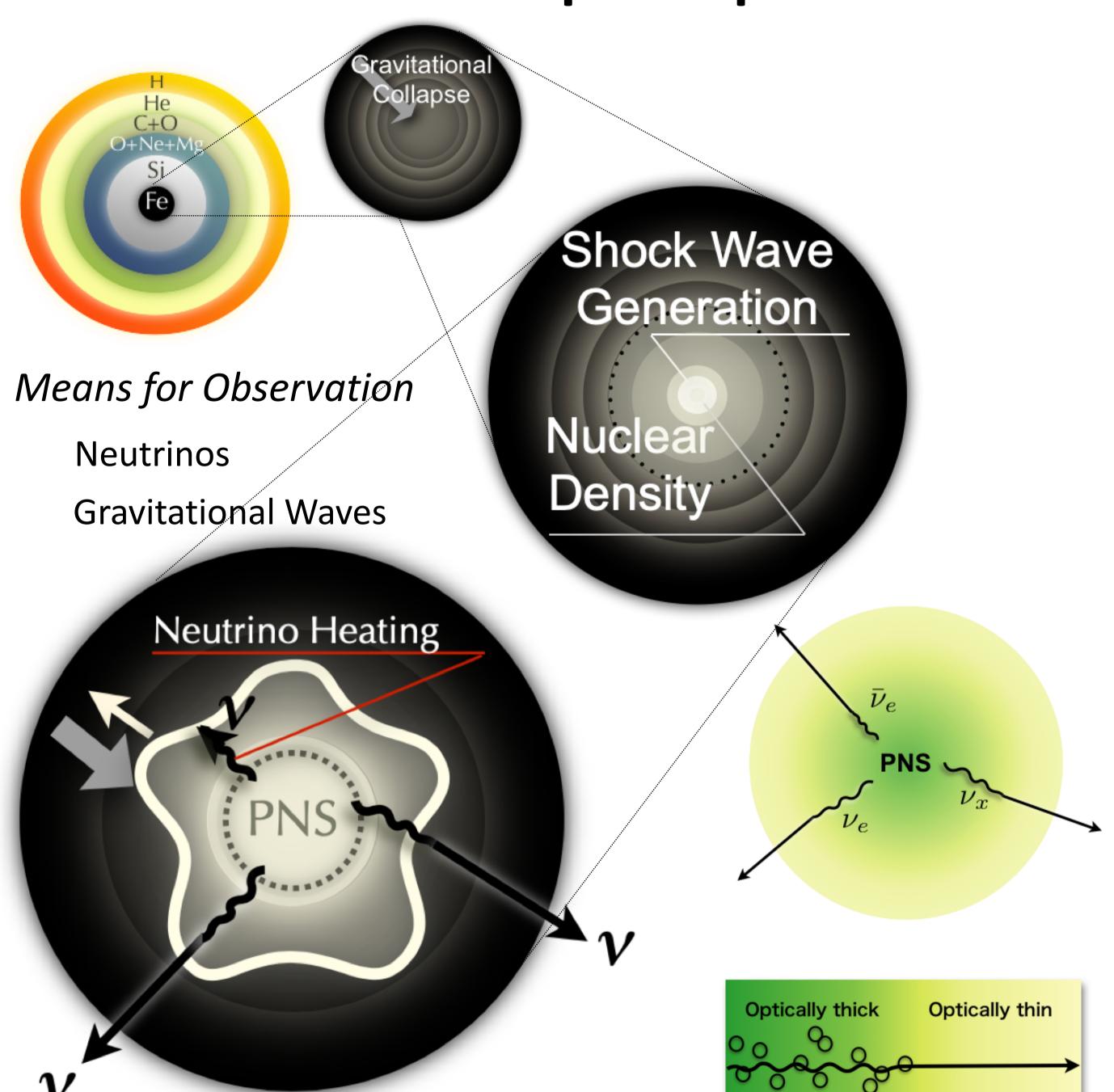
A Simulation of Core-collapse Supernovae in Three-dimensional Space with Full Boltzmann Neutrino Transport on the Supercomputer FUGAKU

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Neutrino Heating Mechanism for Core-collapse Supernovae



Euler Equations

Hydrodynamics

Continuity Equation:

$$\frac{\partial \rho}{\partial t} + \frac{\partial}{\partial x^{j}} \left(\rho v^{j} \right) = 0$$

