Glueball dark matter in SU(N) lattice gauge theory



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Based on

- N. Yamanaka et al., arXiv:1910.01440 [hep-ph]
- N. Yamanaka et al., arXiv:1910.07756 [hep-lat]
- N. Yamanaka et al., arXiv:1911.03048 [hep-lat]



Many evidences of Dark matter



Galactic rotation curve



N-body simulation : large-scale structure



DM density extracted from CMB



Bullet cluster : collision of galaxies

$$\mathcal{L}_{\rm YM} = -\frac{1}{4} F_a^{\mu\nu} F_{\mu\nu,a}$$
(a =1,...,Nc²⁻1)

 \Rightarrow The simplest interacting theory

Important properties:

X_M does not have apparent scale, but scale is dynamically generated (dimensional transmutation)

Renormalizable theory, running coupling has logarithmic scale variation,

difference of N_c can generate Λ_{YM} 's which differ by orders of magnitude

No scalars and massive fermions \Rightarrow Free from quadratic divergences

 \Rightarrow No important fine-tuning problem in the choice of Λ_{YM} !

(Suppose a GUT which generates SM and DM, the difference of mass scales between SM and DM is not serious)

 \Rightarrow Theory with very high naturalness

Dark matter in hidden YM theory:

Lightest particles are glueballs ! \Rightarrow SU(N) glueballs are candidate of DM

(summarized in the report of USQCD Collaboration : arXiv:1904.09964 [hep-lat])

Self-interacting dark matter

The DM distribution can be predicted in N-body simulation with gravity only

 \Rightarrow Successful in describing the large scale structure (scale > Mpc)

Introducing DM self-interaction changes the structure smaller than Mpc (= DM-DM scattering)

There are (were?) several problems in the galactic DM distribution:

Core vs Cusp problem:

N-body simulation predicts cuspy DM distribution near the galactic center, whereas observations suggest flat ones.

<u>Too-big-to-fail problem:</u>

Satellite galaxies are less dense than those predicted by the N-body simulation.

Missing satellite problem:

More satellite galaxies than those predicted by the N-body simulation are observed.



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Resolved thanks to improvement of observation?

dy simulation are

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DM density

radius

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Still under debate, but this shows the importance of the investigation of DM-DM scattering

Glueballs of SU(N) Yang-Mills theory are good candidates of dark matter

In this work, we study the interglueball scattering on lattice which is the only way to quantify nonperturbative physics of nonabelian gauge theory.

The Yang-Mills theory depends only on the scale parameter Λ (given N_c): can we determine Λ from observation?

Object:

In this work, we study the interglueball scattering of SU(2) Yang-Mills theory on lattice, and set constraint on its scale parameter Λ .

Nambu-Bethe-Salpeter amplitude

The information of the scattering is included in the following n-point correlator (Nambu-Bethe-Salpeter amplitude):



2-glueball (0++) state mixes with all other multi-glueball states:

 \Rightarrow The source may be chosen as 1-body, 2-body, etc, on convenience

The NBS amplitude obeys the Schroedinger equation below inelastic threshold

Extract the interglueball potential from the NBS amplitude by inversely solving Schroedinger equation

$$\frac{1}{4m_{\phi}}\frac{\partial^2}{\partial t^2} - \frac{\partial}{\partial t} + \frac{1}{m_{\phi}}\nabla^2 + \frac{(\mathbf{r}\times\nabla)^2}{2m_{\phi}r^2} R(t,\mathbf{r}) = \int d^3\mathbf{r}' U(\mathbf{r},\mathbf{r}')R(t,\mathbf{r}')$$
$$R(t,\mathbf{r}) \equiv \frac{C_{\phi\phi}(t,\mathbf{r})}{e^{-2m_{\phi}t}}$$

N. Ishii et al., PLB 712 (2012) 437.

Crucial advantage : do not need ground state saturation

Almost mandatory to use time-dependent HAL method for the glueball analysis, since the glueball correlator becomes very noisy before ground state saturation

Inelastic threshold for glueball = $3m_{\phi}$: high enough to use low t

Subtract centrifugal force for removing higher angular momenta

<u>Result</u>



We test two fitting forms: • Yukawa fit: $V_Y(r) = V_1 \frac{e^{-m_{\phi}r}}{4\pi r}$ $V_1 = -231 \pm 8$ x² d.o.f. = 1.3 • 2-Gaussian fit: $V(r) = V_1 e^{-\frac{(m_{\phi}r)^2}{8}} + V_2 e^{-\frac{(m_{\phi}r)^2}{2}}$

$$V_1 = (-8.5 \pm 0.5)\Lambda$$

 $V_2 = (-26.6 \pm 2.6)\Lambda$ x² d.o.f. = 0.9

DM cross section is derived from phase shift calculated with the potentials

•
$$\sigma_{tot} = \frac{4\pi}{k^2} \sin^2[\delta(k \to 0)]$$

Yukawa: $\sigma_{tot} = (2.5 - 4.7)\Lambda^{-2}$ (stat.)
2-Gaussian: $\sigma_{tot} = (14 - 51)\Lambda^{-2}$ (stat.)
• $\sigma_{tot} = (2 - 51)\Lambda^{-2}$ (stat.and sys.)

(sys. due to fitting forms)

Constraint on SU(N) YM scale parameter from DM X section

Observational constraints:

 $\frac{\sigma_{\rm tot}}{m_{\Phi}} < 1.0 \ \rm cm^2/g$

Robust constraint from galactic cluster shape, collisions (upper limit)

A. H. Peter et al., MNRAS 430, 81 (2013), 430, 105 (2013); S. W. Randall et al., APJ 679, 1173 (2008).



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Constraint on SU(N) YM scale parameter from DM X section

Observational constraints:

0.45 cm²/g <
$$\frac{\sigma_{tot}}{m_{\Phi}}$$
 < 1.0 cm²/g

Robust constraint from galactic cluster shape, collisions (upper limit)

A. H. Peter et al., MNRAS 430, 81 (2013), 430, 105 (2013); S. W. Randall et al., APJ 679, 1173 (2008).

Constraint from Spergel et al. (lower limit), under discussion?

D. N. Spergel et al., PRL 84, 3760 (2000).



- Glueballs of the SU(N) Yang-Mills theory are good candidates of dark matter : study of self-interaction is important.
- We studied the glueball cross section in the SU(2) Yang-Mills theory on lattice. HALQCD method is used to extract the interglueball potential.
- We could constrain the scale parameter of SU(2) YMT for the 1st time from observational data : $\Lambda > 60$ MeV.

Homeworks:

- Extract glueball effective lagrangian and predict other cosmologically important observables.
- Calculations for $N_c > 2$: extrapolate to large N_c .