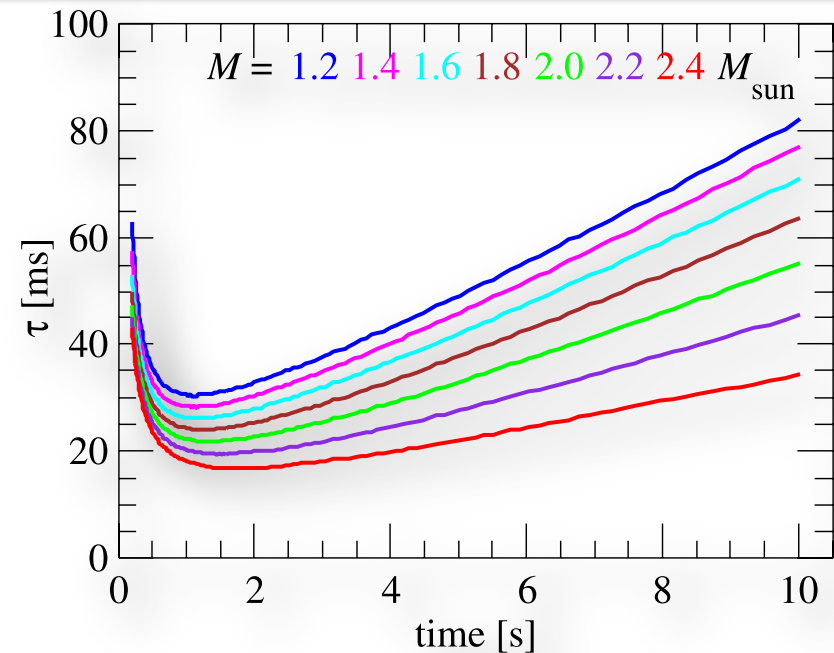
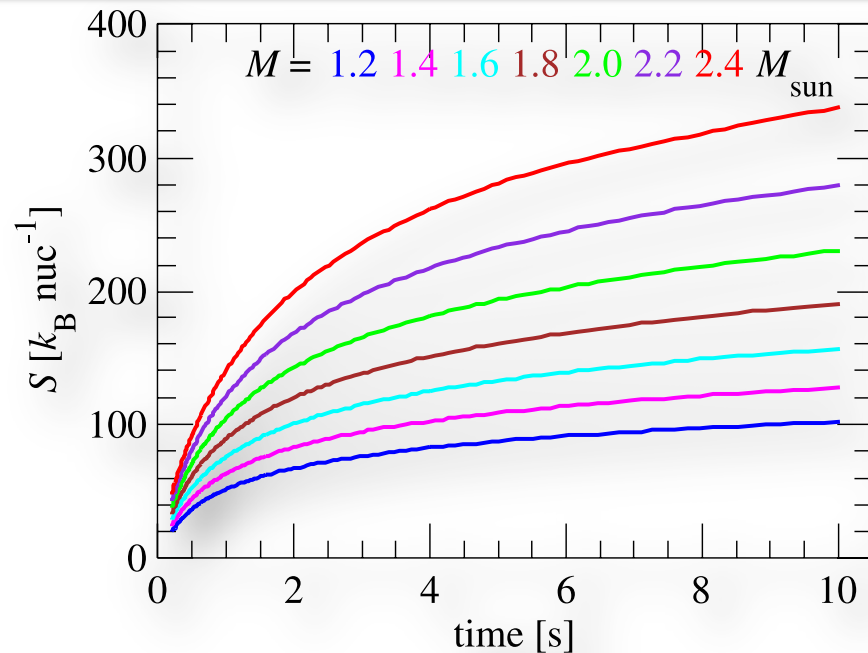
❖ for $\rho > 10^{13}$ g cm $^{-3}$, symmetry energy enhances $\nu_e + n$ and suppress $\bar{\nu}_e + p$ (Reddy+1998; Roberts +2012; M.-Pinedo+2012)

❖ proto-NS wind can be n-rich (down to $Y_e \sim 0.4$?) in the early phase, NOT in the late phase (Pauli-blocking for charged current reactions; Fischer +2012)

“history” of Y_e evolutions: who is right?



