

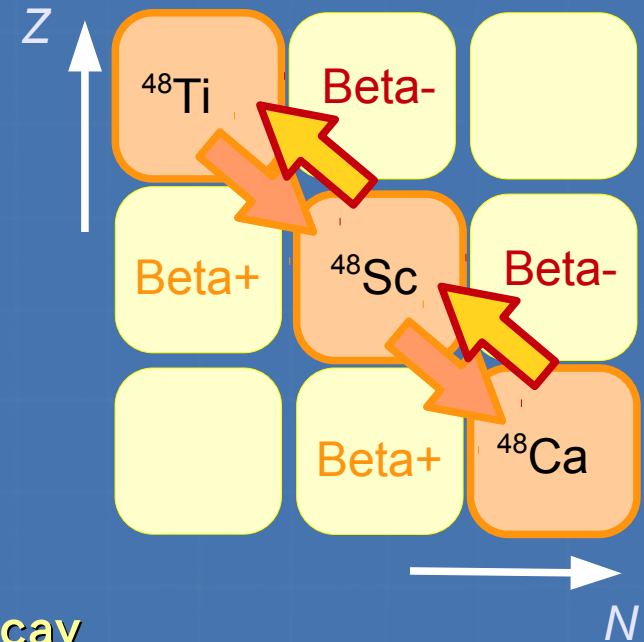
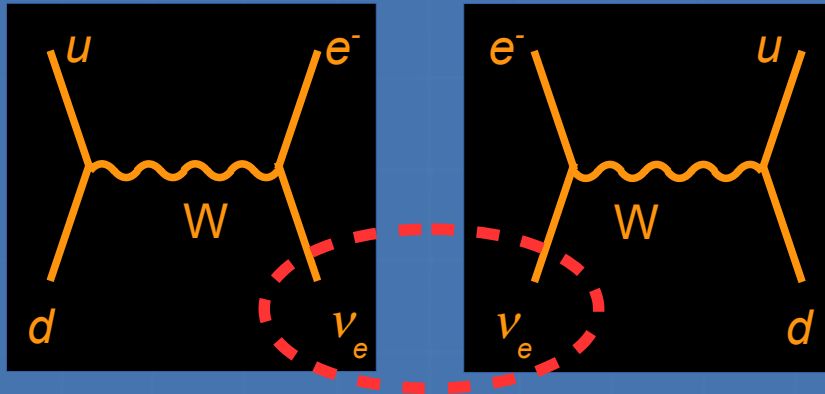


Constraining the neutrino mass with low-energy experiments on nuclei

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Double beta decay of ^{48}Ca



Two modes associated with double beta decay

2ν process (neutrino emission):

Take place when neutrinos is Majorana or Dirac particles:

$$(Z, A) \rightarrow (Z+2, A) + 2e^- + 2\bar{\nu}$$

0ν process (neutrinoless):

Take place when neutrinos is Majorana particles:

$$(Z, A) \rightarrow (Z+2, A) + 2e^-$$

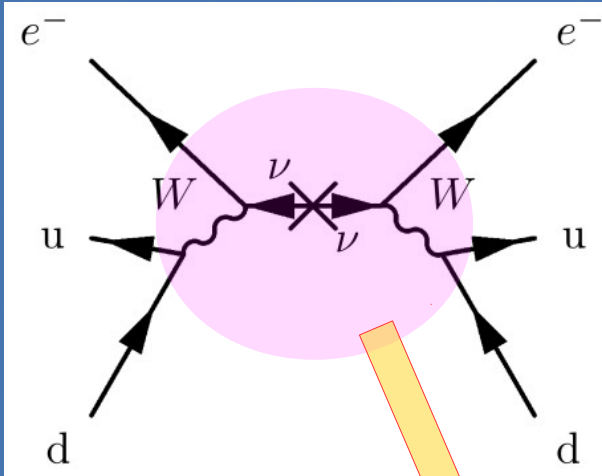
Nuclear matrix element (NME)

Majorana particle or not?