Equations of Motion:

$$\frac{\partial}{\partial t} (\rho v_i) + \frac{\partial}{\partial x^j} (\rho v_i v^j + P \delta_i^j) = -\rho \frac{\partial \psi}{\partial x^j} - G^i$$

Energy Equation:

$$\frac{\partial}{\partial t} \left(\frac{1}{2} \rho v^2 + e \right) + \frac{\partial}{\partial x^j} \left[\left(\frac{1}{2} \rho v^2 + e + P \right) v^j \right] = -\rho v^j \frac{\partial \psi}{\partial x^j} - G^0$$

Time-Evolution Equation of Electron Number:
$$\frac{\partial}{\partial t} \left(\frac{\rho Y_e}{m_{\star}} \right) + \frac{\partial}{\partial x^j} \left(\frac{\rho Y_e}{m_{\star}} v^j \right) = -\Gamma$$

Poisson's equation for gravity:

$$\Delta \psi = 4\pi G \rho$$

EOS table of Nuclear Matter

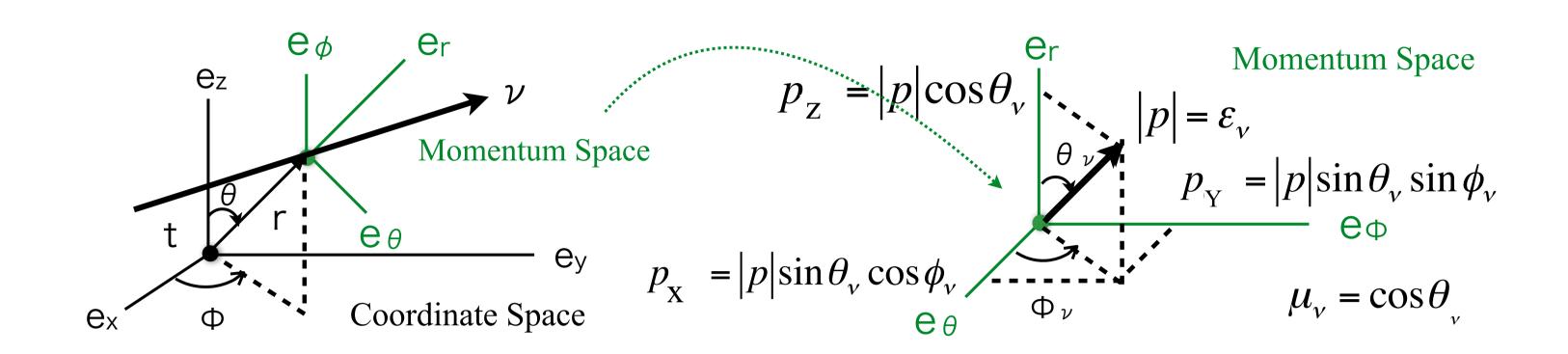
(LS K=220MeV, Furusawa+Togashi, etc.)

 ρ : density, v: velocity, P: pressure, e: internal energy, ψ : the gravitational potential, G: the gravitational constant (=6.67 × 10^{-8} [cm³g⁻¹s²]), Y_e: electron fraction, m_A : the atomic mass unit, G^0 : neutrino radiation energy, G^i : neutrino radiation pressure, Γ : deleptonization rate ($\equiv \Gamma_{v} - \Gamma_{\bar{v}}$), Γ_{s} : neutrino number density

$\frac{\partial f}{\partial t} + \cos\theta_v \frac{\partial f}{\partial r} + \frac{\sin\theta_v \cos\theta_v}{r} \frac{\partial f}{\partial \theta} + \frac{\sin\theta_v \sin\phi_v}{r} \frac{\partial f}{\partial \phi}$