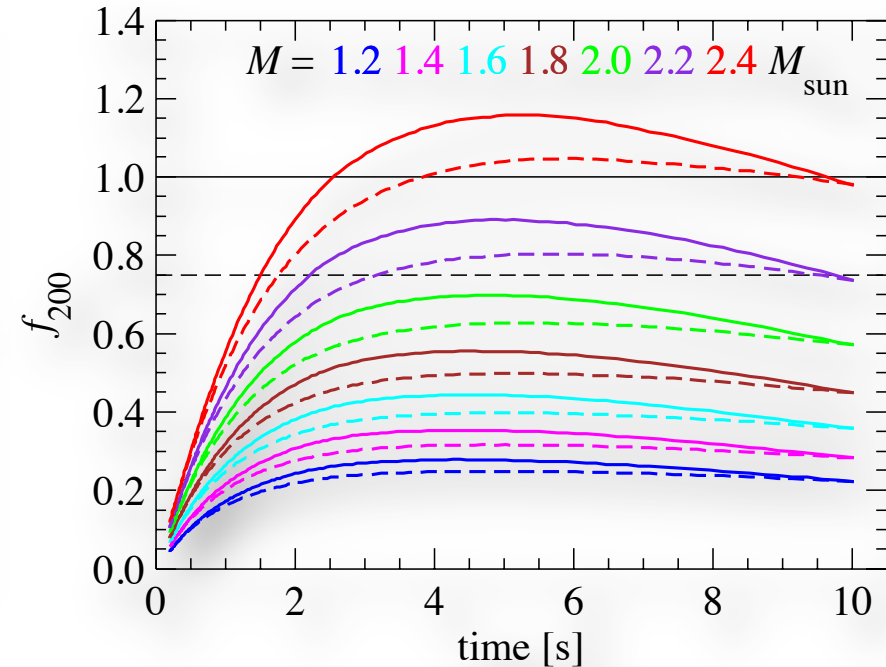
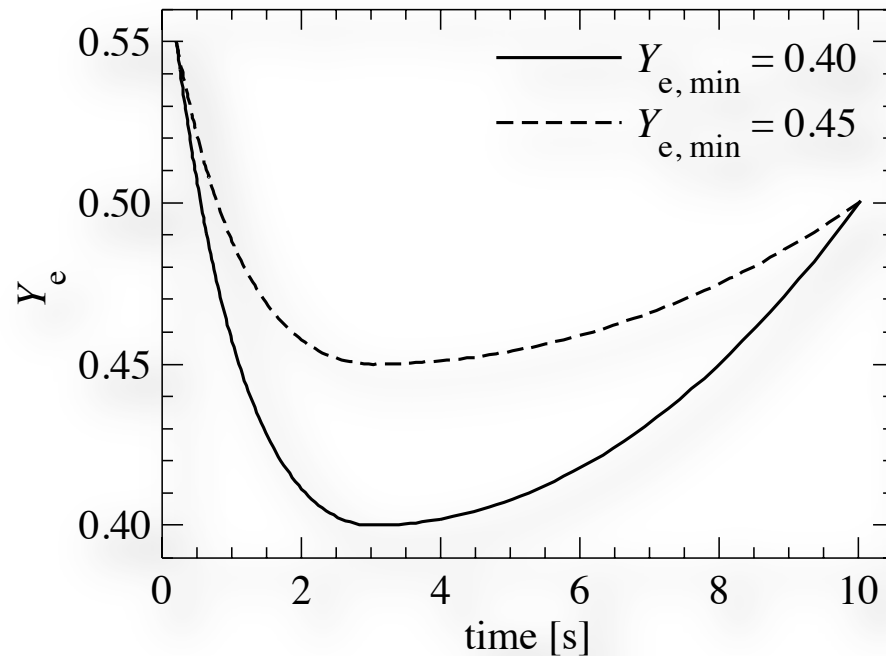
is the answer blowing in the wind?



semi-analytic full-GR wind model (Wanajo 2013)

- ❖ $M_{\text{NS}} = 1.2, 1.4, \dots, 2.4 [M_{\odot}]$
- ❖ $R_{\text{NS}} = 30 (t/t_0)^{-1} + 10$ [km]; $t_0 = 0.2 \leq t \leq 10.2$ [s]
- ❖ $L_{\nu} = 40 (t/t_0)^{-1} [10^{51} \text{ erg s}^{-1}]$ for all flavors
- ❖ $\langle \varepsilon_{\nu} \rangle = 12, 14, 14$ [MeV] for $\nu_e, \bar{\nu}_e$, and the others