Effective neutrino mass?

Neutrinoless double beta decay

Relation between neutrino mass and decay half life:



$$[T_{1/2}^{0\nu}]^{-1} = G_1^{0\nu} |M^{0\nu}|^2 \left(\frac{\langle m_{\beta\beta} \rangle}{m_e} \right)^2$$

Half life

NME

Effective neutrino mass



lepton number violation
(beyond the standard model)


Nuclear Matrix element under the closure approx.

$$M^{0\nu} = \left\langle f \left| \sum_{a,b} \tau_a^+ \tau_b^+ \left\{ \underbrace{-(g_V/g_A)^2 H_F(r)}_{\text{Fermi}} + \underbrace{\sigma_a \cdot \sigma_b H_{GT}(r)}_{\text{Gamow-Teller}} - \underbrace{(3(\sigma_a \cdot r)(\sigma_a \cdot r) - \sigma_a \cdot \sigma_b) H_T(r)}_{\text{Tensor}} \right\} \right| i \right\rangle$$

$= H\alpha$ ($\alpha: F, GT, T$)

NME component: Neutrino potential

Neutrino potential under closure approx. is calculated within the precision of "0.0010 [MeV fm].

Precise calculation provided by "MAXIMA" 

$$H_{\alpha}(r) = \frac{2R}{\pi} \int_0^{\infty} f_{\alpha}(qr) \frac{h_{\alpha}(q)}{q + \langle E \rangle} q dq$$

$\langle E \rangle$ = averaged energy of virtual intermediate state

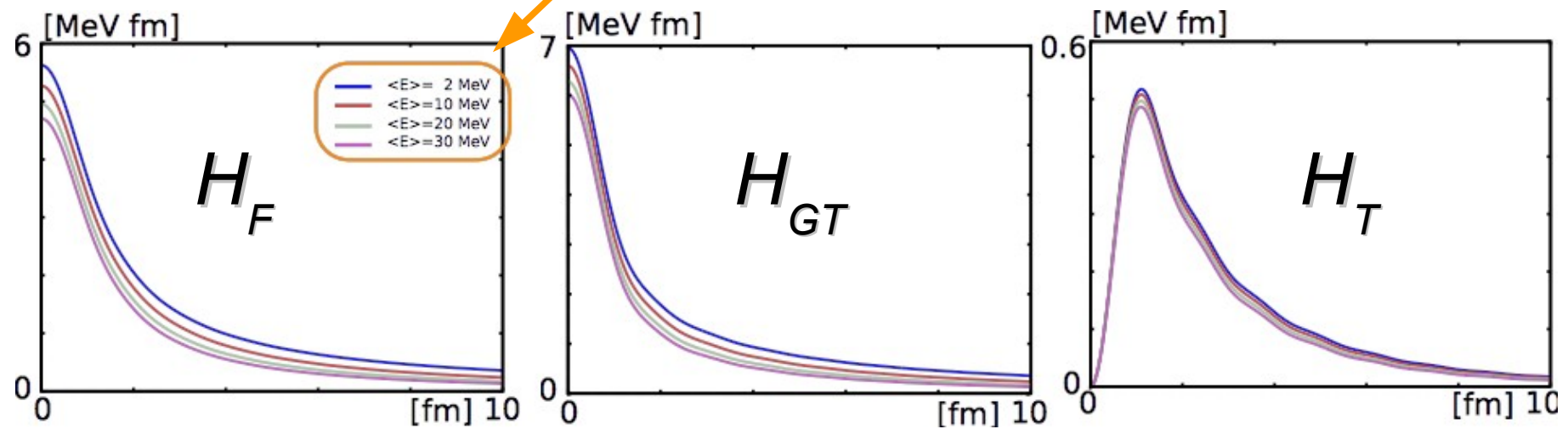
q: momentum of virtual neutrino
f $_{\alpha}$: spherical Bessel function ($\alpha=0,2$)