Boltzmann Equation

Neutrino Radiation



Neutrino distribution function

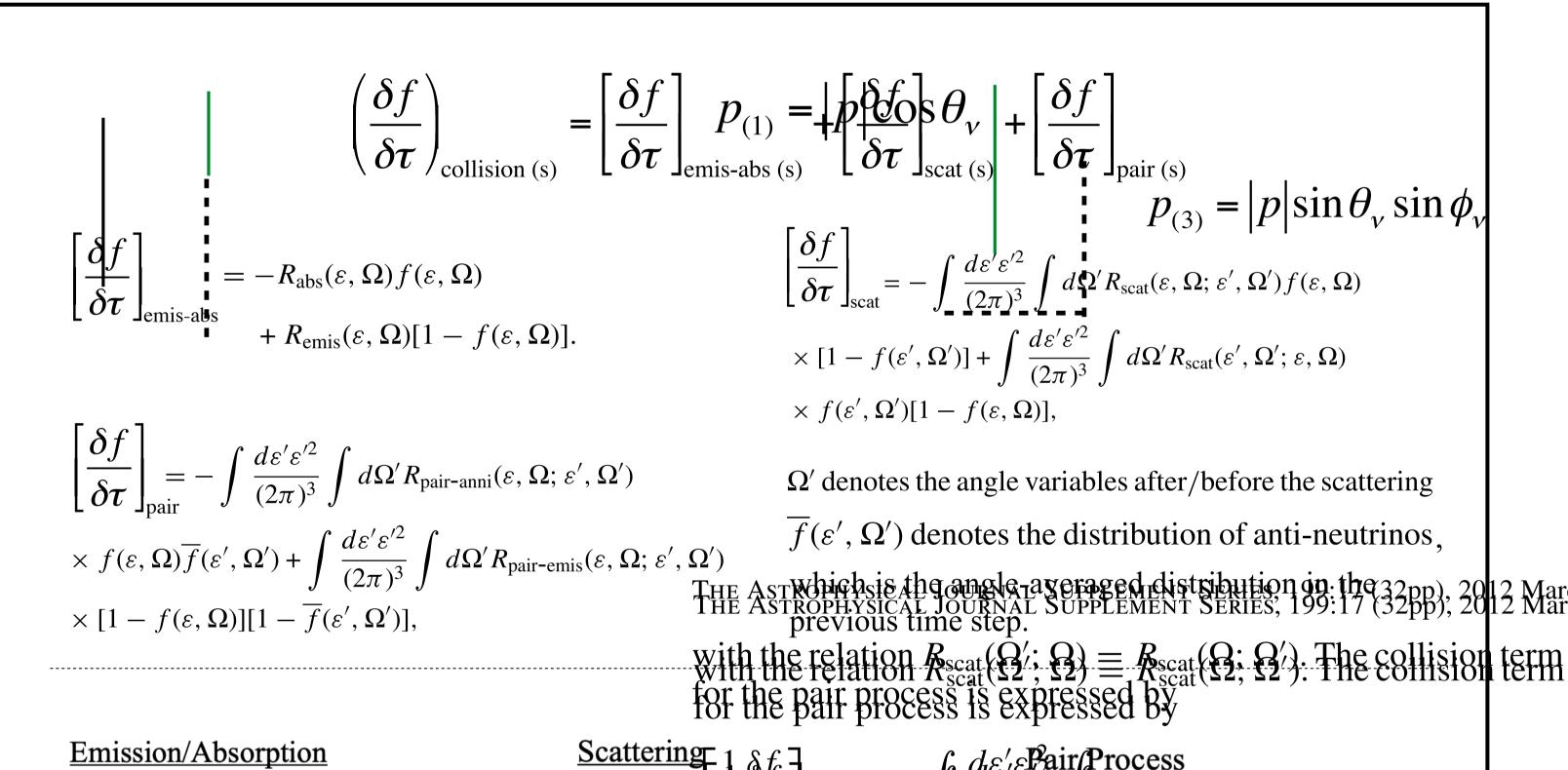
$$f(t, r, \theta, \phi; \varepsilon_v, \mu_v, \phi_v)$$