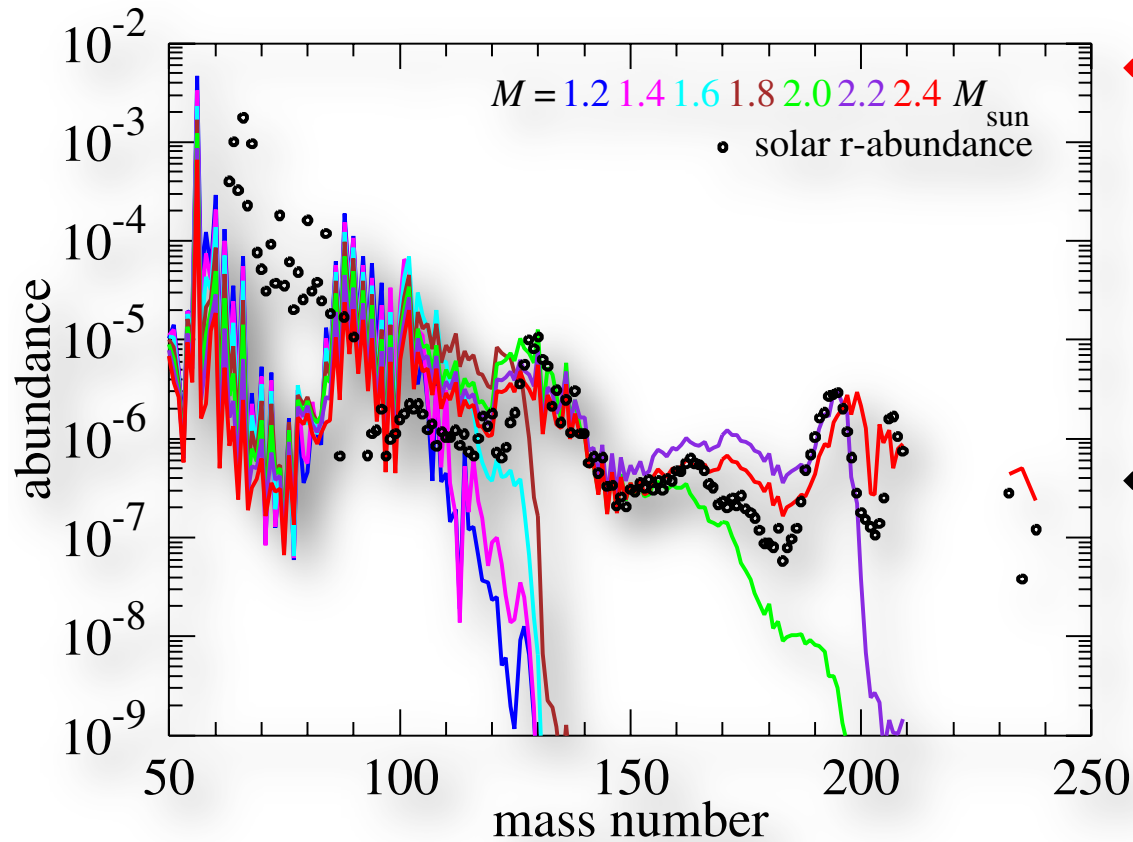
is the answer blowing in the wind?



ad hoc Y_e evolution (to mimic Roberts+2012)

❖ only very massive proto-NSs ($> 2.2 M_{\odot}$) satisfy $f_{200} \geq 1$

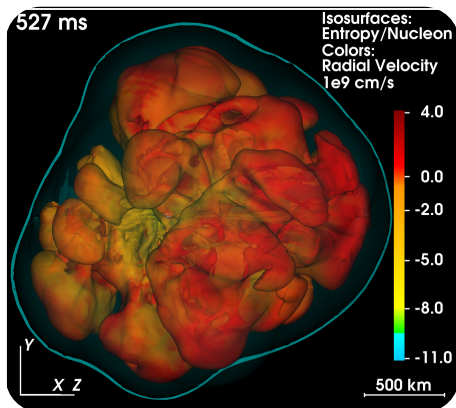
is the answer blowing in the wind?