$\alpha = F, GT, T$ (Fermi, Gamow-Teller, Tensor parts)

$\langle E \rangle$ dependence

$\langle E \rangle = 2 \sim 30$ MeV

$\langle E \rangle$: closure parameter = average energy



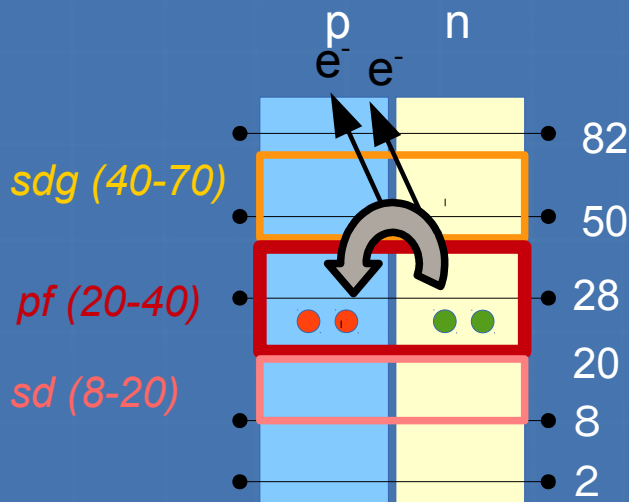
In the following $\langle E \rangle = 0.50$ MeV, suggested by calc. w/o the closure approximation. Senkov-Horoi-Brown PRC (2013)

NME component: Initial/final state

Y.I.-Shimizu-Otsuka-Utsuno-Menendez-Honma-Abe, Phys. Rev. Lett. 116 (2016) 112502
Large-scale shell model analysis on nuclear matrix element

$$\left[T_{1/2}^{0\nu} \left(0_i^+ \rightarrow 0_f^+ \right) \right]^{-1} = G^{0\nu} |M^{0\nu}|^2 \left(\frac{\langle m_\nu \rangle}{m_e} \right)^2$$

$G^{0\nu}$: phase space factor
 m_e : electron mass
These values are well known.



Inclusion rate of 2nd major shell components :

^{48}Ca (22%) , ^{48}Ti (33%)

$sd + pf$

^{48}Ca (~2%) , ^{48}Ti (~2%)

$pf + sdg$

This result shows that
It should be necessary to take into account *sd* shell

In case of *sd + pf* ...

$$^{48}\text{Ca} (p,n) = (20, 28) = (20+0, 20+8) = (8+12, 8+20)$$

$$^{48}\text{Ti} (p,n) = (22, 26) = (20+2, 20+6) = (8+14, 8+18)$$

1shell

+1 shell

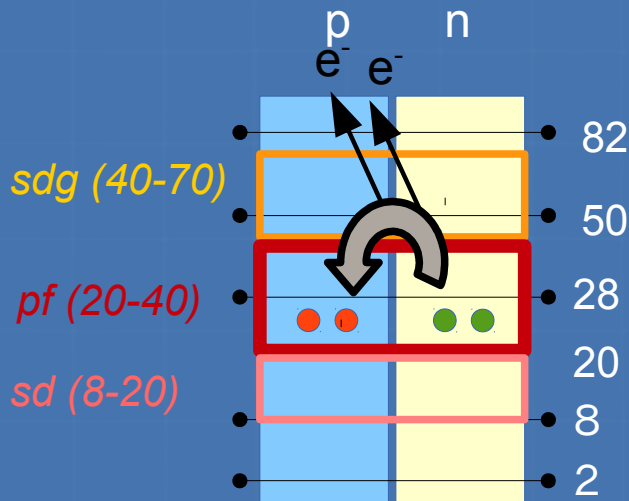
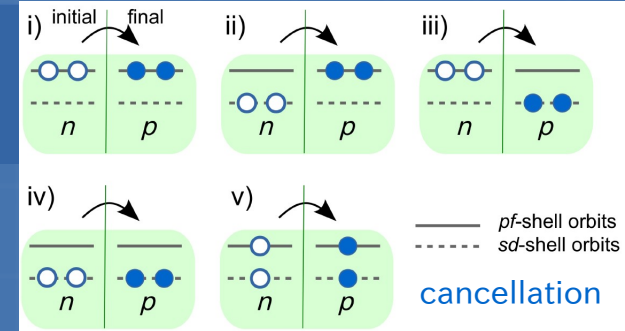
10^6 dim