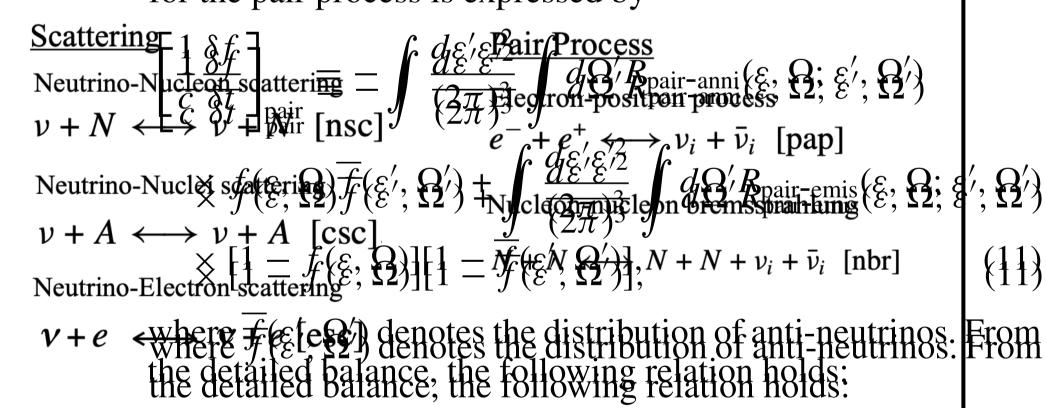
Boltzmann Equation

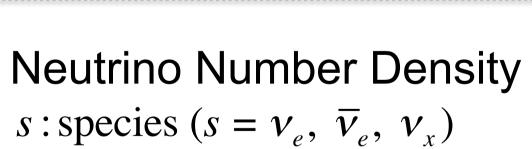
$$\frac{dx^{\mu}}{d\lambda} \frac{\partial f}{\partial x^{\mu}} + \frac{dp^{i}}{d\lambda} \frac{\partial f}{\partial p^{i}} = \left(\frac{\delta f}{\delta \lambda}\right)_{\text{collision}}$$

Boltzmann Equation in the spherical coordinate
$$\frac{\partial f}{\partial t} + \cos \theta_{v} \frac{\partial f}{\partial r} + \frac{\sin \theta_{v} \cos \theta_{v}}{r} \frac{\partial f}{\partial \theta} + \frac{\sin \theta_{v} \sin \phi_{v}}{r \sin \theta} \frac{\partial f}{\partial \phi}$$

$$-\frac{\sin \theta_{v}}{r} \frac{\partial f}{\partial \theta_{v}} - \frac{\cos \theta}{\sin \theta} \frac{\sin \theta_{v} \sin \phi_{v}}{r} \frac{\partial f}{\partial \phi_{v}} = \left(\frac{\delta f}{\delta t}\right)_{\text{collision}}$$







 $e^- + p \longleftrightarrow \nu_e + n \text{ [ecp]}$

 $e^+ + n \longleftrightarrow \bar{\nu}_e + p \text{ [aecp]}$

 $e^- + A \longleftrightarrow \nu_e + A'$ [eca]

Electron Capture

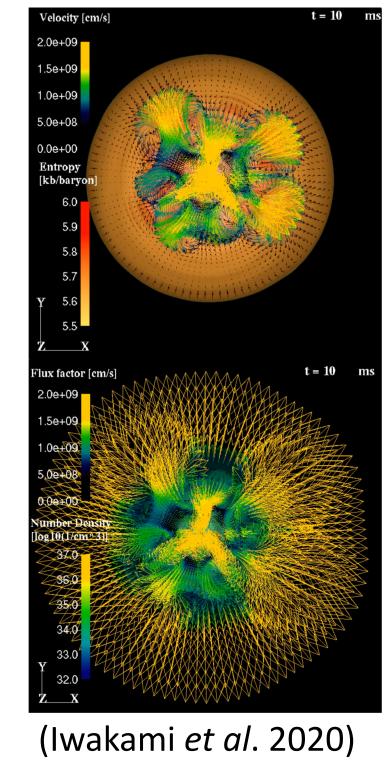
Anti-Electron Capture

Electron Capture on nuclei

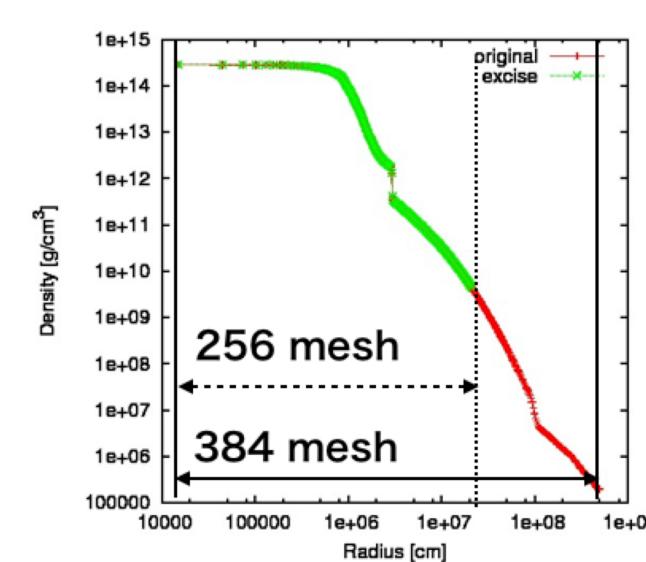
Neutrino Energy Density ($\mu = 0$) Radiation Pressure ($\mu = 1, 2, 3$)