❖ only (unrealistically) massive proto-NSs ($> 2.2 M_{\odot}$) make the heavy r-elements

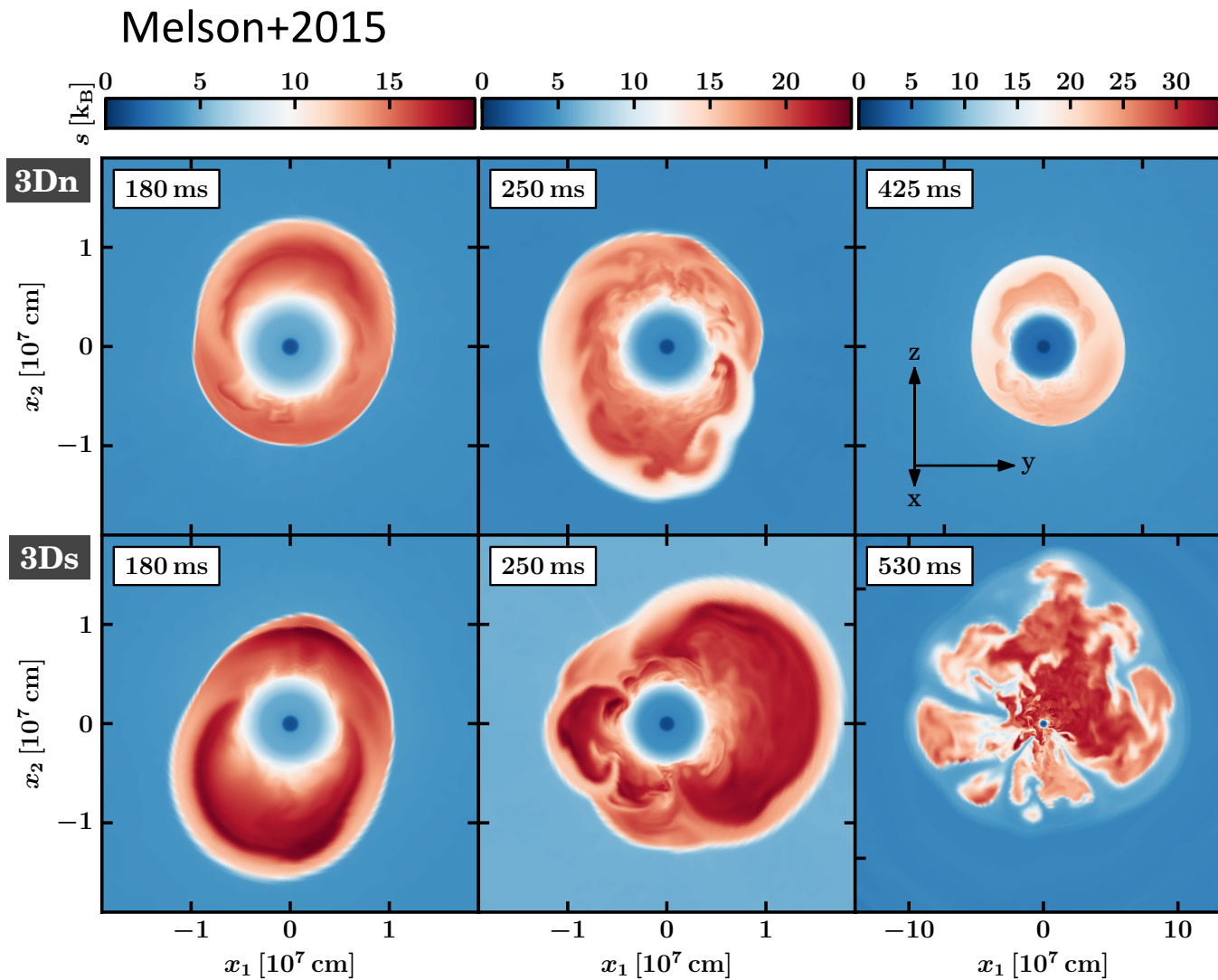
❖ typical proto-NSs ($< 2.0 M_{\odot}$) make weak r-elements ($A \sim 90 - 130$) only

Wanajo 2013



3. how close to successful supernova explosions?
sensitivity to neutrino opacity (Melson+2015)

strange quarks help the explosion?



3D SN simulation
(20 M_{\odot})

❖ without s
($g_a^s = 0$)
explosion failed

❖ with s quarks
($g_a^s = -0.2$)
explosion
succeeded!!!

strange quarks help the explosion?