~ 10^9 dim

Neutrinos of ordinary type

Y.I.-Shimizu-Otsuka-Utsuno-Menendez-Honma-Abe, Phys. Rev. Lett. 116 (2016) 112502
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Inclusion rate of 2nd major shell components :

^{48}Ca (22%) , ^{48}Ti (33%)

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This result shows that
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$M^{0\nu}$ (1 shell) 0.833

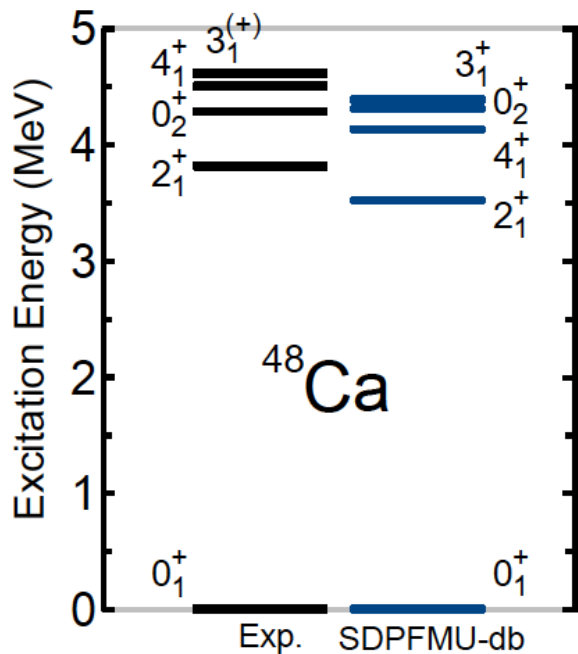
$M^{0\nu}$ (2 shells) 1.118

34.2 %
increased

Due to $(1/1.34)^2 \sim 0.56$,
it means that
the half-life is almost halved
for the same neutrino mass.

Energy spectra, as a test of the nuclear structure calculation

As an evidence of good description, the energy spectra made by SDPFMU-db is compared to the experiment; SDPFMU-db is an effective interaction made for 2 major shell description.



