

# Quo vadis neutrinoless double beta decay?

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Colorado State University

Unraveling the History of the Universe and  
Matter Evolution with Underground Physics

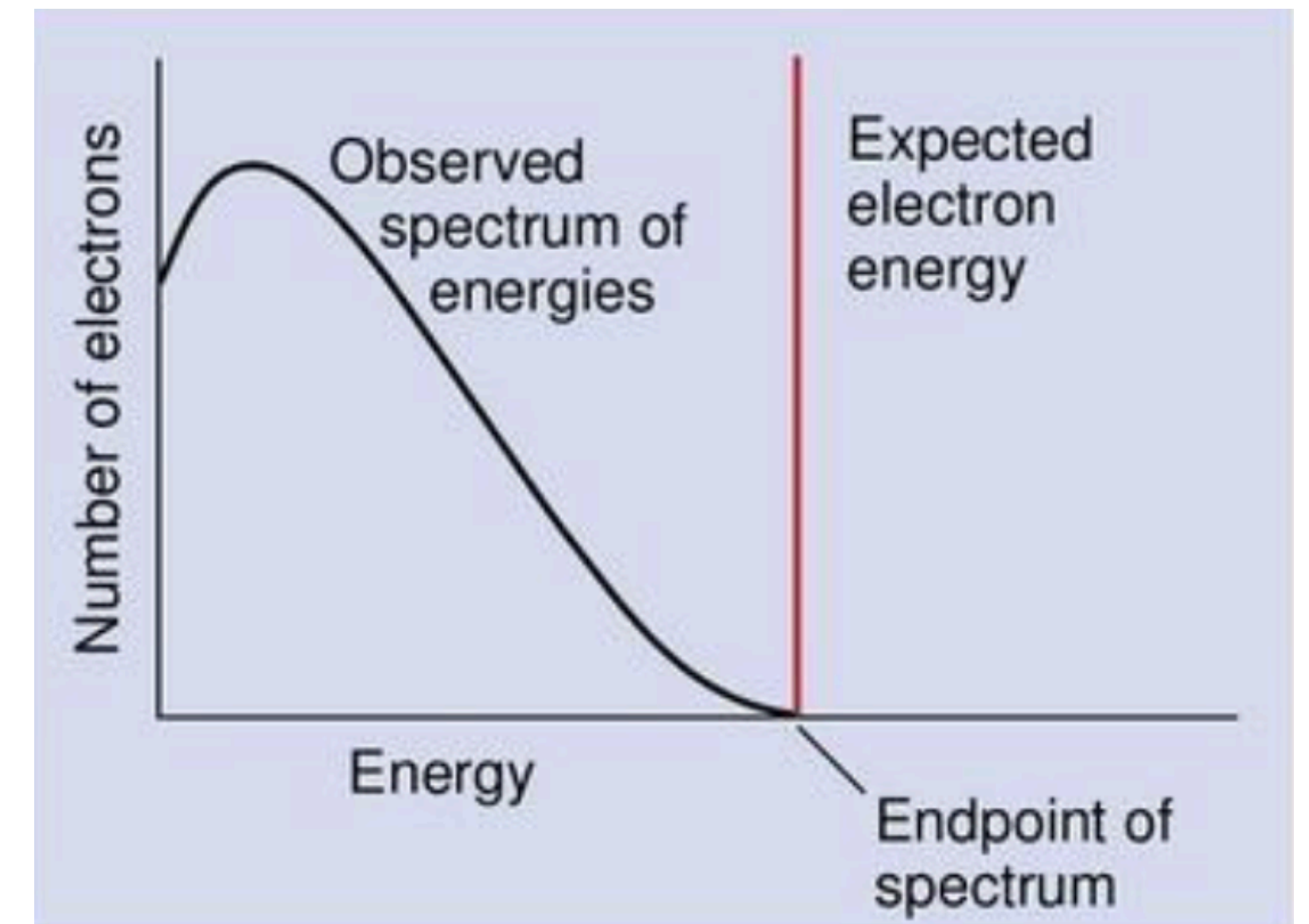
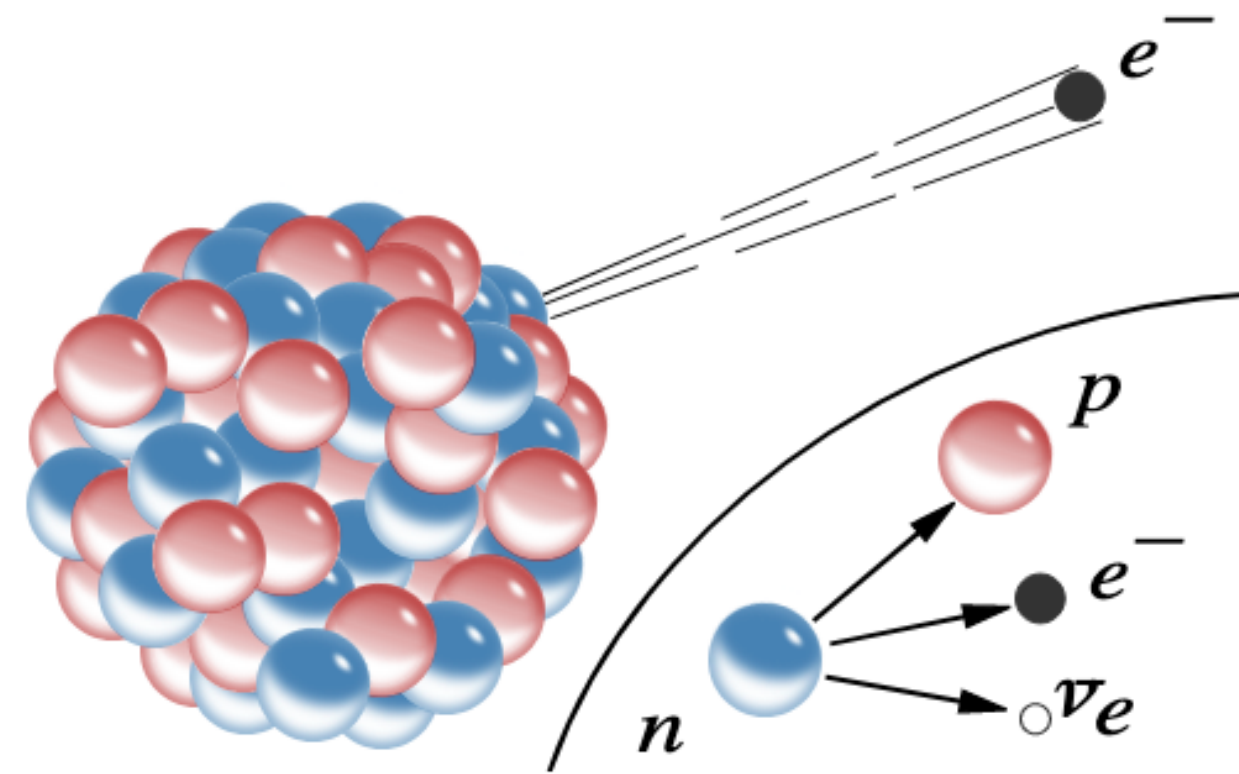
March 4, 2024



**COLORADO STATE  
UNIVERSITY**

# Neutrinos and beta decay

**Beta decay:**  $(Z, A) \rightarrow (Z + 1, A) + e^{-} + \bar{\nu}_e$   
→ **need** to introduce neutrino



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**Double beta decay:**

Maria Goeppert-Mayer (1935) (Nobel prize in 1963)

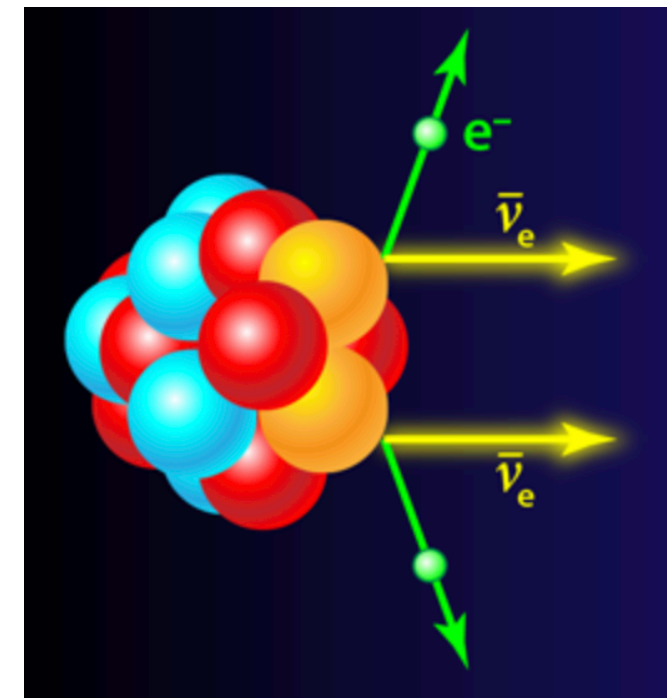
**Simultaneous beta decay** of

two neutrons inside of atomic nucleus:  $(Z, A) \rightarrow (Z + 2, A) + 2 e^- + 2 \bar{\nu}_e$   
SM process!

**Observed** in several isotopes (Ge, Xe, Te, Se)

Double beta decay happens for elements where single beta decay is forbidden  
by energy conservation:

elements with an even atomic number and even neutron number



# Neutrinos and beta decay

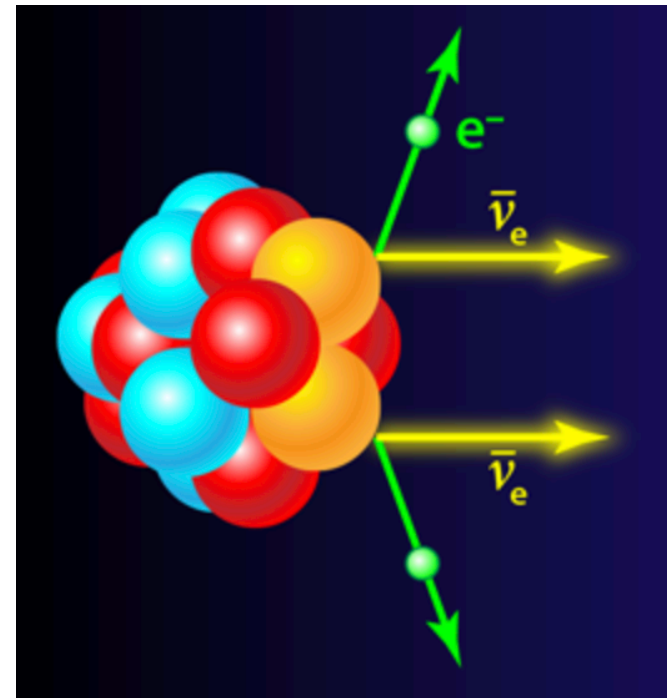
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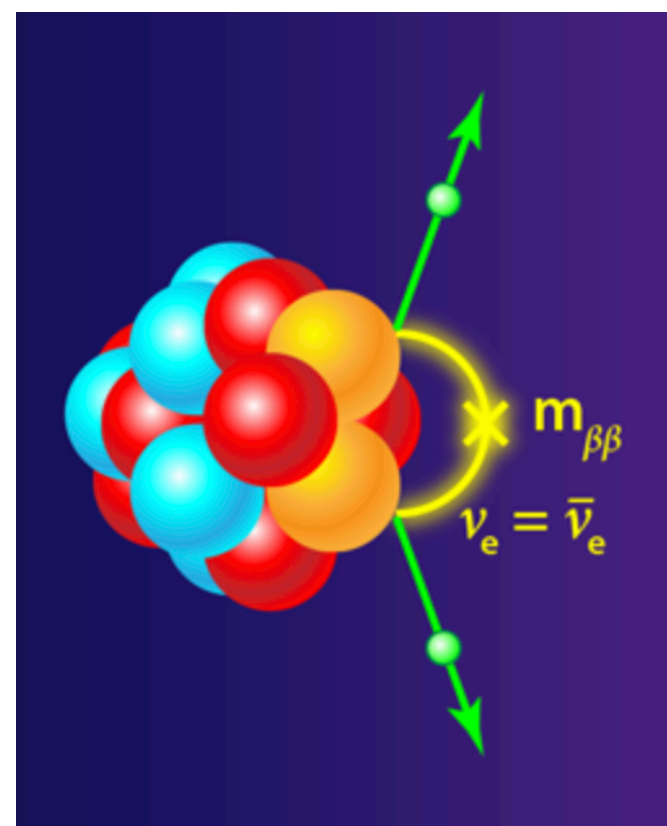
## Neutrinoless double beta decay:

W. Furry (1939)

Neutrinos inside of nucleus emitted and absorbed if they are their own antiparticles: **lepton number violation!**

$$(Z, A) \rightarrow (Z + 2, A) + 2 e^-$$

**BSM** process!

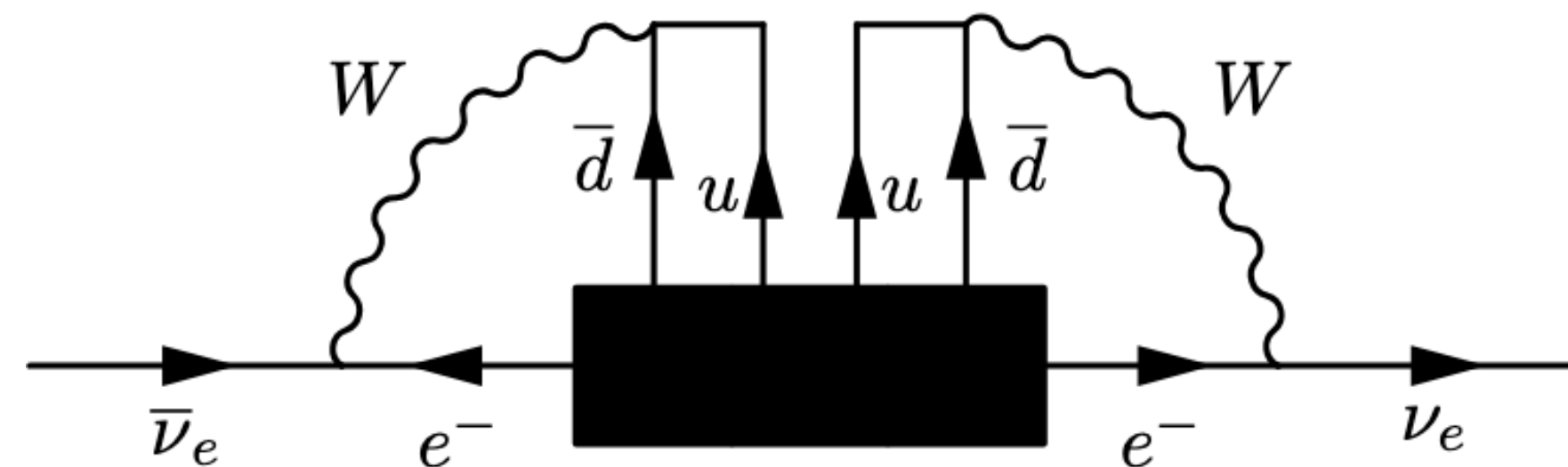


# Neutrinoless double beta decay

Observation of neutrinoless double beta decay:  
⇒ **Lepton number violated!**

## Schechter-Valle theorem

Any  $\Delta L = 2$  operator contributing to  $0\nu\beta\beta$  will generate **Majorana neutrino mass** contribution Schechter, Valle '82



Duerr, Lindner, Merle '11

Neutrino masses might still have (large) Dirac mass term  
Majorana mass induced by this operator tiny:  $\lesssim \mathcal{O}(10^{-28} \text{ eV})$   
Lower limit on second-lightest neutrino mass  $m_\nu \gtrsim 8 \times 10^{-3} \text{ eV}$

# Neutrinoless double beta decay

**Powerful way** of testing lepton number violation!

Individual lepton number ( $L_e, L_\mu, L_\tau$ ) **violated** in neutrino oscillations

total lepton number ( $L_e + L_\mu + L_\tau$ ) (and baryon number) is **accidental**  
symmetry of the SM



# Neutrinoless double beta decay

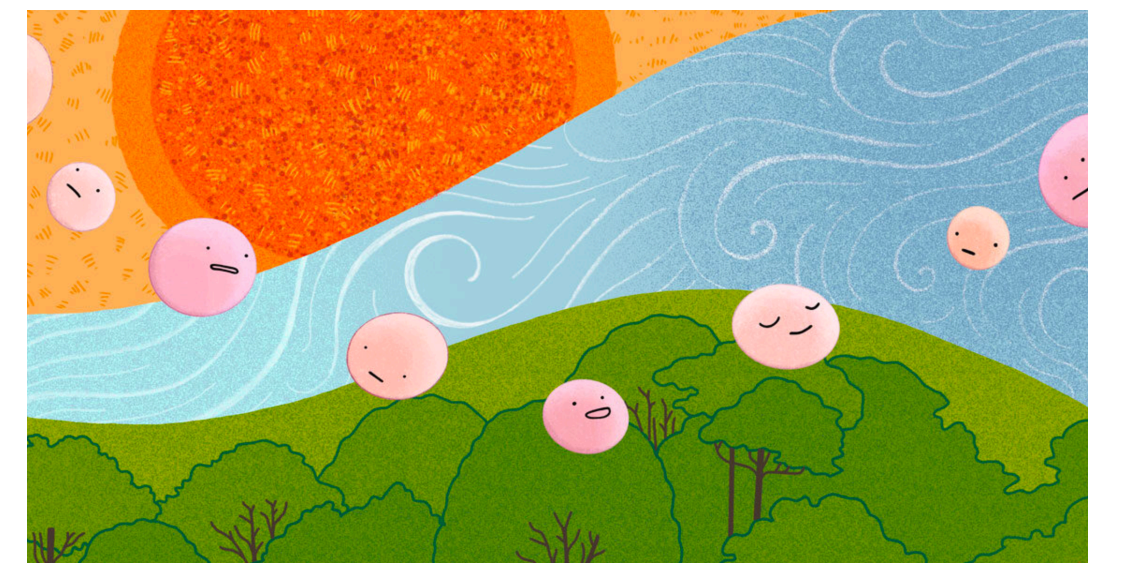
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**Leptogenesis** scenarios to generate matter-antimatter  
asymmetry of the universe rely on lepton number violation