- ❖ lowest-order differential neutrino-nucleon scattering cross section

$$\frac{d\sigma_0}{d\Omega} = \frac{G_F^2 \epsilon^2}{4\pi^2} \left[c_v^2 (1 + \cos \theta) + c_a^2 (3 - \cos \theta) \right]$$

- ❖ strange quark contribution $g_a^s (\leq 0)$

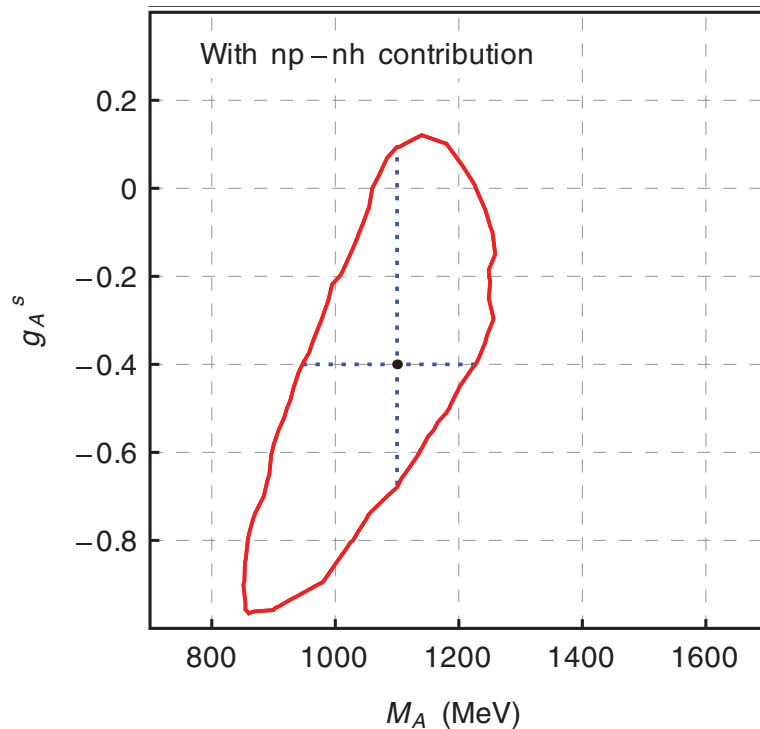
$$c_a = \frac{1}{2} (\pm g_a - g_a^s) \quad \text{where the plus sign is for } \nu p \text{ and the minus sign for } \nu n$$

- ❖ c_a^2 decreases in the (n-rich) neutrinospheric region
 - ~10% reduction of neutral-current opacity ($g_a = 1.26$, $g_a^s = -0.2$)
 - ~10% increase of neutrino luminosities and mean energies
 - more efficient neutrino heating in the gain region

- ❖ only a 10% change of neutrino opacity can help the explosion !!!

strange quarks help the explosion?