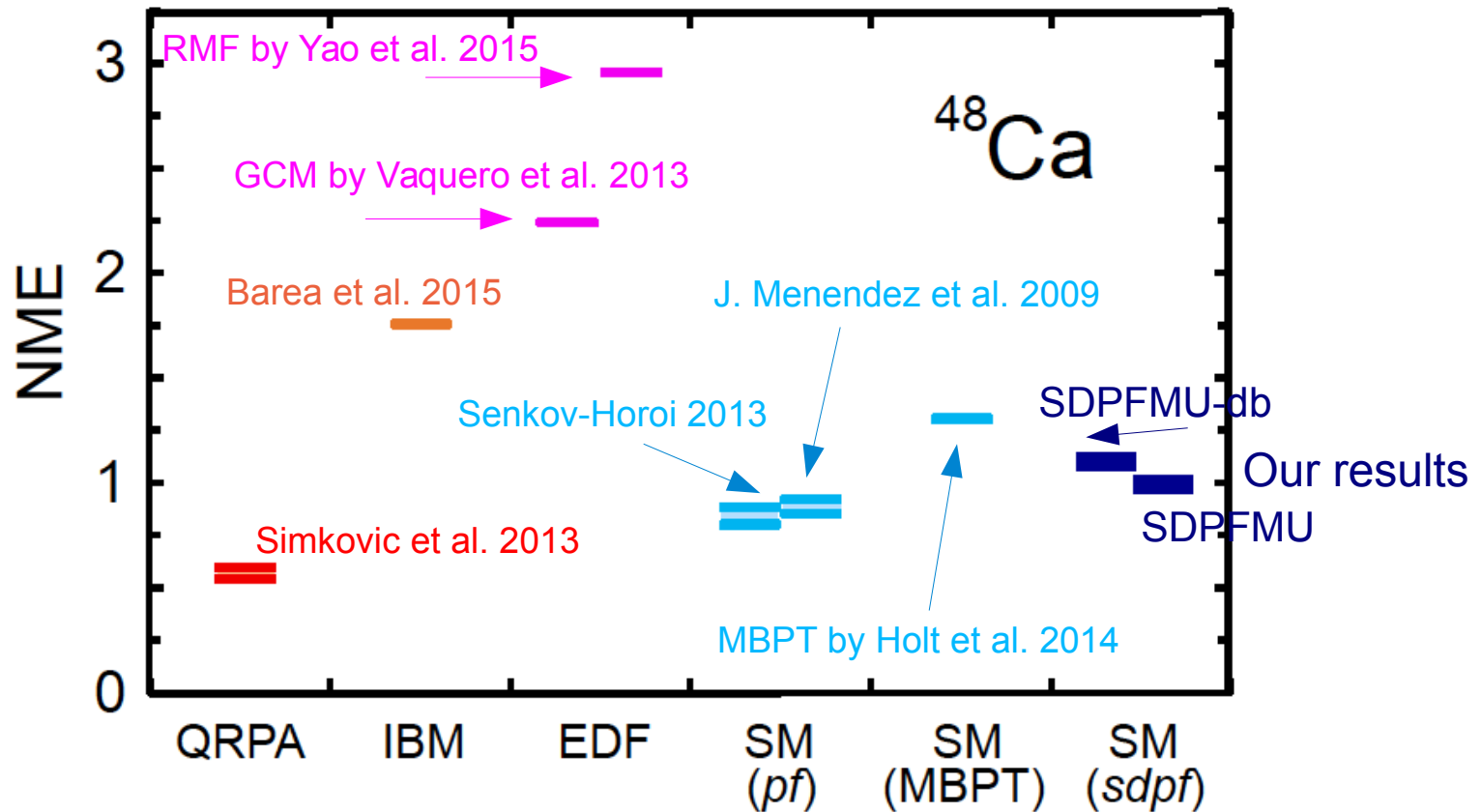
2hw component ratio

SDPFMU-db

Ca48 (g.s.) , Ti48 (g.s.): 22%, 33%

Summary: NME for $0\nu\beta\beta$ of ^{48}Ca

Comparison of neutrinoless double beta decay NME (with ranges)



Constraint on the neutrino mass

$$\langle m_\nu \rangle = \sqrt{\frac{m_e^2}{|M^{0\nu}|^2 G^{0\nu}} \left[T_{1/2}^{0\nu} \left(0_i^+ \rightarrow 0_f^+ \right) \right]^{-1}}$$

Constants:

$$G^{0\nu} = 1.27^4 \times 0.2989 \times 10^{-15} \text{ y}^{-1} \text{ (for Ca48, Kotila-Iachello, PRC 2012)}$$

$$m_e = 5.110 \times 10^5 \text{ eV}$$

$$T^{0\nu} > 5.8 \times 10^{22} \text{ y (ELEGANT IV, 2008)}$$

Increase of the nuclear matrix element (NME) makes the experiment **sensitive to** smaller neutrino masses (for the same half-life).