→  $0\nu\beta\beta$  probes history of the Universe and Matter  
Evolution



# Neutrinoless double beta decay

**Powerful way** of testing lepton number violation!

total lepton number ( $L_e + L_\mu + L_\tau$ ) (and baryon number) is  
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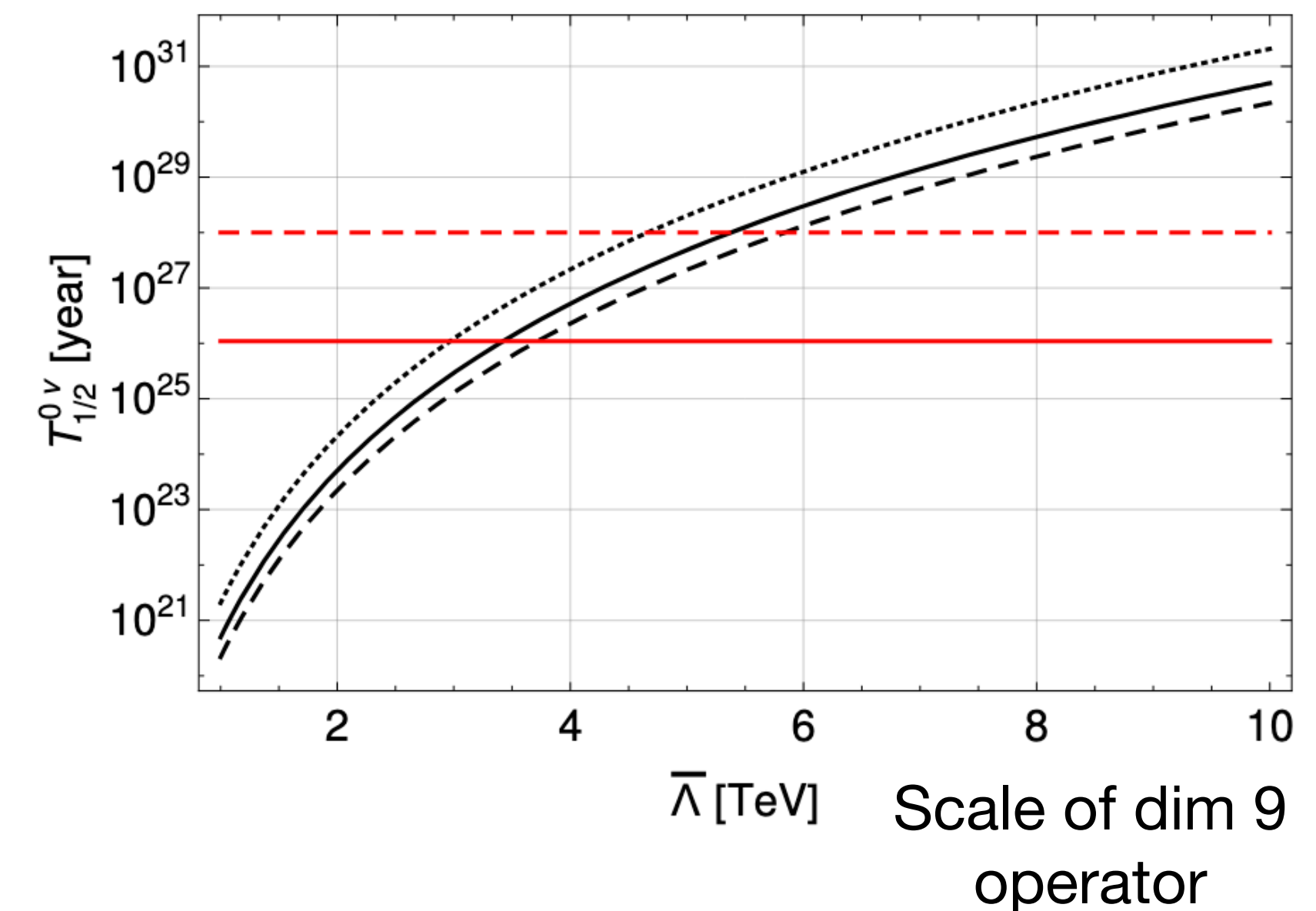


**Leptogenesis** scenarios to generate matter-antimatter asymmetry of the universe rely on lepton number violation

→  $0\nu\beta\beta$  probes history of the Universe and Matter Evolution

lepton number violation could come from an odd-dimensional (dim 5, 7, 9, ...) EFT operator

Lowest dimensional SMEFT operator → Majorana neutrino mass term



Kochbach '16



# Neutrinoless double beta decay

Observable: half-life of isotope

$$T_{1/2}^{-1} = |m_{\beta\beta}|^2 G^{0\nu} |M^{0\nu}|^2$$

Phase space

Particle physics quantity

Nuclear matrix element

# Neutrinoless double beta decay

See talks today and posters tomorrow

**Observable: half-life of isotope**

$$T_{1/2}^{-1} = |m_{\beta\beta}|^2 G^{0\nu} |M^{0\nu}|^2$$

**Phase space**

Under control

**Particle physics quantity**

Under control

**Nuclear matrix element**

Source of uncertainty

See talk by N. Hinohara

# Neutrinoless double beta decay

## Nuclear matrix element

Engel, Menendez '16

**Disagreements** between determinations using different nuclear models

**New idea:**

**Ab-initio** many body methods

start with interactions and

operators determined from QCD

and/or fit to data in very light nuclei

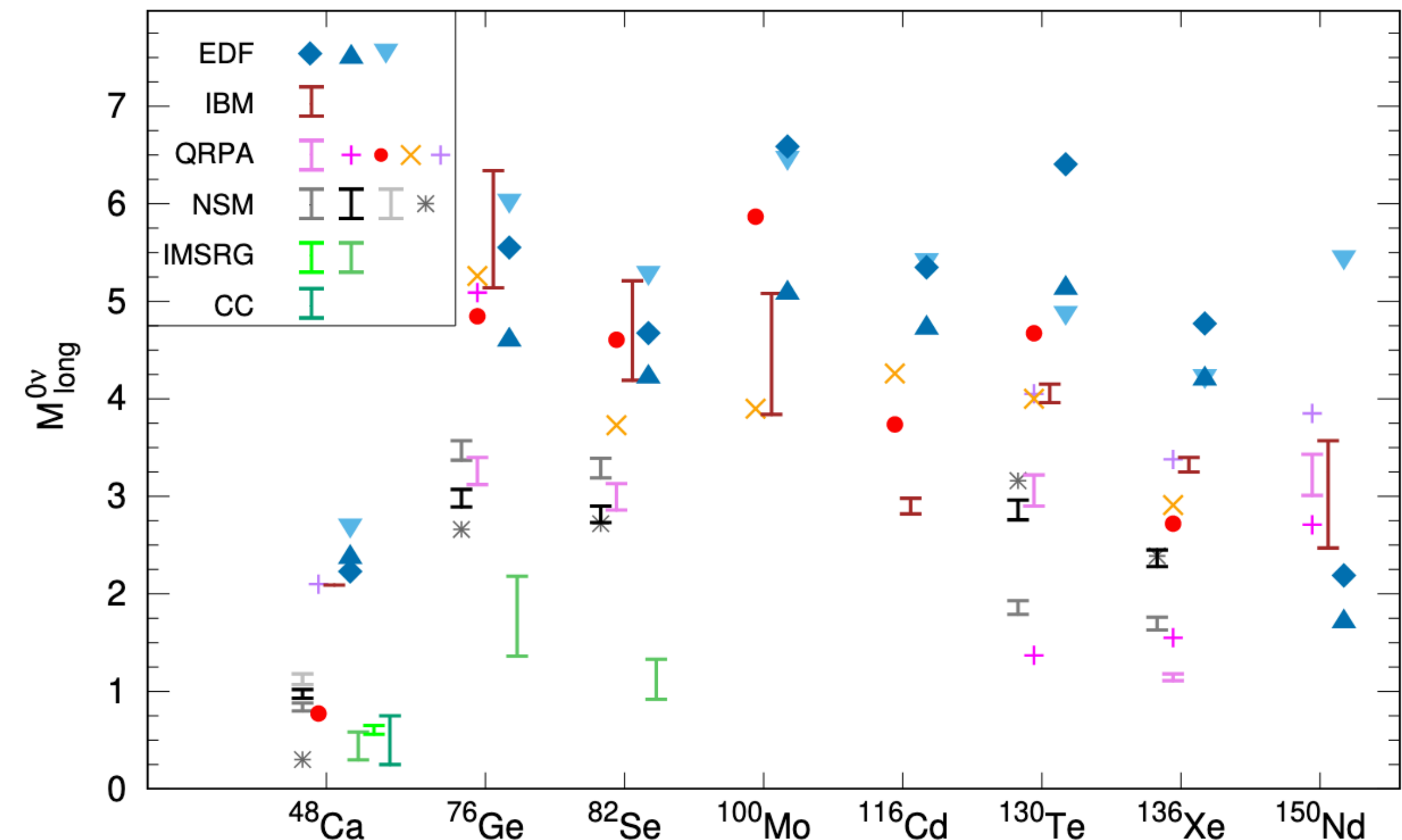
produce solutions to the Schroedinger equation

in heavier nuclei, with systematically

improvable approximations

**Goal: Reliable uncertainty quantification**

Agostini, Benato, Detwiler, Menendez, Vissani '22



# Neutrinoless double beta decay

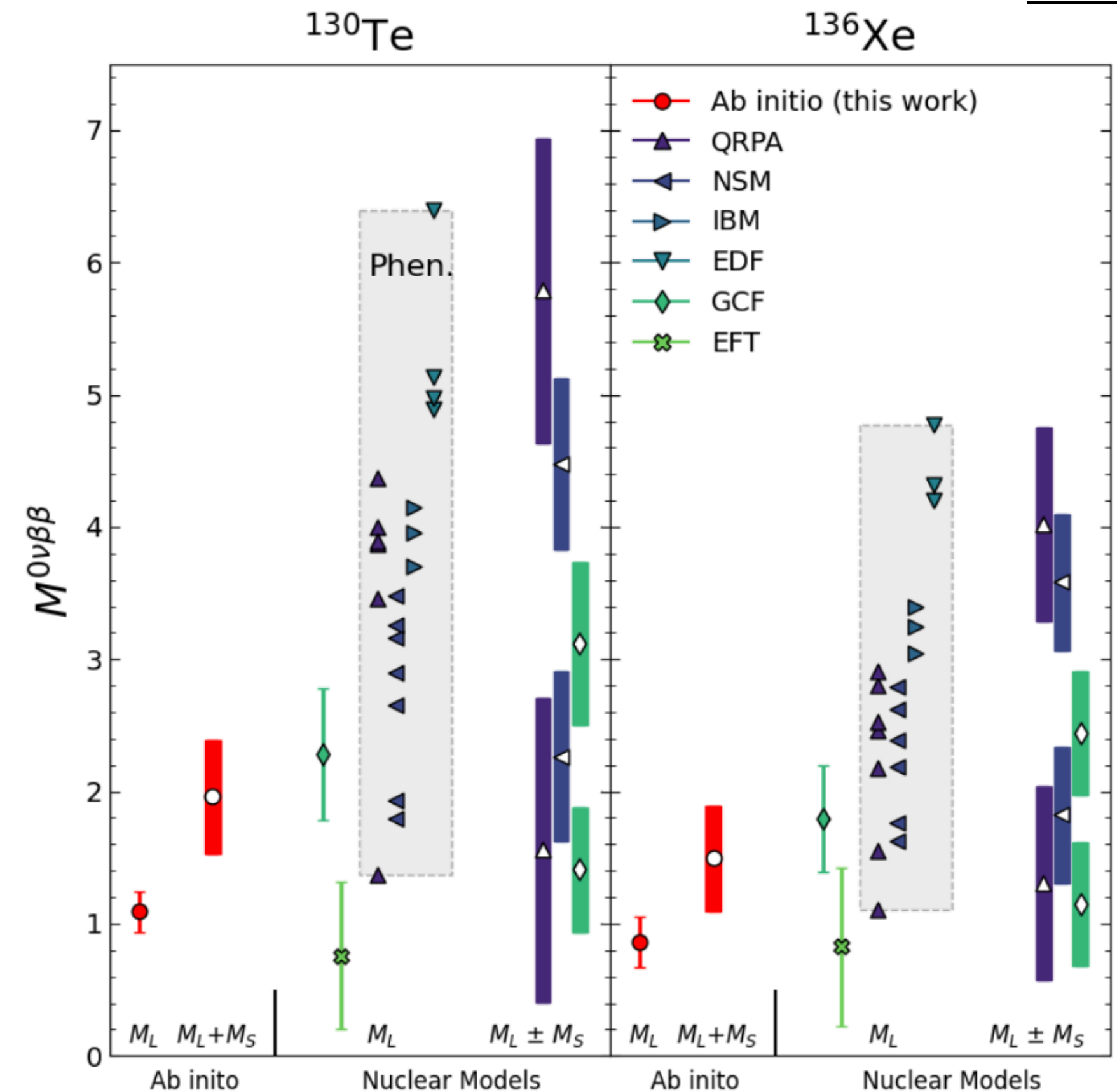
## Nuclear matrix element

Belley et al '23

**Disagreements** between determinations using different nuclear models

**New idea:**

**Ab-initio** many body methods start with interactions and operators determined from QCD and/or fit to data in very light nuclei produce solutions to the Schroedinger equation in heavier nuclei, with systematically improvable approximations



Goal: Reliable uncertainty quantification

first comprehensive ab initio uncertainty quantification last year [Belley et al '23](#)

# Neutrinoless double beta decay

Particle physics quantity

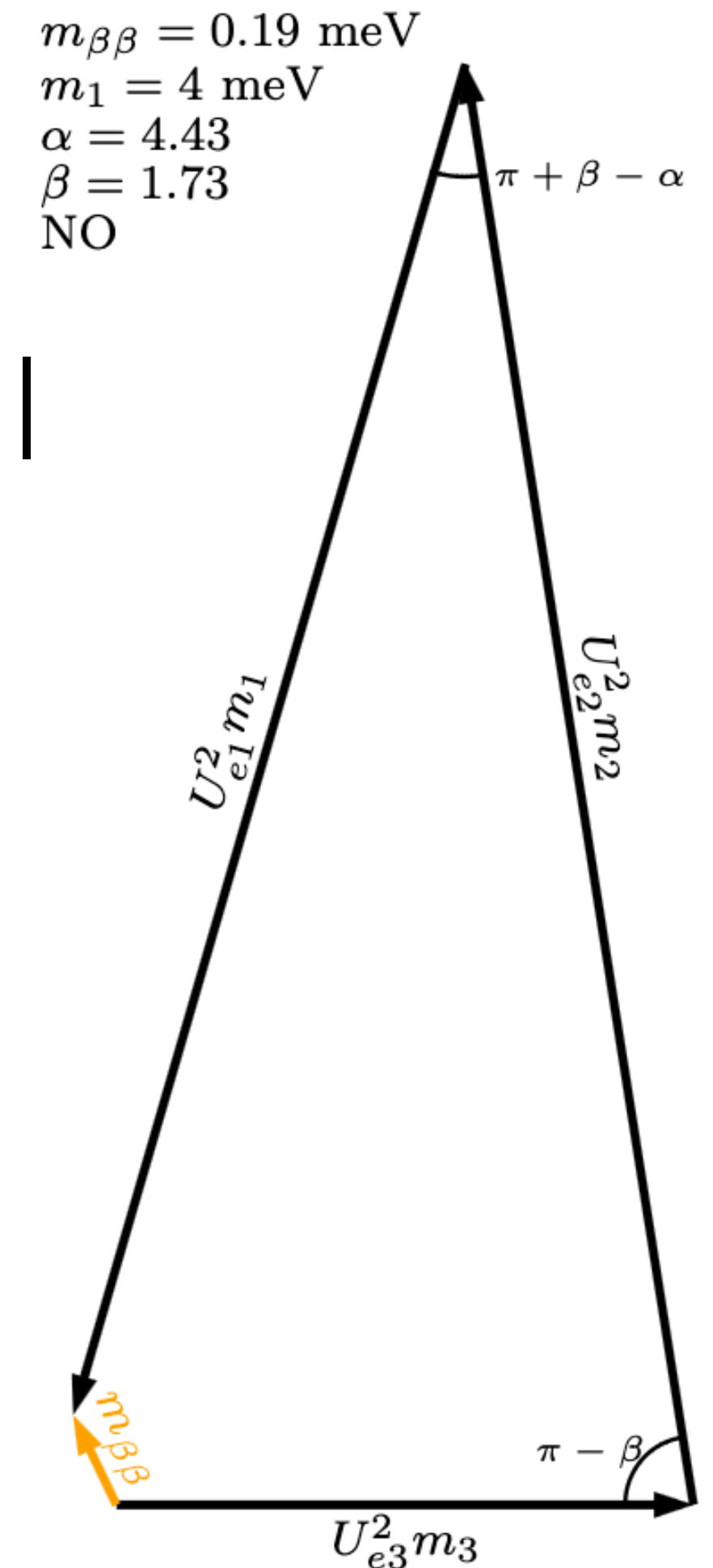
$$|m_{\beta\beta}| = \left| \sum U_{ei}^2 m_i \right|$$

$$= \left| \cos^2 \theta_{12} \cos^2 \theta_{13} m_1 e^{-i\alpha} + \sin^2 \theta_{12} \cos^2 \theta_{13} m_2 e^{-i\beta} + \sin^2 \theta_{13} m_3 \right|$$

With measurement of mixing angles and mass splittings

→ **3 unknowns**: Majorana phases ( $\alpha$ ,  $\beta$ ), mass of lightest neutrino

Only sensitive to a combination of Majorana phases!