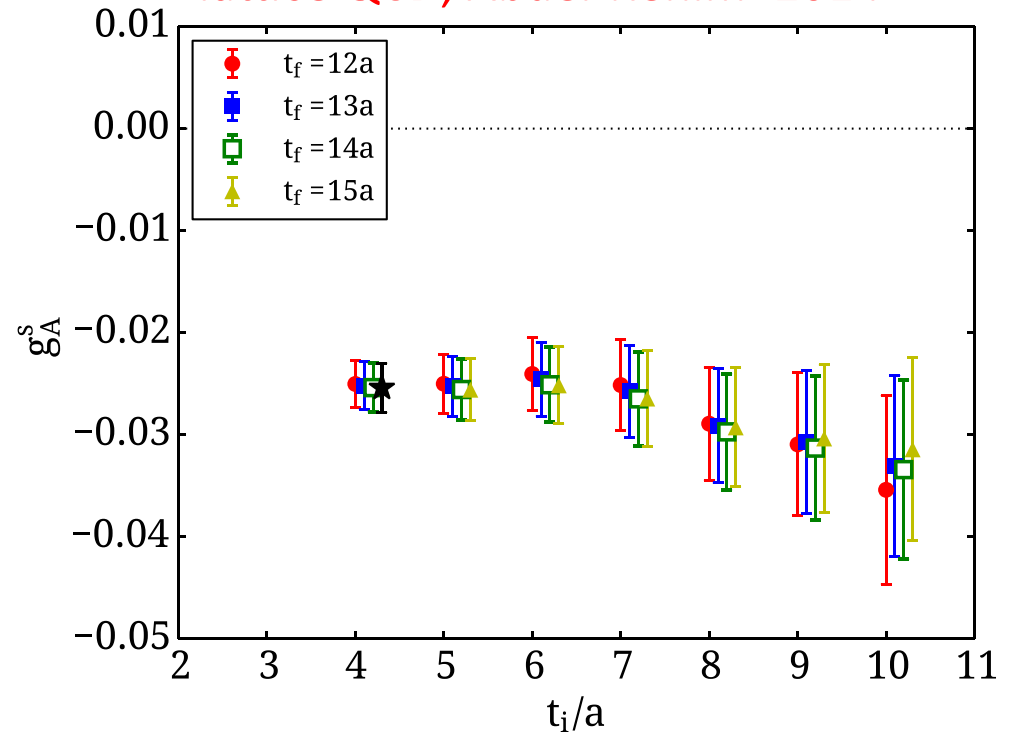
MiniBooNE; Golan+2013



experiments

- ❖ large (negative) value with large errors:
 $g_a^s = -0.8$ to 0

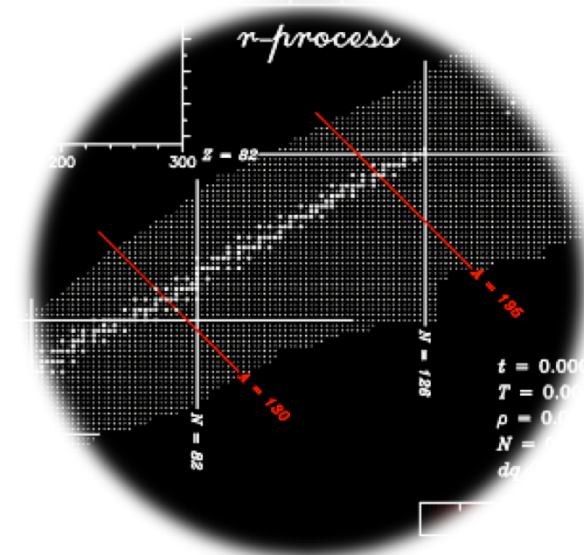
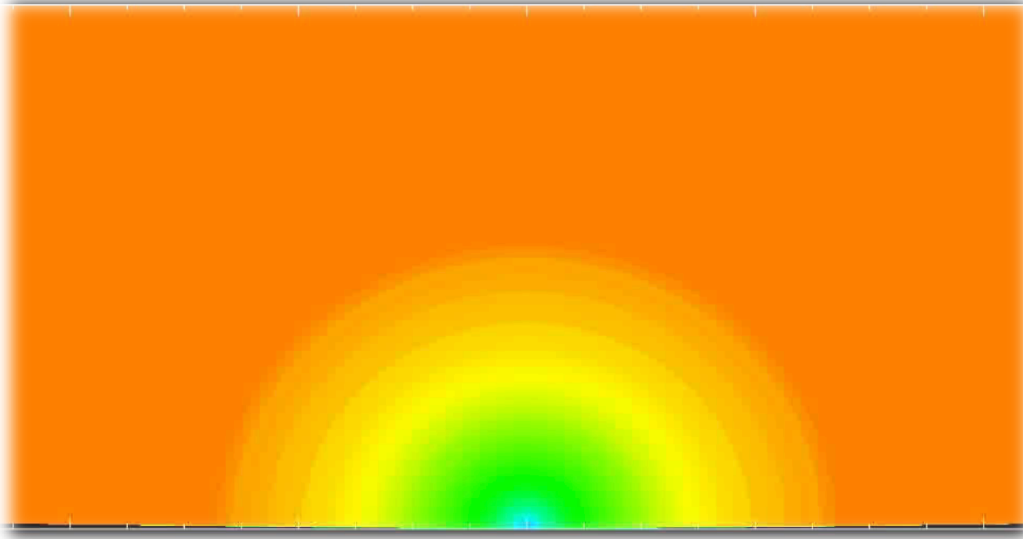
lattice QCD; Abdel-Rehim+2014



theory

- ❖ small (negative) value with small errors:
 $g_a^s = -0.05$ to -0.02

summary



- ❖ core-collapse SNe can make light trans-iron elements (Zn to Zr) but r-process elements (see NS merger works; Wanajo+2014)
- ❖ nucleosynthesis in the innermost ejecta is highly sensitive to neutrino spectra; a few percent accuracy for Y_e is needed
- ❖ 10% change of neutrino opacity can help SN explosions