Constraint on the neutrino mass

$$\langle m_\nu \rangle = \sqrt{\frac{m_e^2}{|M^{0\nu}|^2 G^{0\nu}} \left[T_{1/2}^{0\nu} \left(0_i^+ \rightarrow 0_f^+ \right) \right]^{-1}}$$

Constants:

$G^{0\nu} = 1.27^4 \times 0.2989 \times 10^{-15} \text{ y}^{-1}$ (for Ca48, Kotila-Iachello, PRC 2012)

$m_e = 5.110 \times 10^5 \text{ eV}$

$T^{0\nu} > 5.8 \times 10^{22} \text{ y}$ (ELEGANT IV, 2008)

10.1 eV
(old)

Upper limit for the effective mass:

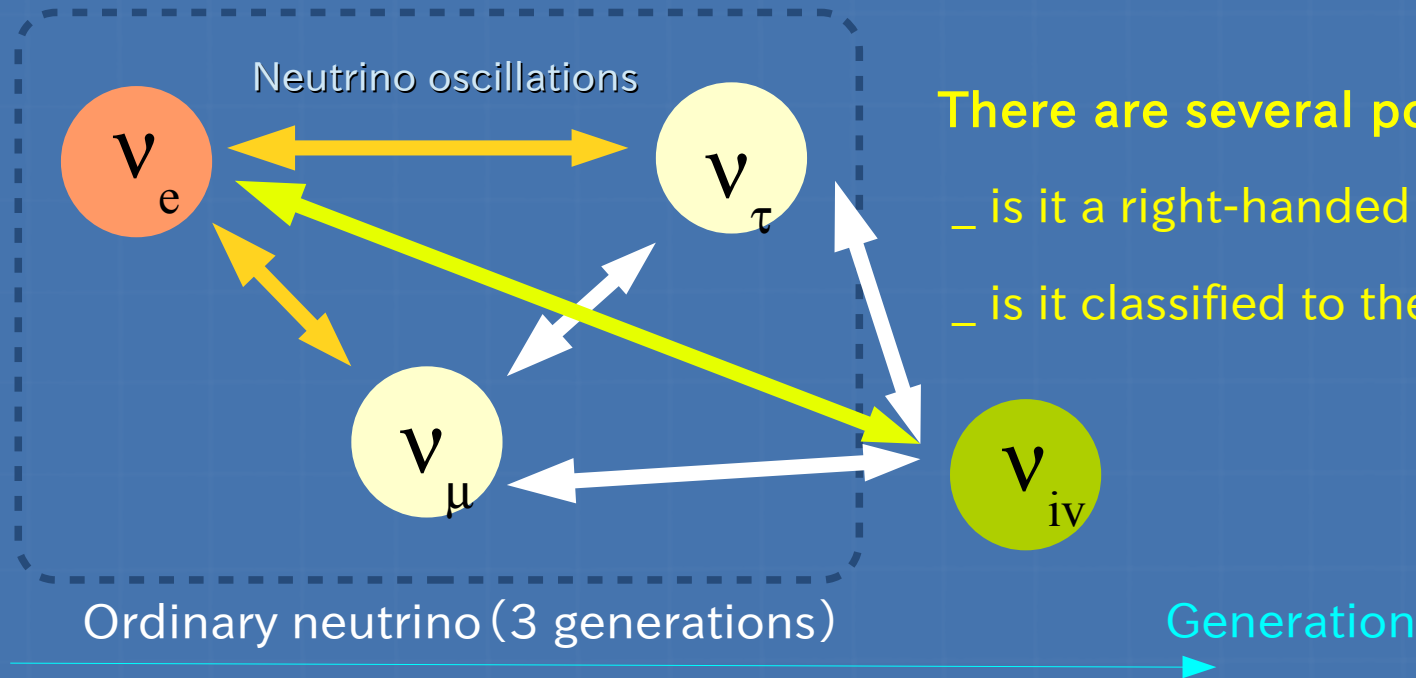
$$\langle m_\nu \rangle < 7.5 \text{ eV}$$

(new)

The latest value
using ^{48}Ca

Sterile neutrino

Mass range (at present):
 10^{-10} to 10^{20} GeV/c²



There are several possibilities:

- _ is it a right-handed neutrino ?
- _ is it classified to the 4th generation ?