# Neutrinoless double beta decay

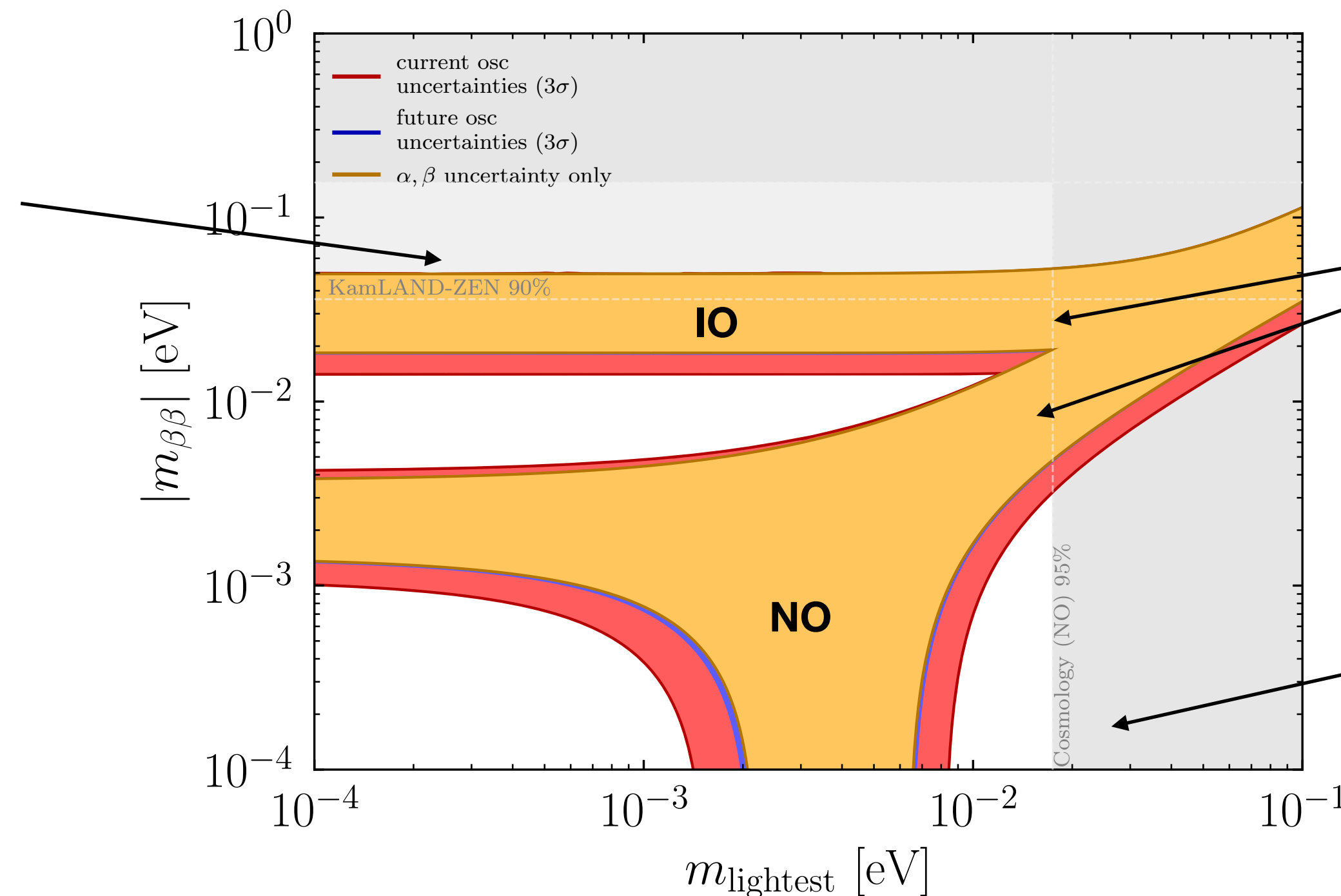
## Particle physics quantity

Denton, JG '23

$$|m_{\beta\beta}| = \left| \sum U_{ei}^2 m_i \right|$$

$$= \left| \cos^2 \theta_{12} \cos^2 \theta_{13} m_1 e^{-i\alpha} + \sin^2 \theta_{12} \cos^2 \theta_{13} m_2 e^{-i\beta} + \sin^2 \theta_{13} m_3 \right|$$

Current bound on  $|m_{\beta\beta}|$



Upcoming oscillation experiments will select MO and slightly decrease parameter space

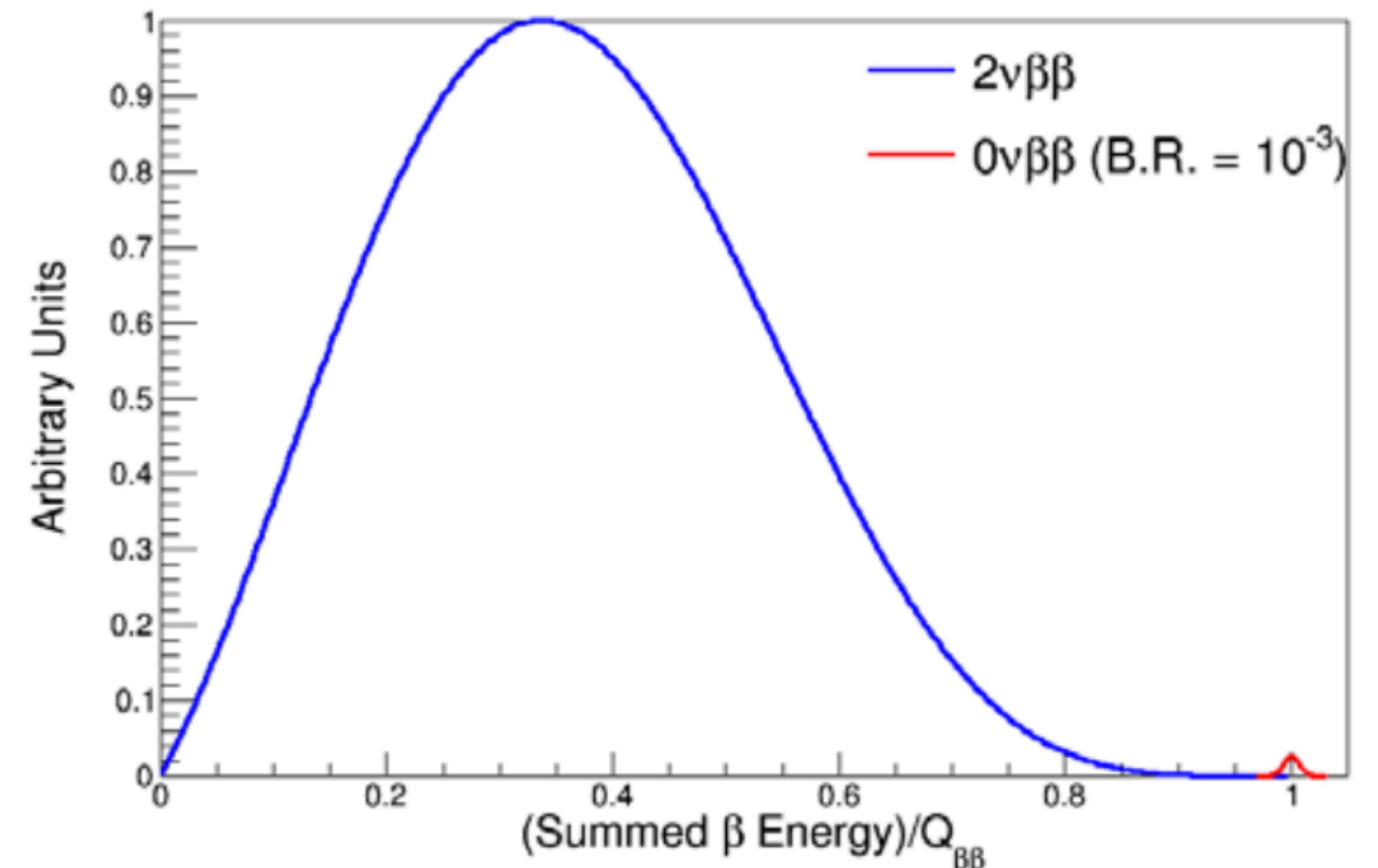
Interplay with cosmology: sum of neutrino masses

# Neutrinoless double beta decay

## Experiment

Observable: 2 emitted electrons (+daughter nucleus)

intrinsic, irreducible background:  $2\nu\beta\beta$



# Neutrinoless double beta decay

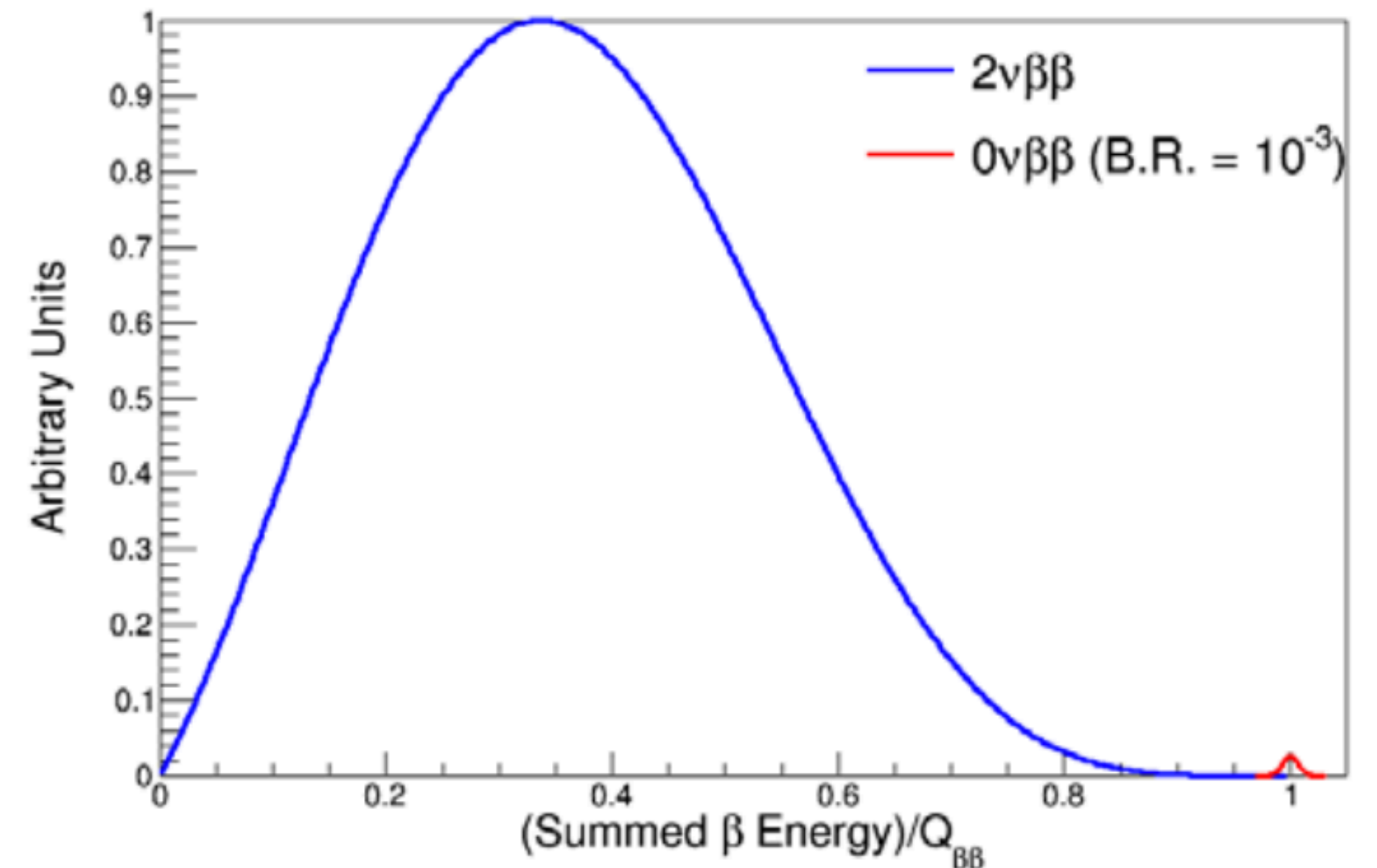
## Experiment

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### Experimental requirements:

- excellent energy resolution
- low backgrounds
- Large detectors (expect one decay per ton-year)
- Long exposure
- Topological information of signal and background





# Neutrinoless double beta decay

## Experiment

KamLAND-Zen '22

**No** observation

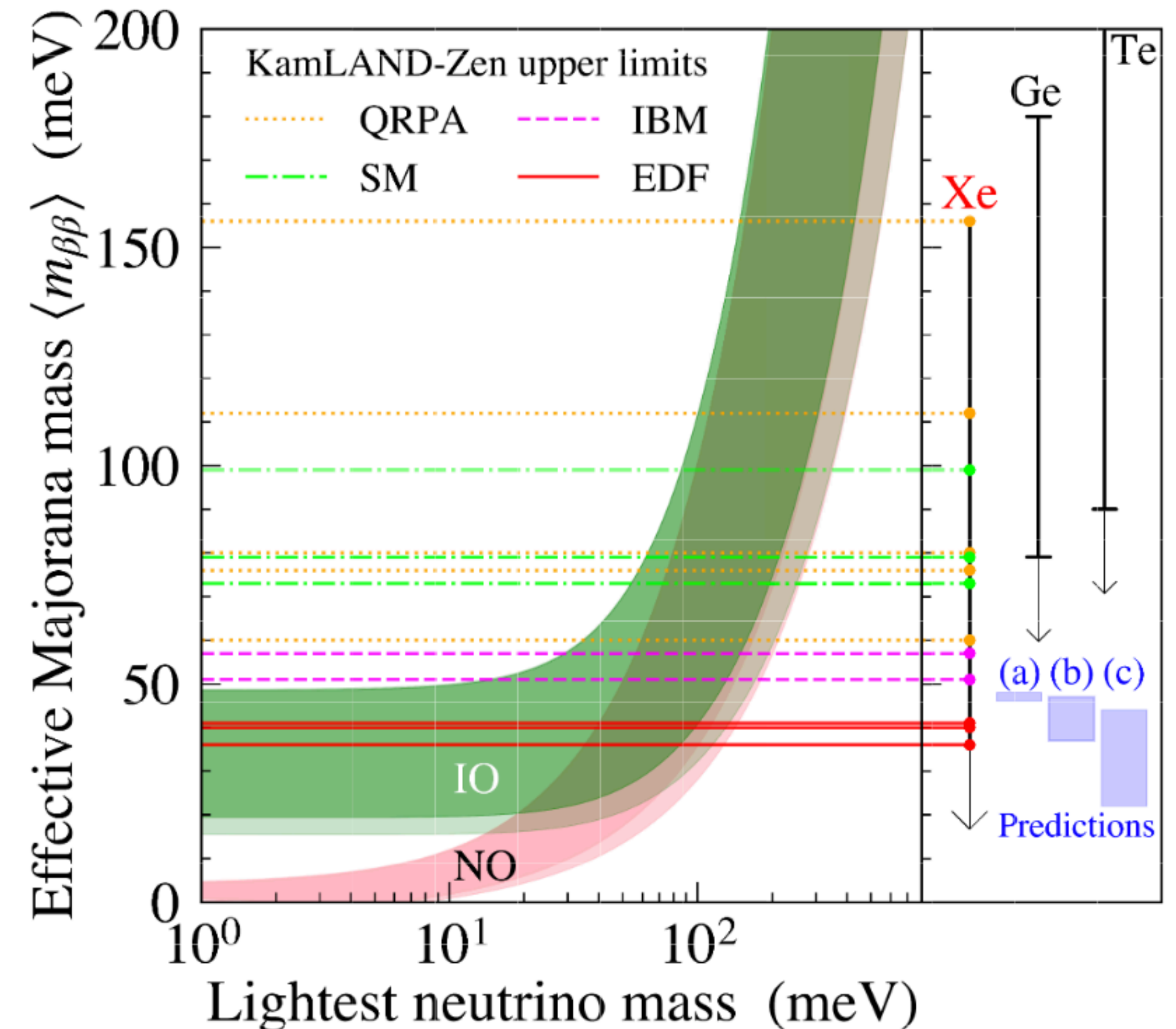
Best constraint from KamLAND-Zen:

$$T_{1/2}^{0\nu} > 2.3 \times 10^{26} \text{ yr}$$

$$\rightarrow m_{\beta\beta} < 36 - 156 \text{ meV}$$

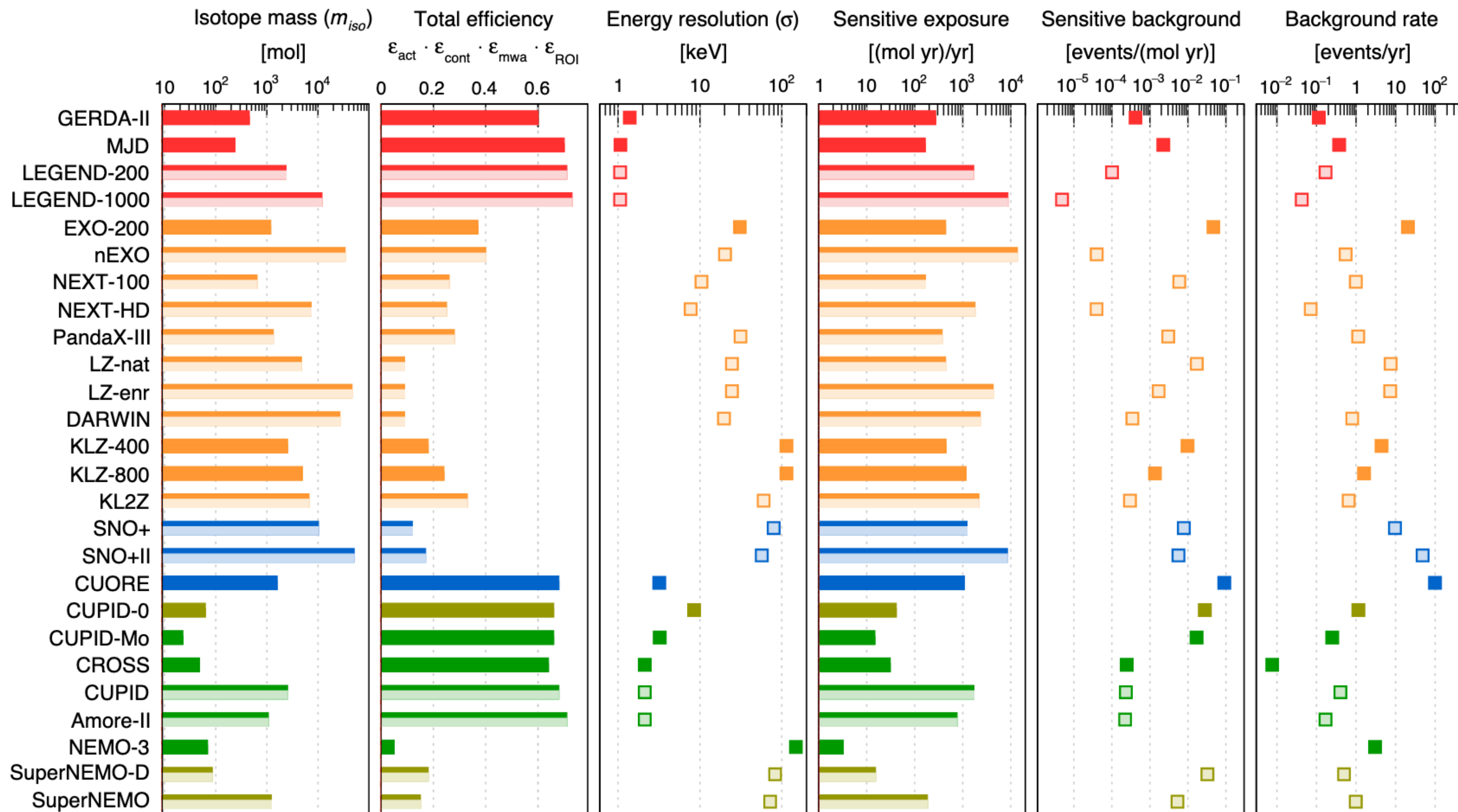
See previous talks

Several other experiments start to probe the IO



# Neutrinoless double beta decay

## Experiment



Experimental prospects  
are **rich**:  
Different isotopes, different  
detection techniques,...

One of top priorities of US  
Nuclear Science Advisory  
Committee long-range plan:  
pursuit of ton-scale  
neutrinoless double beta decay  
experiments  
KamLAND2-Zen: MEXT roadmap  
2023

See talks today

# Neutrinoless double beta decay

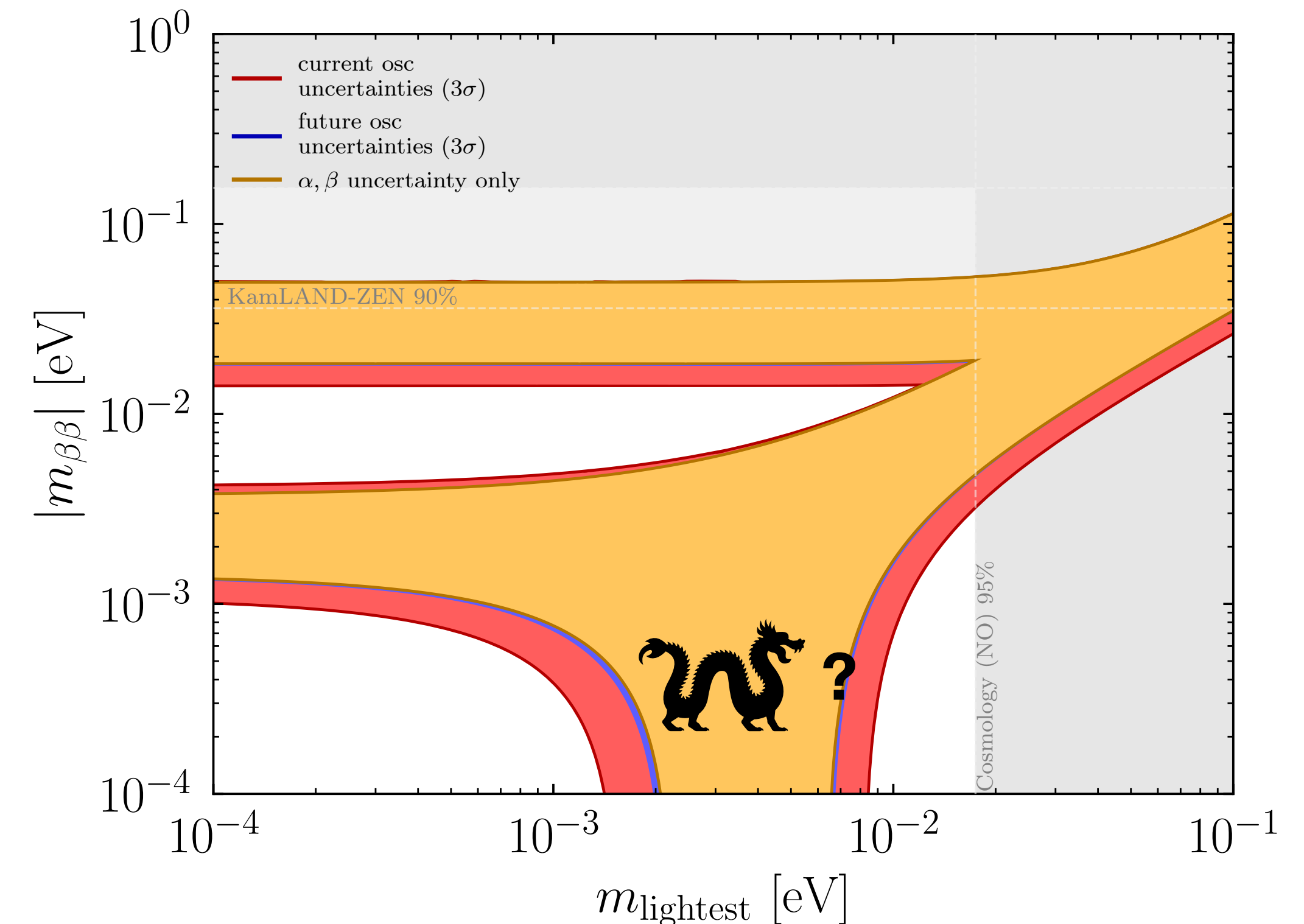
## Future

Where are we going?

Experiments are moving forward

Where are the regions of interest?

Do we need/want to probe down to very small  $m_{\beta\beta}$ ?



# Neutrinoless double beta decay

## Future

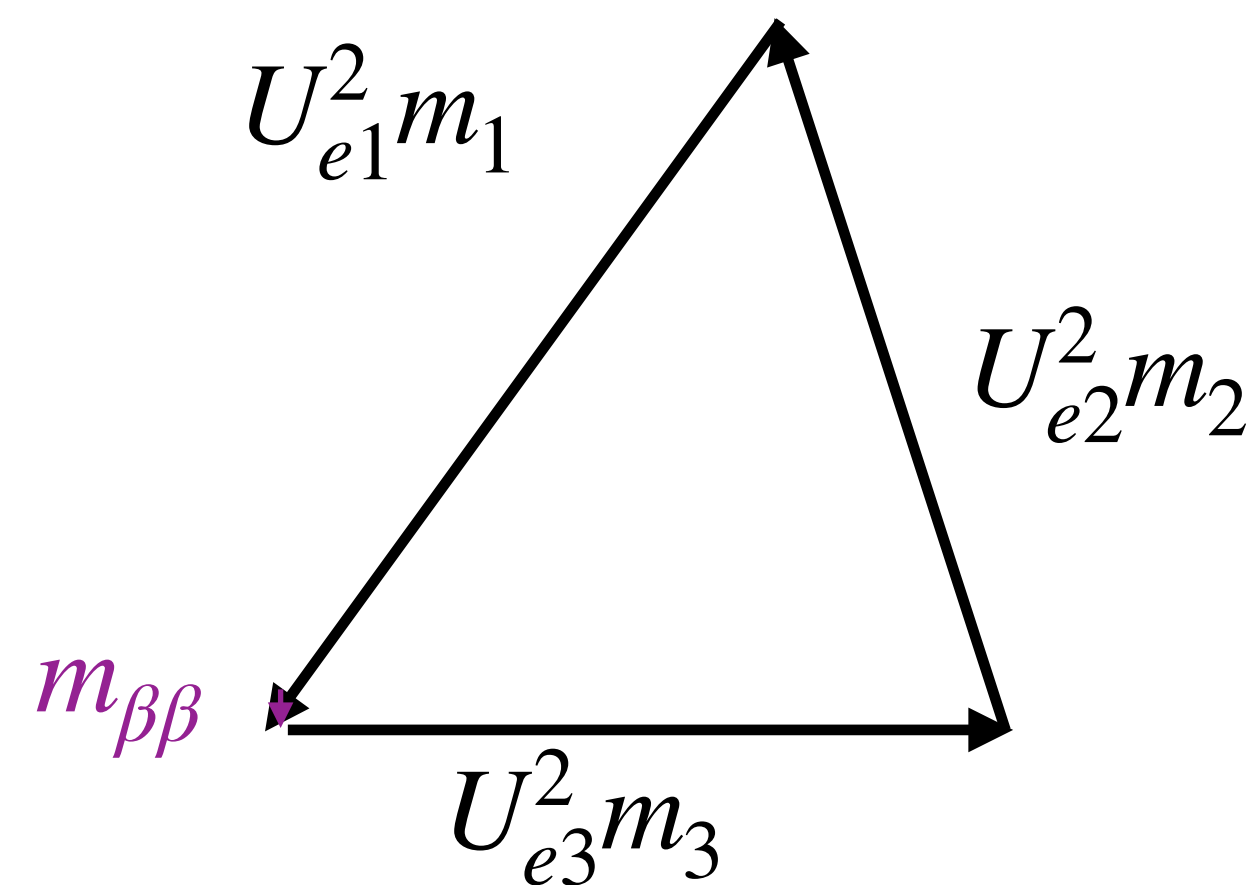
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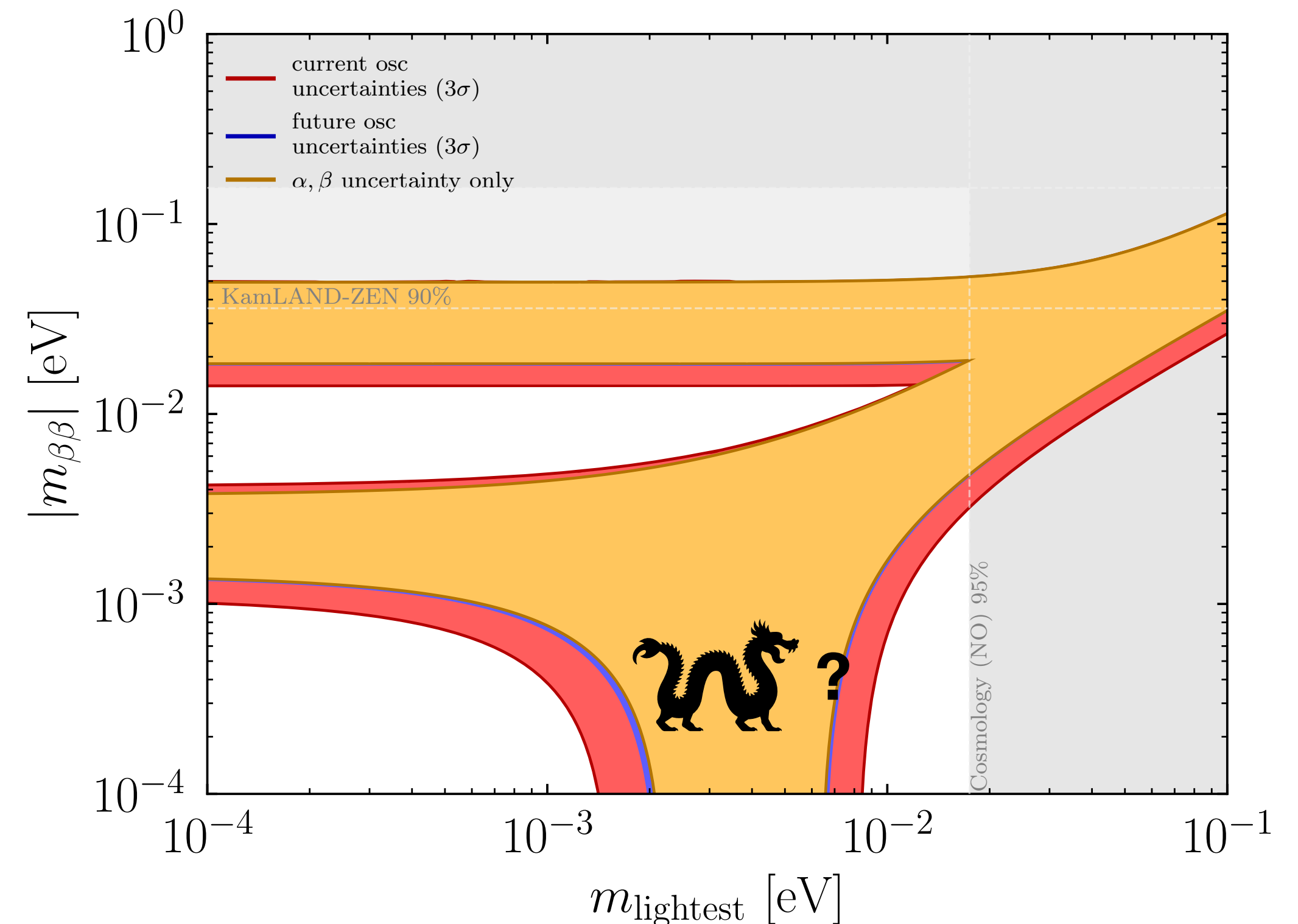
Do we need/want to probe down to very small  $m_{\beta\beta}$ ?

In funnel both Majorana phases  
can be extracted

Ge, Lindner '16



Triangle closes  
from the knowledge of  
length of sides  
we can determine the  
Majorana phases



# Neutrinoless double beta decay

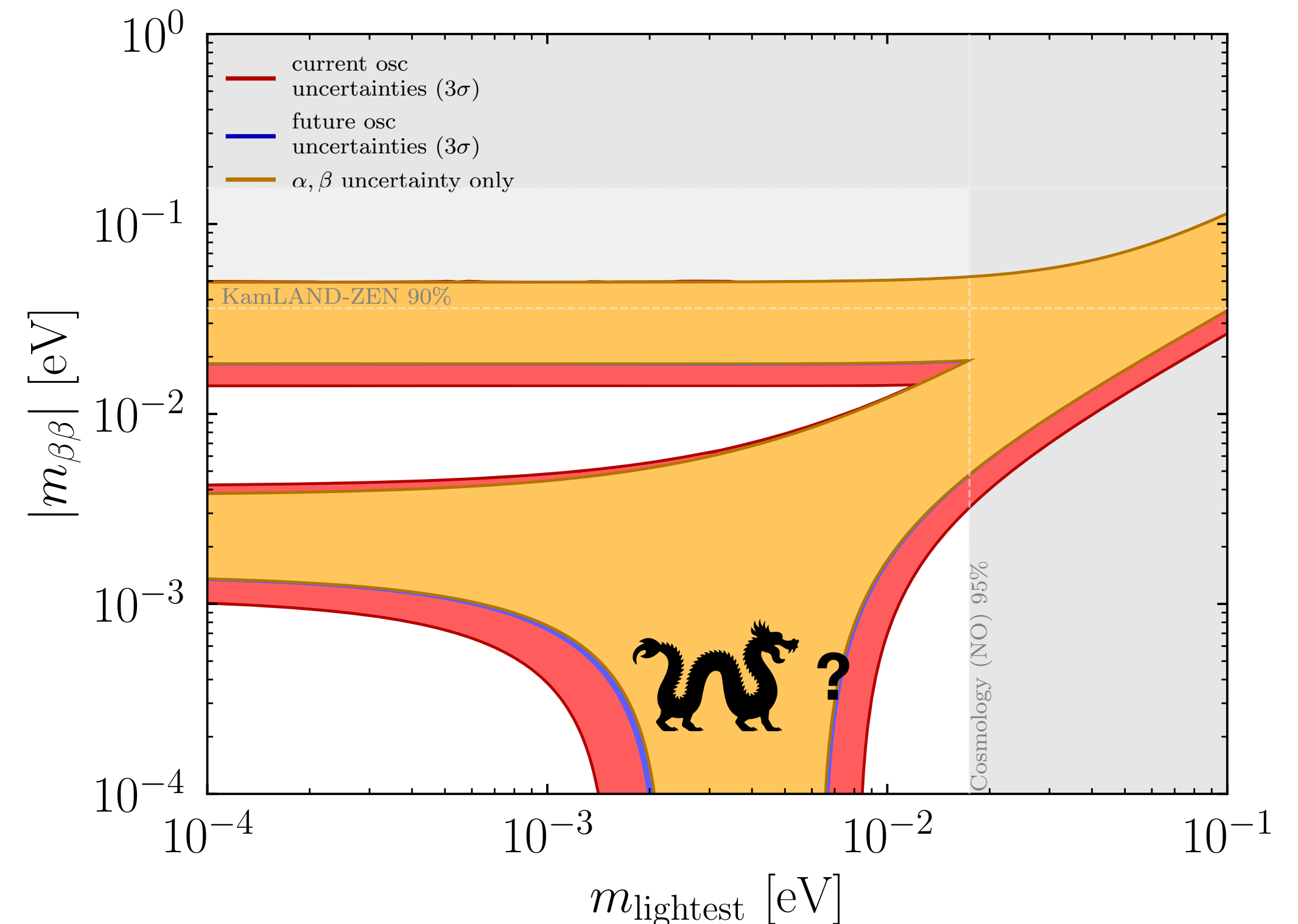
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Do we need/want to probe down to very small  $m_{\beta\beta}$ ?