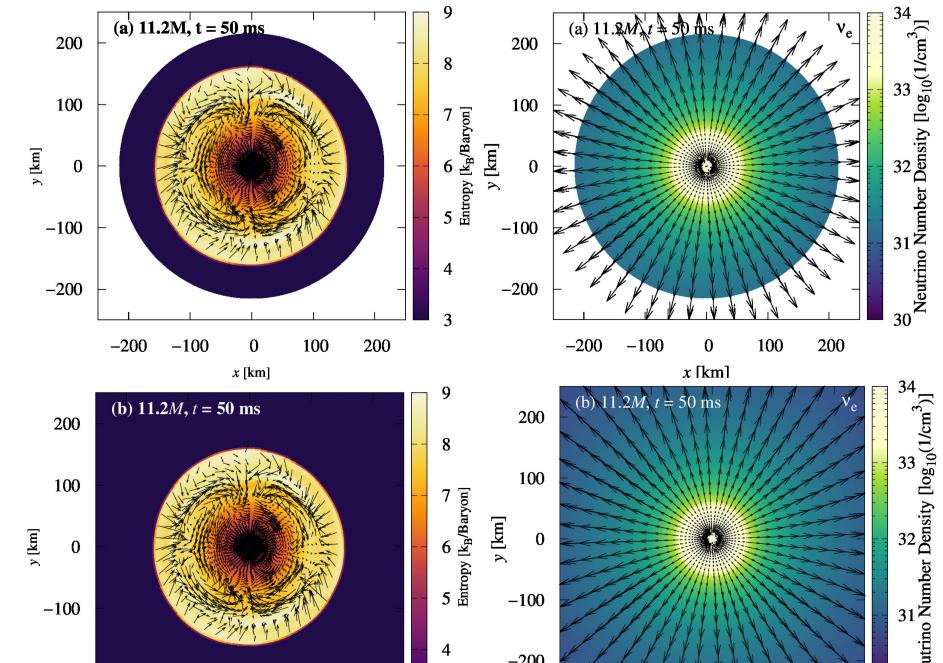
FUGAKU computer

K computer



 $Nr \times N\theta \times N\phi \times Nve \times Nv\theta \times Nv\phi$ $= 256 \times 48 \times 96 \times 16 \times 6 \times 6 \quad 0-50$ ms $= 384 \times 48 \times 96 \times 16 \times 6 \times 6 \times 50 - 300$ ms





 $R_{\text{pair-anni}}(\mathcal{E}, \Omega; \mathcal{E}', \Omega') \equiv R_{\text{pair-emis}}(\mathcal{E}, \Omega; \mathcal{E}', \Omega') \mathcal{E}^{\beta(\mathcal{E}+\mathcal{E}')}$. (12) We linearize the collision term, Equation (11), by assuming that the distribution for anti-peutrings is given by that in the previous time step or the equilibrium distribution. This is a good approximation since the pair process is dominant only in hightemperature regions, where neutrinos are in thermal equilibrium: We adopt the appreach with the distribution in the previous time step in all processes in step in processes in the current study. We mutilize further the angle average of the distribution when we take the isotropic emission rate as we will state: We have also tested that the approach with the equilibrium distribution determined by the local temperature and chemical potential works equally well:

simulations for core-collapse supernovae until t = 50 m/s where the shock wave

arrives at r \sim 150 km. $\mathcal{E}^{+} \downarrow \mathcal{H} \iff \bar{\mathcal{V}}_{e} \downarrow \mathcal{H} \text{ [aseb]};$

Then I extend the computational region from $r = 200 \text{ km}^+\text{ (panel (s))} + 60 5000 \text{ km}$ for the absorption/emission;

 $V \pm N \iff V \pm N \text{ [ASE]}$

The values between 200 km and 5000 km at

(13)

(14)

(13)

 $\{18\}$

t = 50 ms are given by the total tassell. 1D

for the listioners evid is lattering. We do enoughte into account The cheuting election is minor although it contributes to the hermalization (Burrows et al. 2008a): As for the pair process; ve take the electron—positron process and the nucleon—nucleon