Fact implying the existence of sterile neutrino :

- _ LSND experiment (and also MiniBooNe experiment)
- _ WMAP experiment ... number of neutrino generation as 4.3.

Sterile neutrino

Under assuming the following form (assuming nonzero sterile neutrino mass), the NME value is calculated.

$$\left[T_{1/2}^{0\nu} \left(0_i^+ \rightarrow 0_f^+ \right) \right]^{-1} = G^{0\nu} \left\{ |M^{0\nu}|^2 \left(\frac{\langle m_\nu \rangle}{m_e} \right)^2 + |M^{0N}|^2 \langle \eta_N \rangle^2 \right\}$$

ordinary neutrino

Sterile neutrino

Refs. Vergados-Ejiri-Simikovic Rep Prog Phys (2012), Horoi PRC (2013)

View
Points

- 1) The value of NME depends on the sterile neutrino mass.
- 2) Sterile neutrino mass appears in the representation.



These two effect can
enhance or cancel with each other.

Result: sterile neutrino

Neutrino potential (in general form)

$$H_\alpha(r) = \frac{2R}{\pi} \int_0^\infty \frac{f_\alpha(qr) h_\alpha(q) q^2}{\sqrt{q^2 + m_\nu^2} (\sqrt{q^2 + m_\nu^2} + \langle E \rangle)} dq$$

Massless limit

$$\frac{2R}{\pi} \int_0^\infty f_\alpha(qr) \frac{h_\alpha(q)}{q + \langle E \rangle} q dq$$

Heavy mass limit


$$\frac{1}{m_\nu^2} \frac{2R}{\pi} \int_0^\infty f_\alpha(qr) h_\alpha(q) q^2 dq$$

M_{II}

Ordinary neutrino (massless limit)

$M^{0\nu}$ (1 major)	0.833	 34.2 % increased
$M^{0\nu}$ (2 major)	1.118	

Sterile neutrino (heavy mass limit)

M^{0N} (1 major)	81.58	 48.0 % increased
M^{0N} (2 major)	120.7	

Consider “heavy mass limit” as an initialization of our research

Simkovic unit

$$M_{II} / (m_e m_p) :$$

$$m_e = 0.510999 \text{ MeV}$$

$$m_p = 938.2723 \text{ MeV}$$

Impact of sterile neutrino

$$\left[T_{1/2}^{0\nu} \left(0_i^+ \rightarrow 0_f^+ \right) \right]^{-1} = G^{0\nu} \left\{ |M^{0\nu}|^2 \left(\frac{\langle m_\nu \rangle}{m_e} \right)^2 + |M^{0N}|^2 \langle \eta_N \rangle^2 \right\}$$

η_N : effective mass of sterile neutrino
(relative to electron mass)

[NME²]

[NME²]

80 %
increased

119 %
increased

Square of mass ratio: $\langle \eta_N \rangle / \langle m_\nu \rangle$ is decisive (to be studied).

Provided sterile neutrino exists,

The balance between 1st term and 2nd term of r.h.s. is dependent on the masses.

By means of precise NME values,
the possibility of the existence of heavy sterile neutrino is suggested
to be determined by the half life measurement.

Summary and Perspective

Summary:

- ◆ We have carried out large-scale shell model calculation (up to $2 \cdot 10^9$ dim diagonalization).

→ More details are explained in

Y.I. *et al.*, Phys. Rev. Lett. 116 (2016) 112502

in terms of

“how this calculation reproduce 2 neutrino processes ”

- ◆ [Effect of the 2nd major shell]

→ More details are explained in

Y.I. *et al.*, Phys. Rev. Lett. 116 (2016) 112502

in terms of

“what causes the increase of NME”, and “existence of cancellation”.

...[Ordinary neutrino] 30% ↑, [Sterile neutrino] 50% ↑

Perspective:

- ◆ Constant on sterile neutrino mass (in preparation)
- ◆ The mechanism of double beta decay w.r.t. mass-dependence.