Majorana phases, lightest mass, MO crucially  
determine allowed regions of  $m_{\beta\beta}$   
→ Predictions from **flavor models**



# Neutrinoless double beta decay

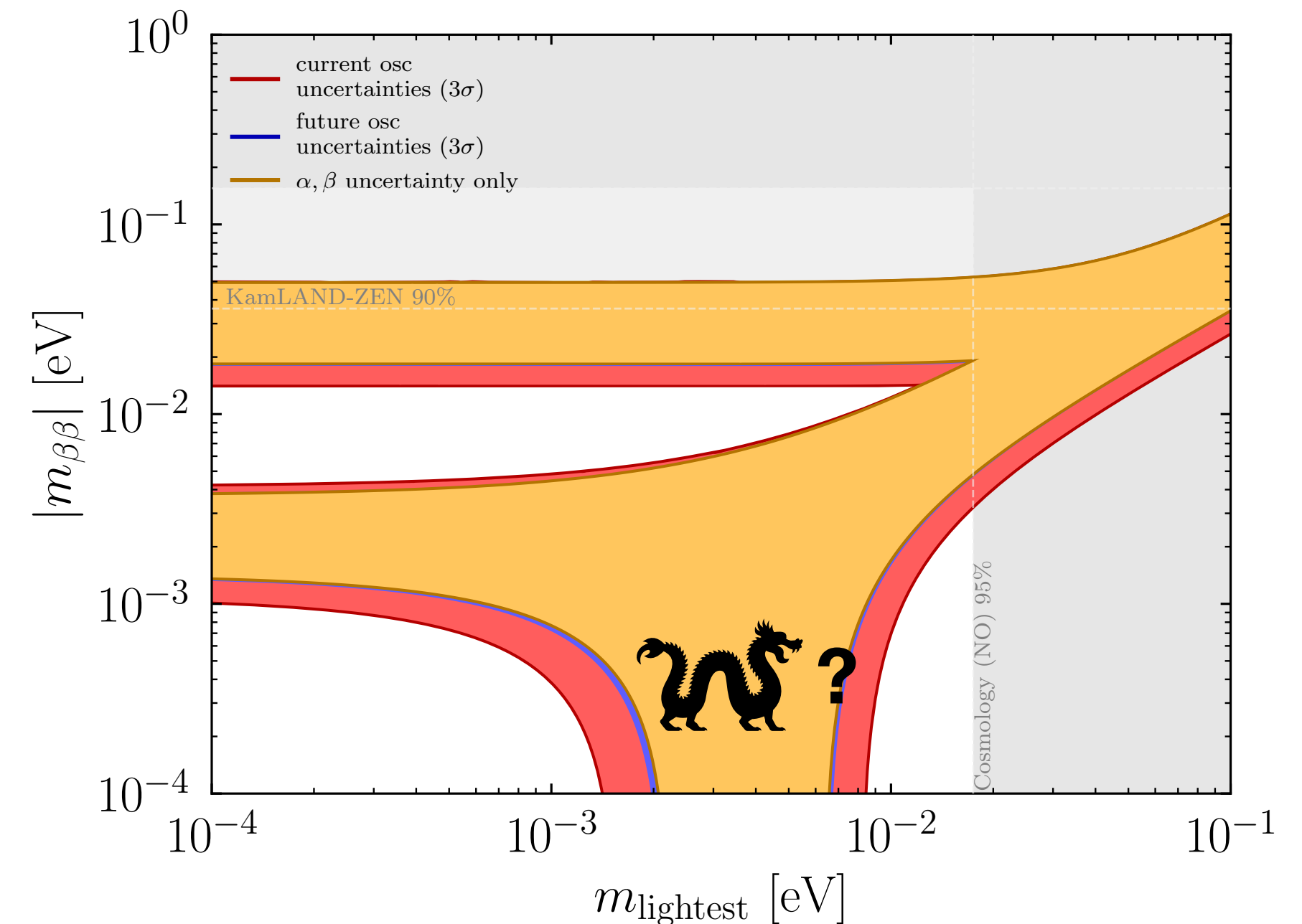
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**A Survey of Neutrino Flavor Models and the  
Neutrinoless Double Beta Decay Funnel**

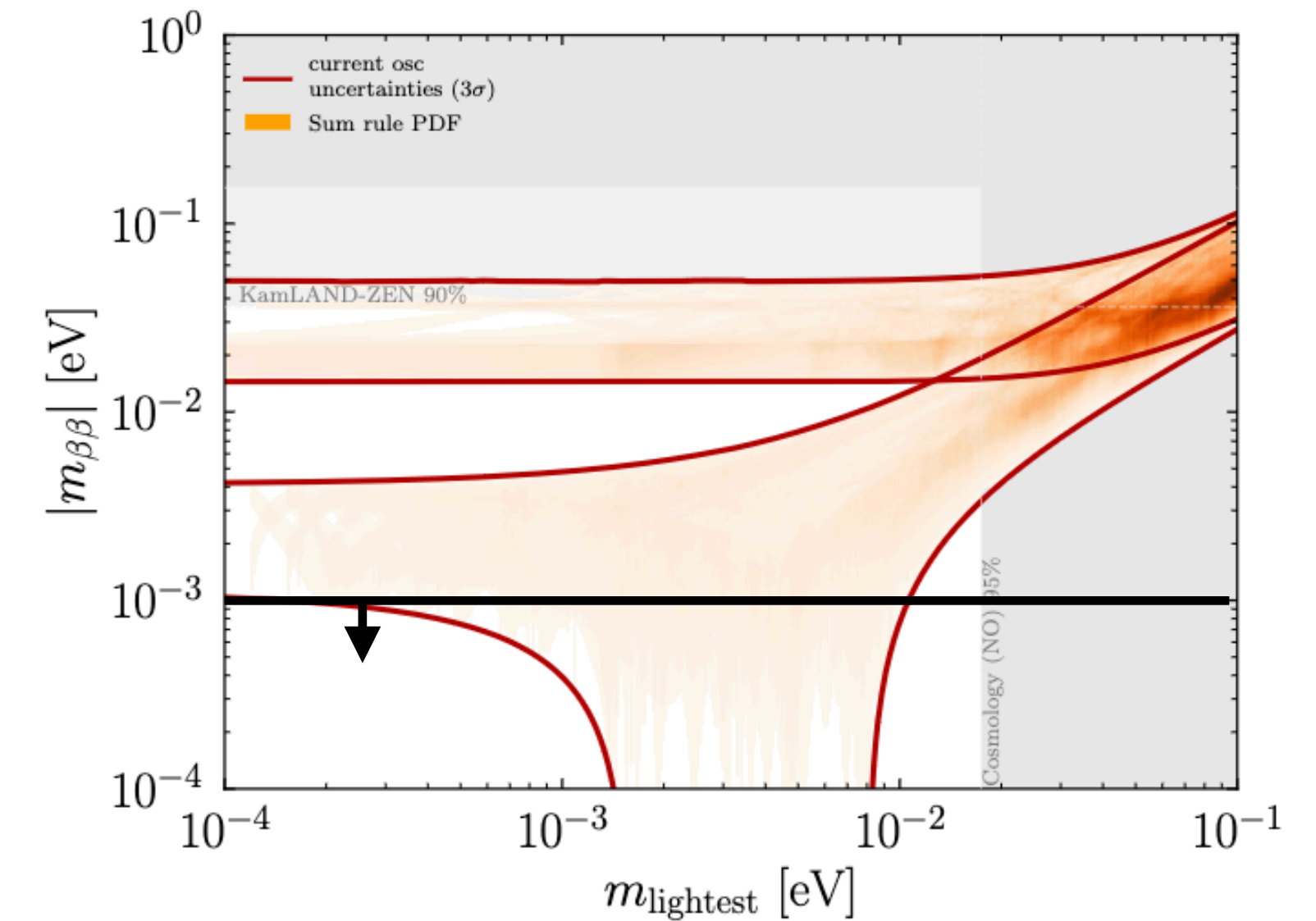
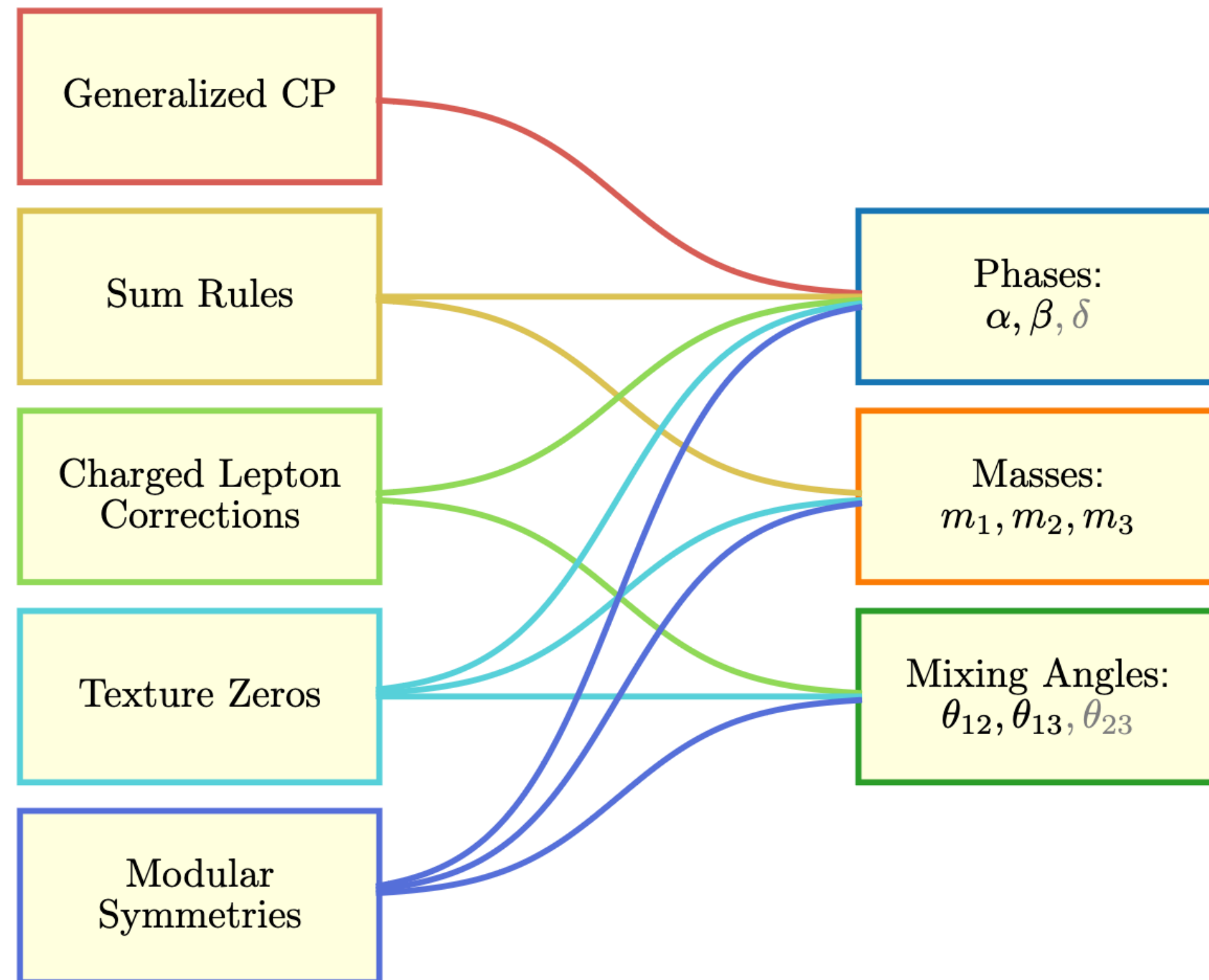
Denton, JG '23

# Neutrinoless double beta decay

## Predictions from flavor models

Denton, JG '23

**Extensive** survey of five broad categories of flavor models  
(>3000 different models)



**non-negligible fraction** of flavor models (14-100%)  
are at least partially in the funnel region  
→ interesting region to probe

# Neutrinoless double beta decay

## Sterile neutrinos

Vanilla scenario:  $0\nu\beta\beta$  due to light Majorana neutrino exchange

Additional neutrino generations can affect  $0\nu\beta\beta$  phenomenology

Phenomenology depends on ratio of sterile neutrino mass  $m_N$  to momentum transfer of process  $\langle p^2 \rangle \sim (100 \text{ MeV})^2$



electron  
neutrino



muon  
neutrino



tau  
neutrino



sterile  
neutrino



# Neutrinoless double beta decay

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$$A \propto \sum_i^{\text{light}} m_i U_{ei}^2 M^{0\nu\beta\beta}(m_i) + \sum_I^{\text{light}} m_I U_{eI}^2 M^{0\nu\beta\beta}(m_I) + \sum_I^{\text{heavy}} m_I U_{eI}^2 M^{0\nu\beta\beta}(m_I)$$

Light active neutrinos

Light sterile neutrinos

Heavy sterile neutrinos

Blennow, Fernandez-Martinez,  
Lopez-Pavon, Menendez '10

# Neutrinoless double beta decay

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$m_N < 100 \text{ MeV}$ : sterile neutrino acts like active neutrinos,  $A$  suppressed as

$$\sum_i^{\text{light}} m_i U_{ei}^2 + \sum_I^{\text{light}} m_I U_{eI}^2 = 0$$

Blennow, Fernandez-Martinez,  
Lopez-Pavon, Menendez '10

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$m_N \gg 100 \text{ MeV}$ : sterile neutrinos heavy and integrated out, **amplitude is 3-flavor amplitude**

Blennow, Fernandez-Martinez,  
Lopez-Pavon, Menendez '10

# Neutrinoless double beta decay

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$m_{N_1} > 100 \text{ MeV}$ ,  $m_{N_2} < 100 \text{ MeV}$ : some sterile neutrinos are heavy, some are light  
→ **cancellation** of light sterile amplitude **with SM amplitude prevented**

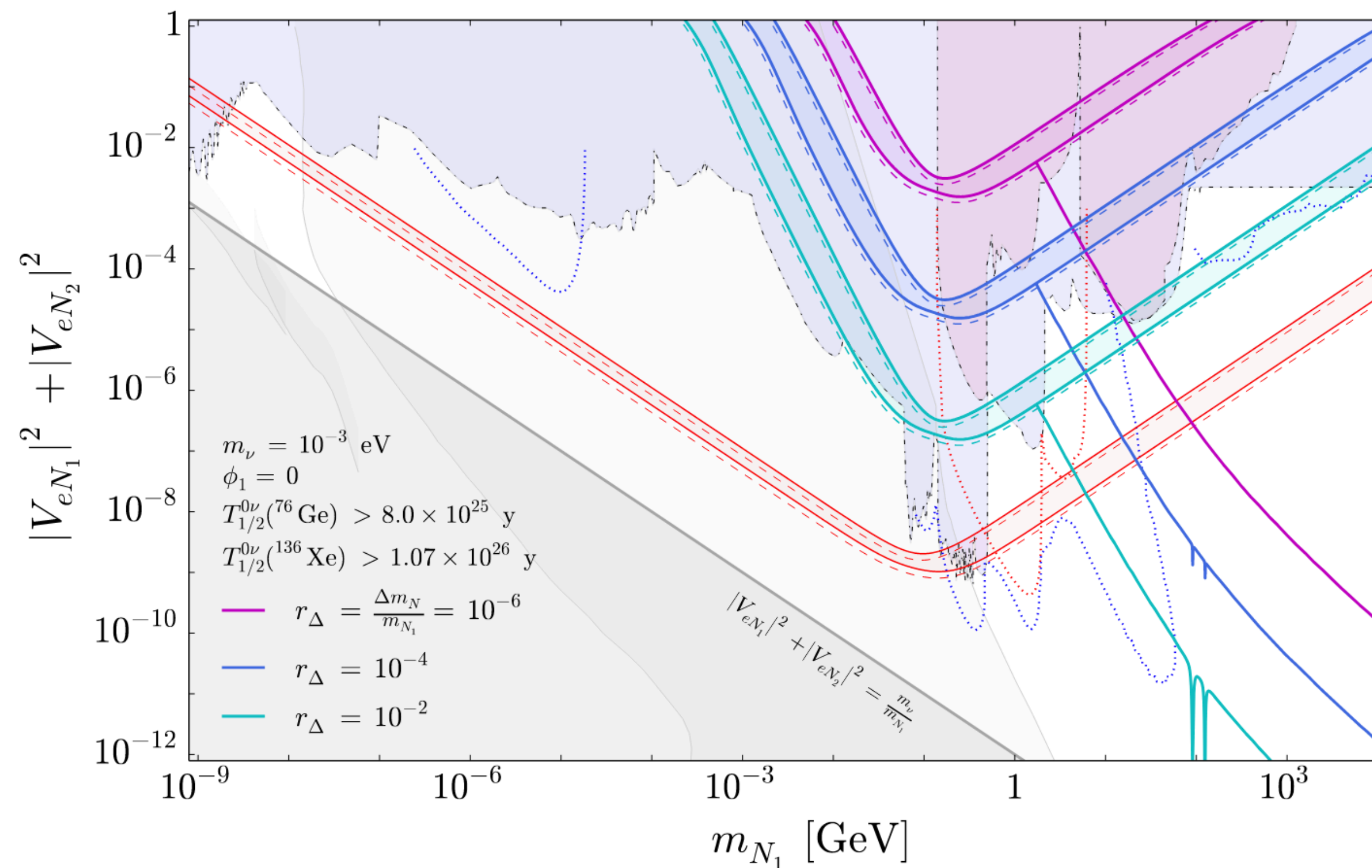
Blennow, Fernandez-Martinez,  
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# Neutrinoless double beta decay

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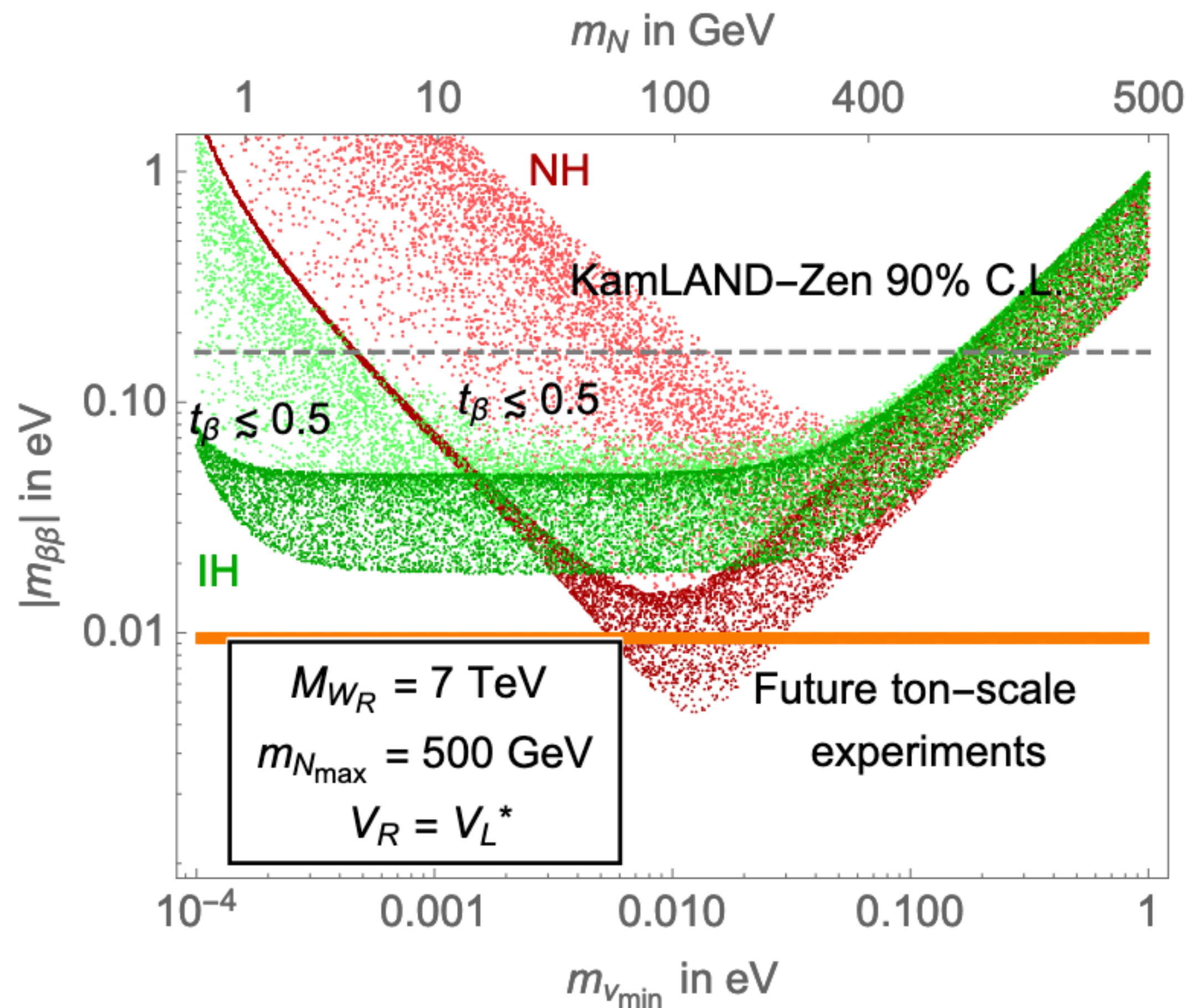
→ **Constraints** on sterile neutrinos



Bolton, Dev, Deppisch '19

# Neutrinoless double beta decay

## Minimal left-right symmetric model



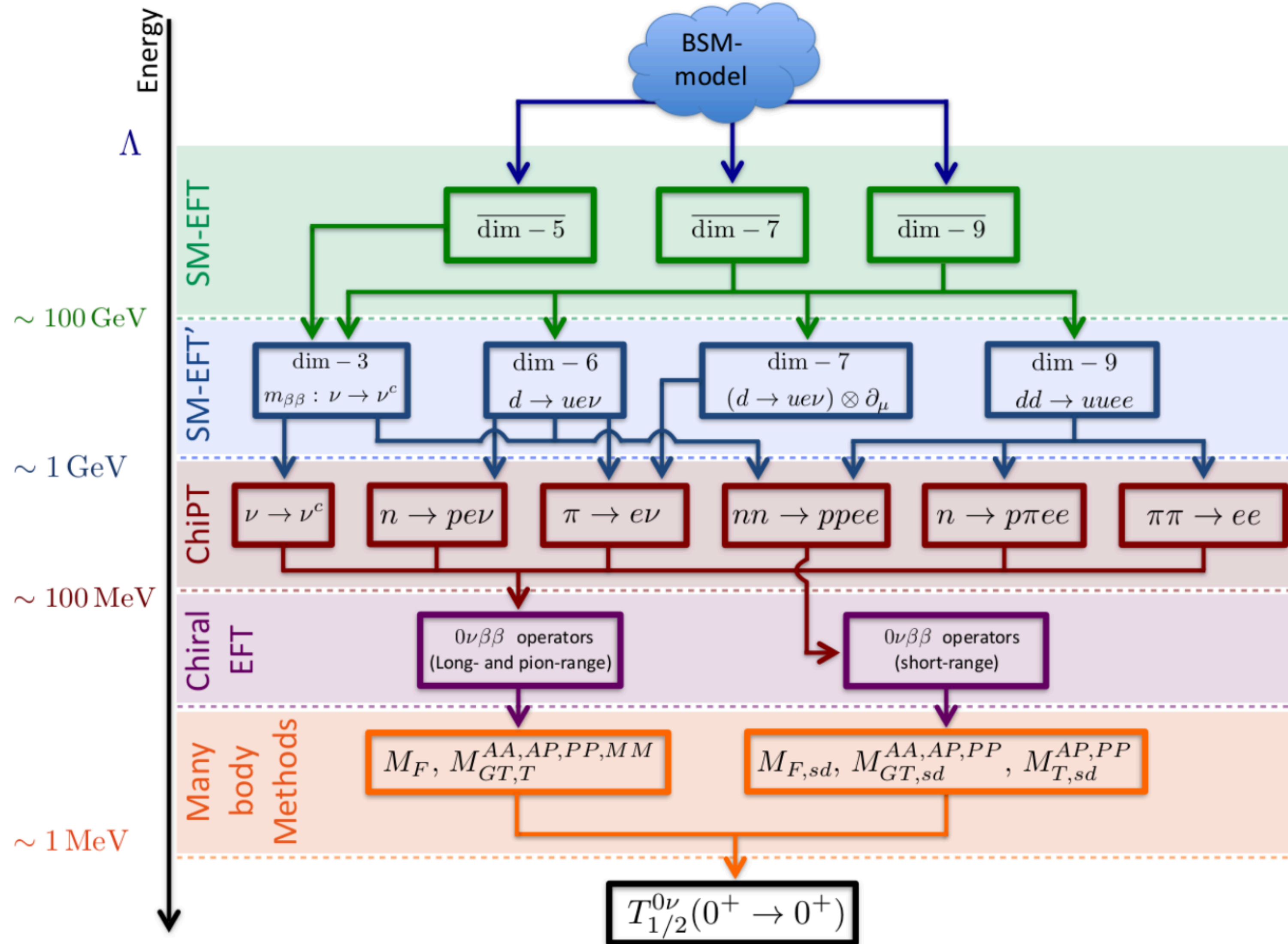
New mediators!

Particle physics quantity is not  
 $|m_{\beta\beta}| = \left| \sum U_{ei}^2 m_i \right|$  anymore

Li, Ramsey-Musolf, Vasquez '20

# Neutrinoless double beta decay

Cirigliano, JG et al '22



Nuclear matrix elements **affected** by new physics

Physics spans a large range of energies

→ need a tower of EFTs to go from high-energy model to the nuclear matrix element

# Neutrinoless double beta decay

## Summary and conclusions

Neutrinoless double beta decay allows to probe lepton number violation and can provide insights into matter-antimatter asymmetry generation, and test symmetries of SM

- Current  $0\nu\beta\beta$  experiments are **ongoing**
- **New** experimental collaborations are forming and will continue in future
- Need to define theoretical goals/targets for these experiments to provide benchmarks
- Sensitivity studies needed for **new physics scenarios** affecting  $0\nu\beta\beta$
- For correct interpretation of results: **theory work on nuclear matrix elements** required  
with **robust** uncertainty quantification

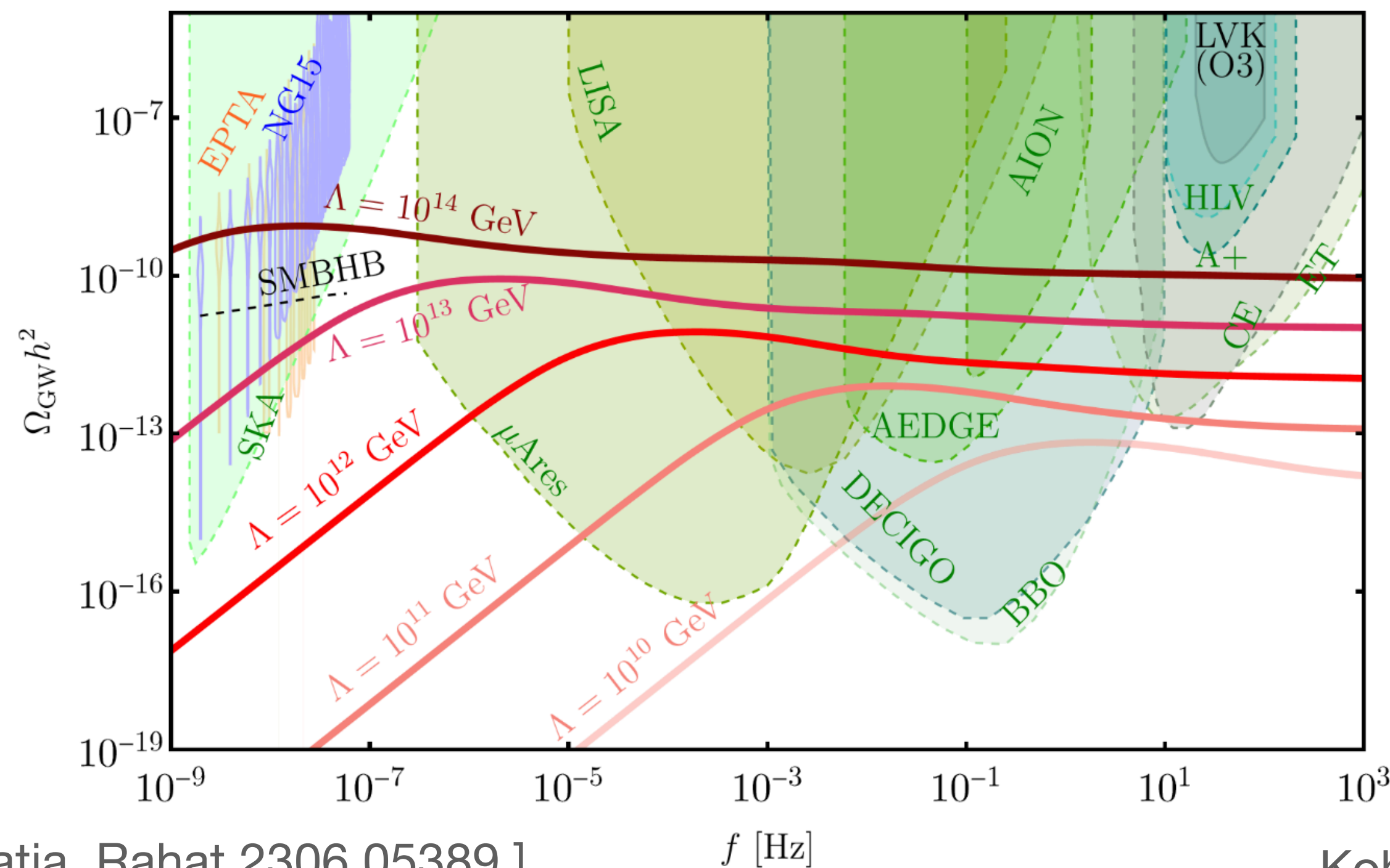


# Thanks for your attention!



# Appendix: Lepton number violation

Signs of **lepton number breaking** in the early Universe  
**Gravitational waves** from decay of cosmic strings from breaking of lepton number symmetry



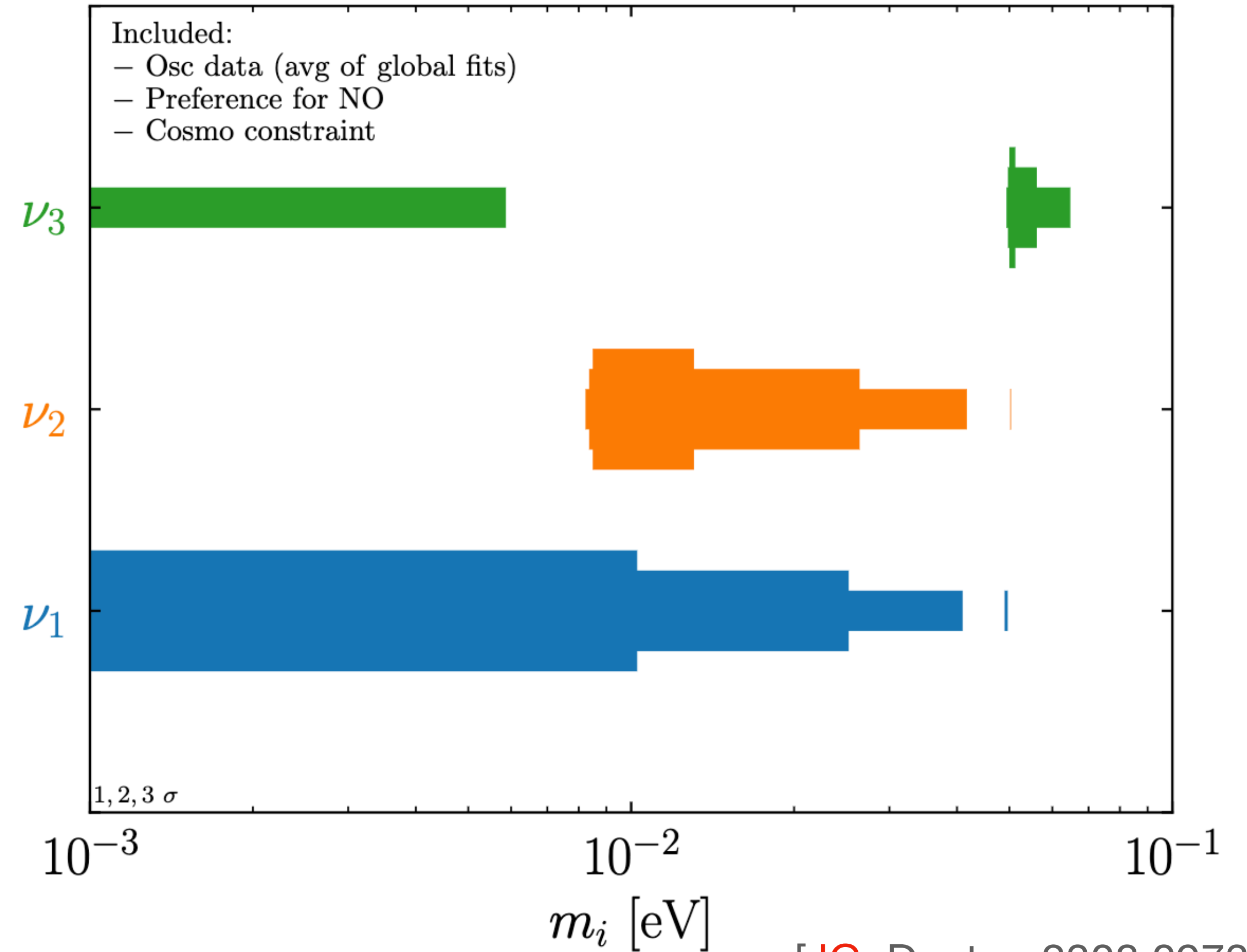
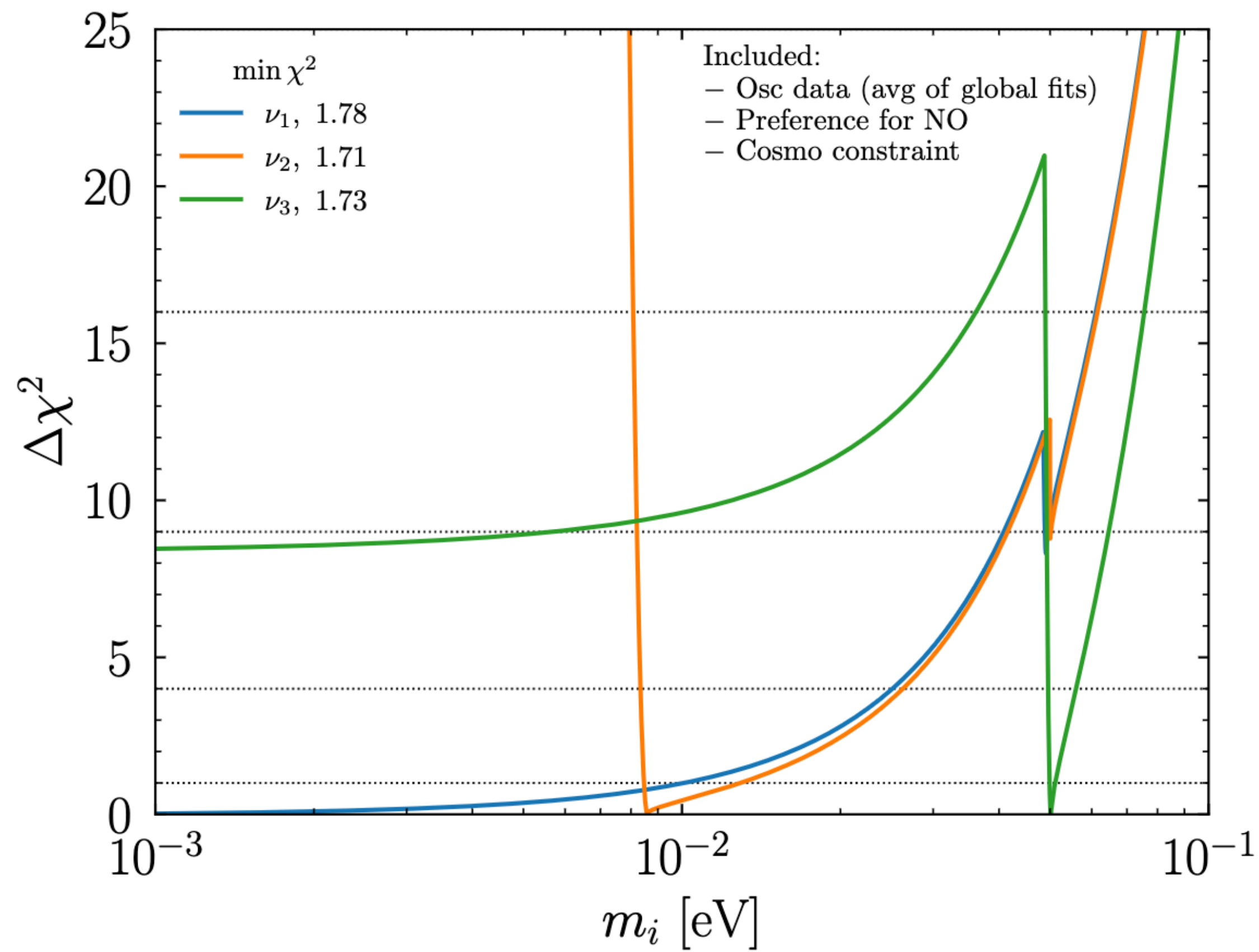
Signature depends on

**Lepton number  
breaking scale**

[King, Marfatia, Rahat [2306.05389](#) ]

[see also Dror, Hiramatsu,  
Kohri, Murayama, White [2306.05389](#) ]

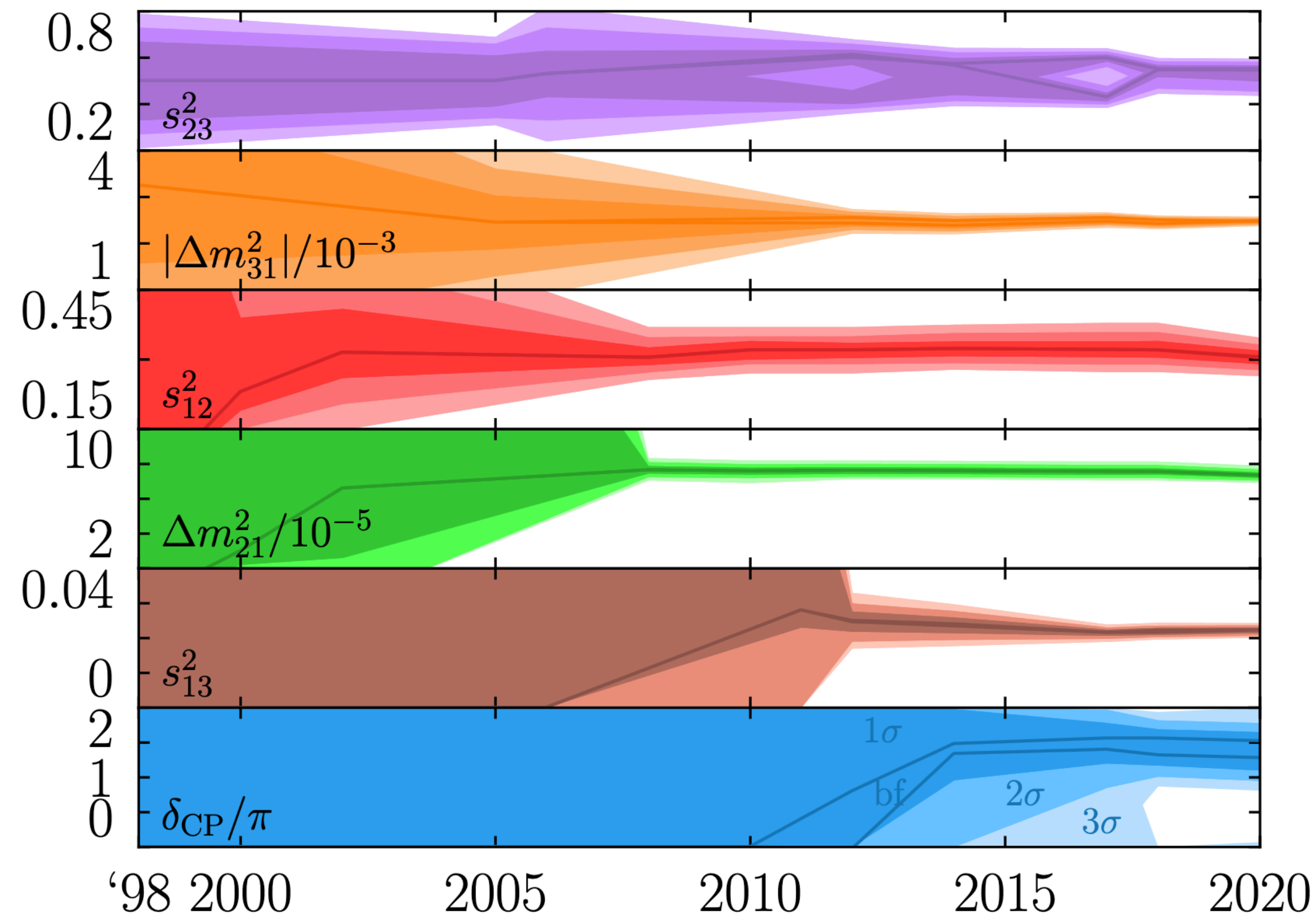
# Appendix: Neutrino mass



[JG, Denton 2308.09737]

# Appendix: Neutrino oscillation parameters

Neutrino oscillation parameters measured over years



[Denton et al [2212.00809](#)]

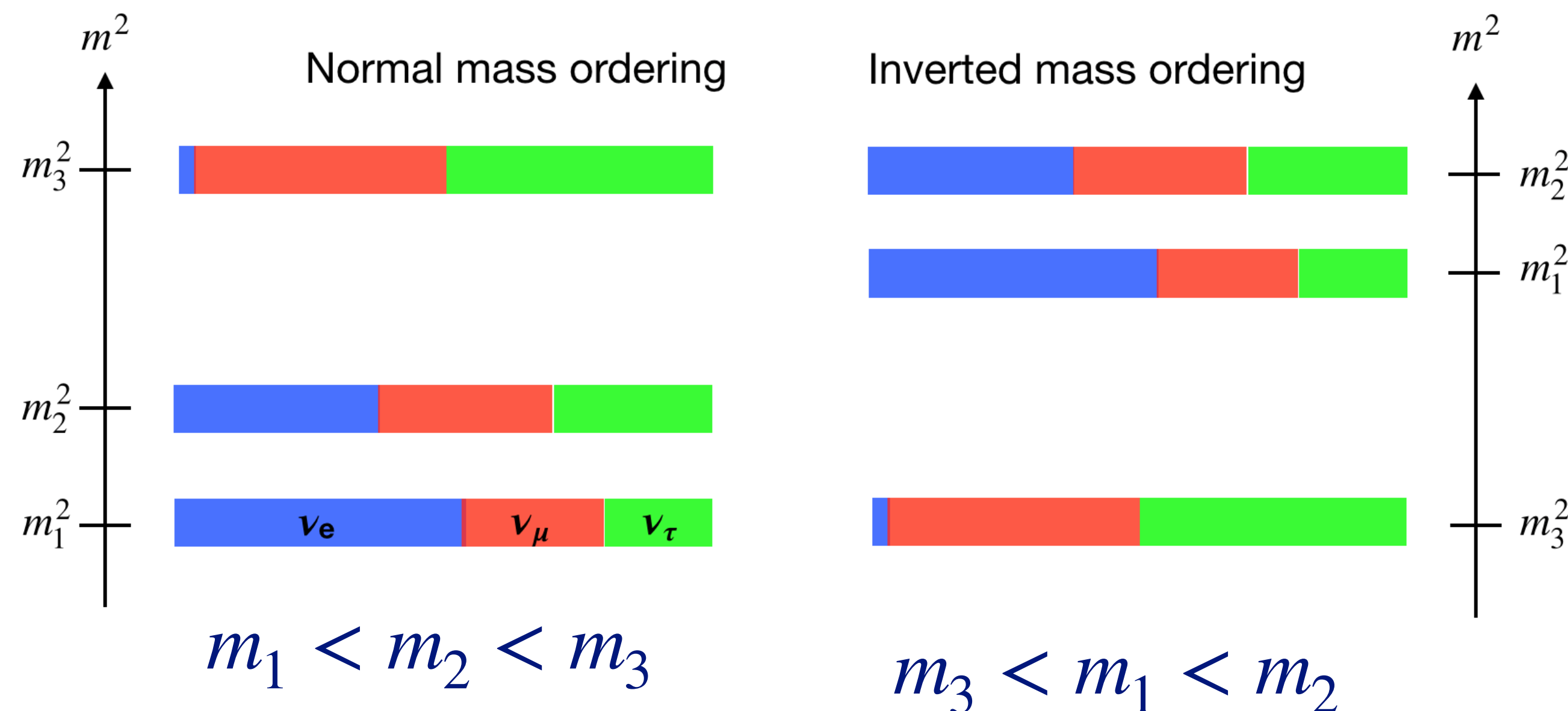
# Appendix: Neutrino oscillation parameters

Global fits to oscillation data:

mass splittings:  $|\Delta m_{32}^2| = 2.5 \cdot 10^{-3} \text{ eV}^2$ ,  $\Delta m_{21}^2 = 7.4 \cdot 10^{-5} \text{ eV}^2$

[nufit v5.1]

mass ordering **unknown**



# Appendix: numerical approach

[JG, Denton [2308.09737](#)]

1. We first calculate the number of models which are viable. These are the models that are in agreement with the oscillation data.
2. Then we determine which of those have any fraction within the funnel which we define to be  $m_{\beta\beta} < 10^{-3}$  eV.
3. Then we determine the fraction of each model that is within the funnel as outlined below.

$$f = \frac{\int_{\text{funnel}} d \log m_{\text{lightest}} d \log m_{\beta\beta}}{\int d \log m_{\text{lightest}} d \log m_{\beta\beta}}$$

# Appendix: Texture zeros

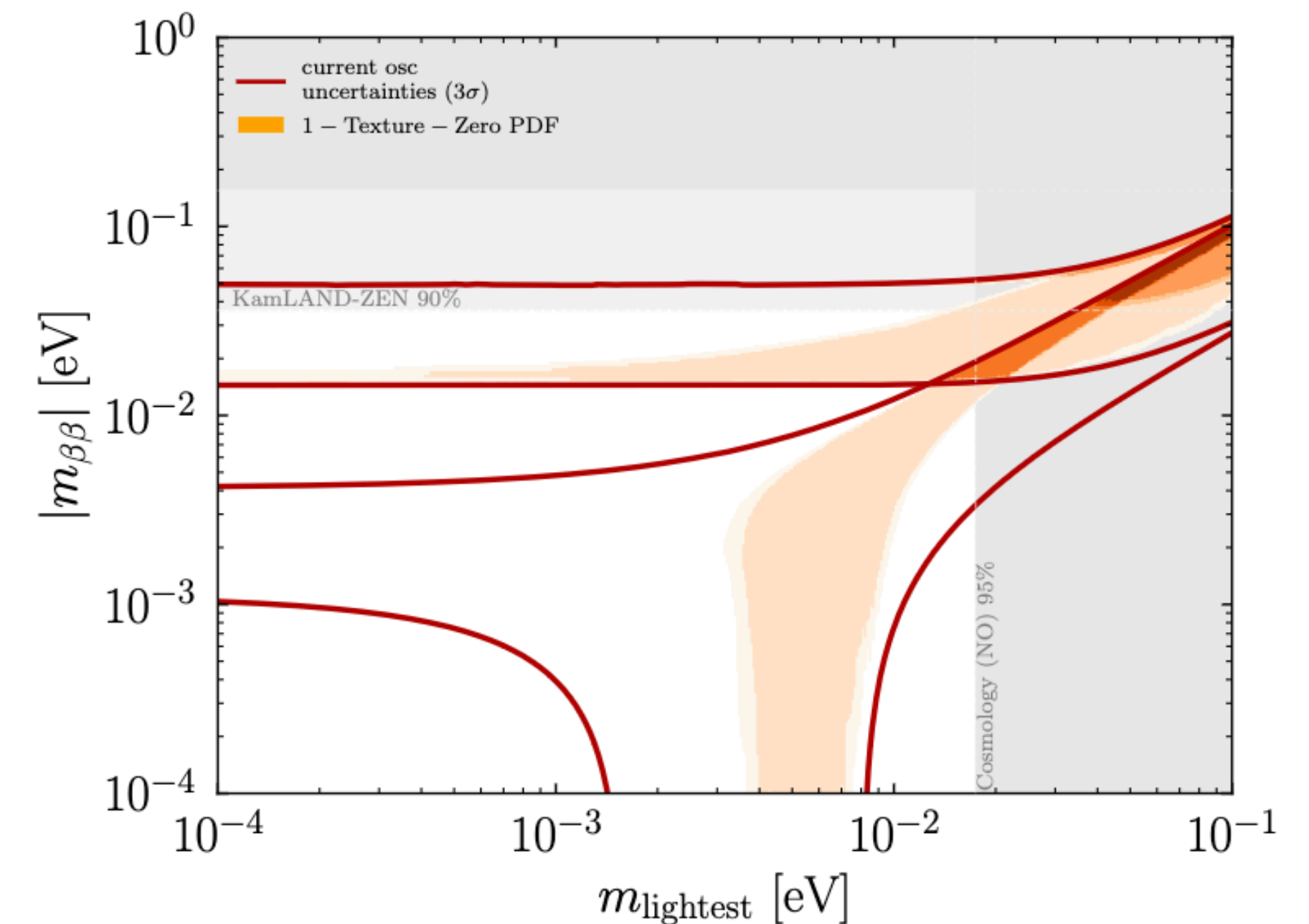
Assume symmetric Majorana mass matrix has vanishing entries

[JG, Denton 2308.09737]

1-1 elements is  $|m_{\beta\beta}|$

All 6 possible one-texture zero mass matrices in **agreement** with data

	Fraction in funnel
$M_{ee}$	1
$M_{e\mu}$	0.31
$M_{e\tau}$	0.30
$M_{\mu\mu}$	0
$M_{\mu\tau}$	0
$M_{\tau\tau}$	0



# Appendix: Texture zeros

Assume symmetric Majorana mass matrix has vanishing entries

[JG, Denton 2308.09737]

1-1 elements is  $|m_{\beta\beta}|$

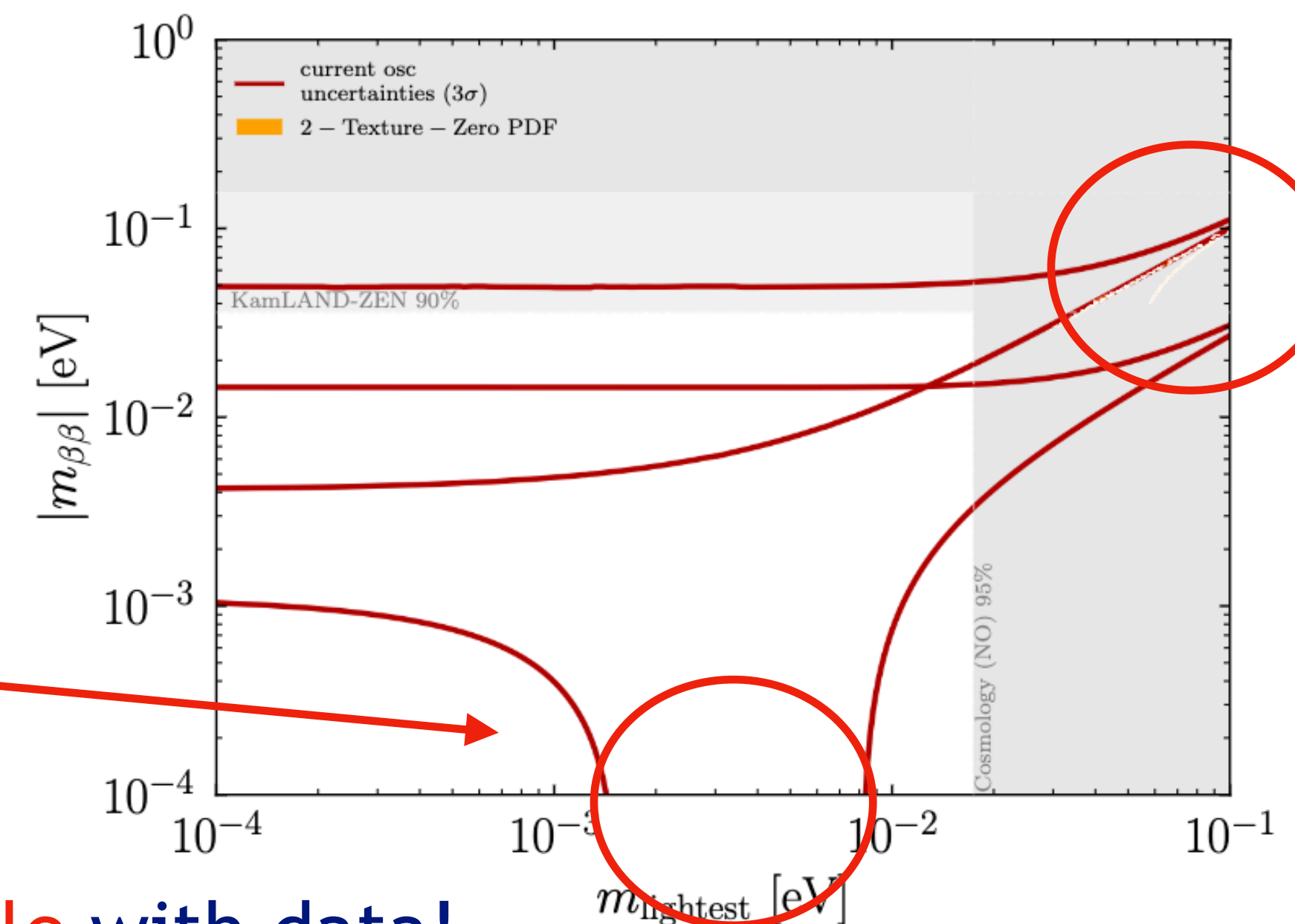
7 of 15 possible two-texture zero mass matrices in agreement with data

New result!

Models fully in funnel

Constrained by cosmology

	$M_{e\mu}$	$M_{e\tau}$	$M_{\mu\mu}$	$M_{\mu\tau}$	$M_{\tau\tau}$
$M_{ee}$	1	1	X	X	X
$M_{e\mu}$		X	0	X	0
$M_{e\tau}$			0	X	0
$M_{\mu\mu}$				X	0
$M_{\mu\tau}$					X



Models with 3+ texture zeros **not compatible** with data!



# Appendix: Mass sum rules

$$c_1 e^{i\chi_1} (m_1 e^{i\alpha})^d + c_2 e^{i\chi_2} (m_2 e^{i\beta})^d + m_3^d = 0$$

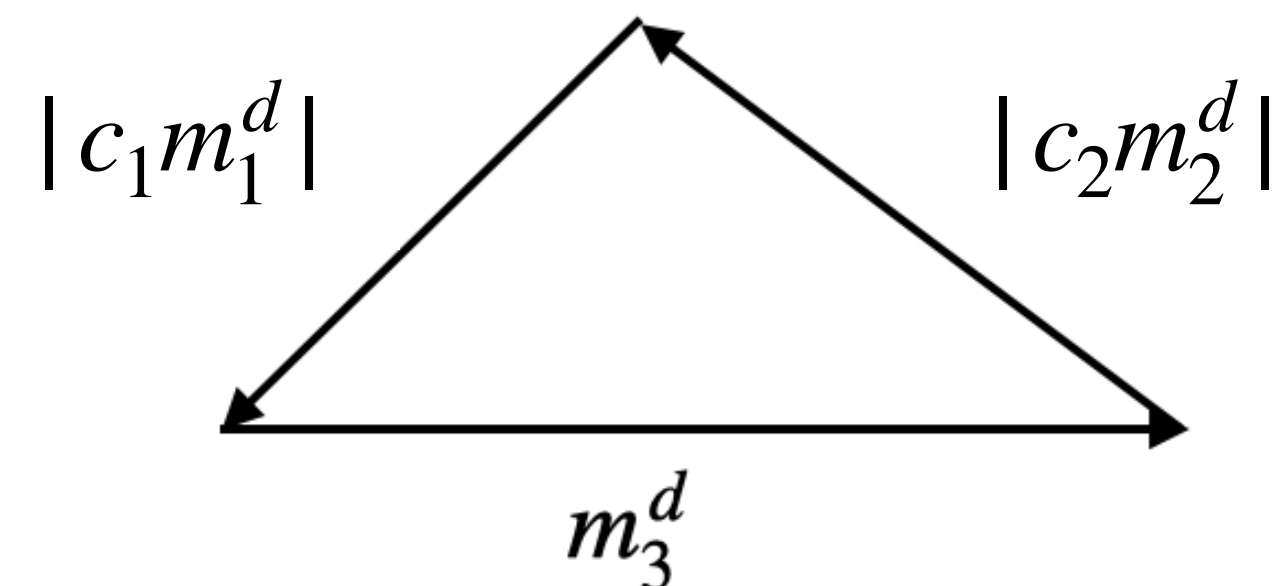
[S. King, A. Merle, A. Stuart '13  
J. Barry, W. Rodejohann '10 ]

12 different SR in over 60 models realized in literature

$$c_i \sim \mathcal{O}(1), \chi_i = (0, \pi, \pm \pi/2), d = (1, -1, \pm 1/2),$$

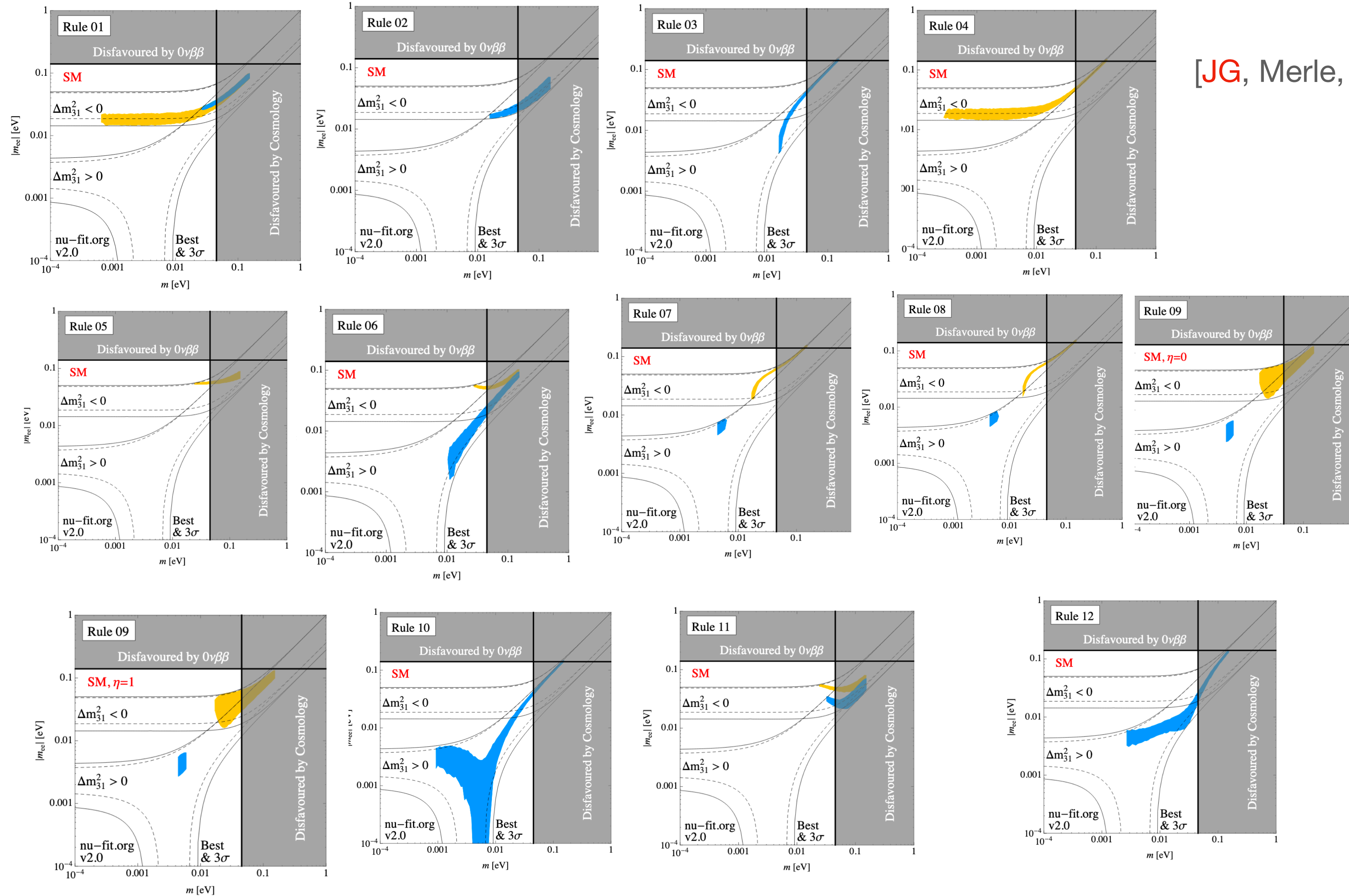
constant and fixed by model

parametrized as triangle in complex plane



# Appendix: Mass sum rules

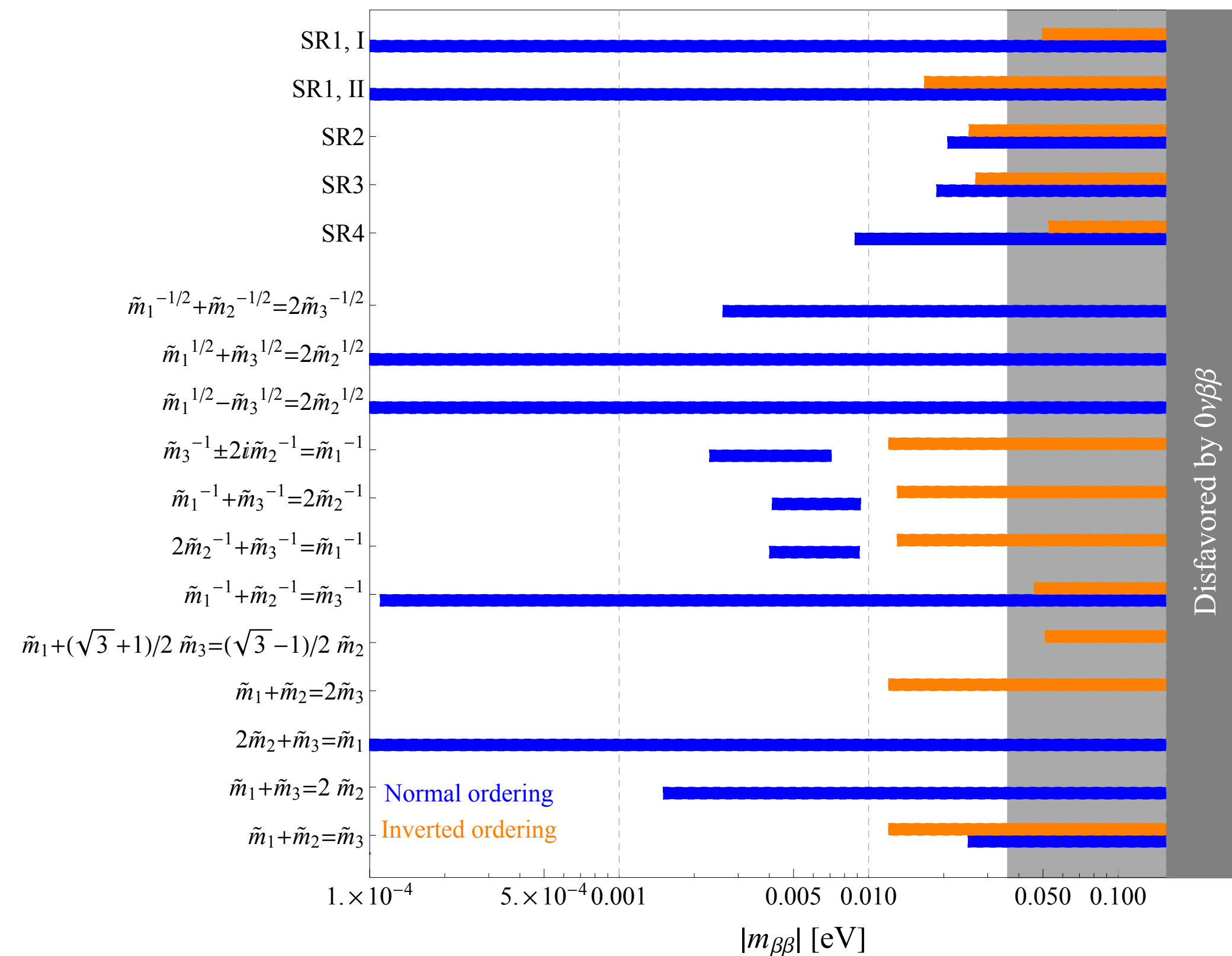
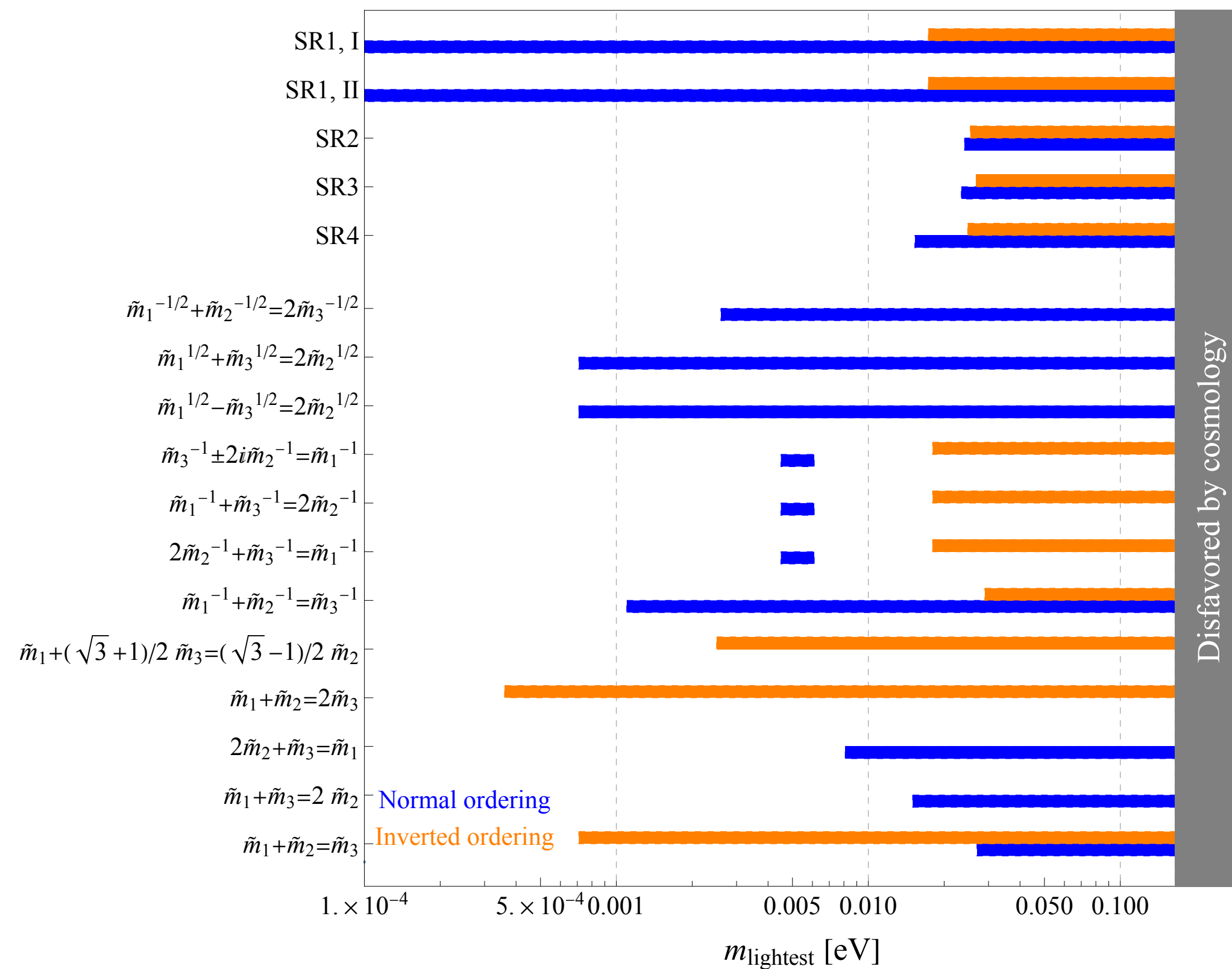
[JG, Merle, Spinrath [1506.06139](#)]



# Appendix: Mass sum rules

Predictions for upcoming experiments

Can be used to plan stages of experiments like in [Merle, Agostini, Zuber 1506.06133]

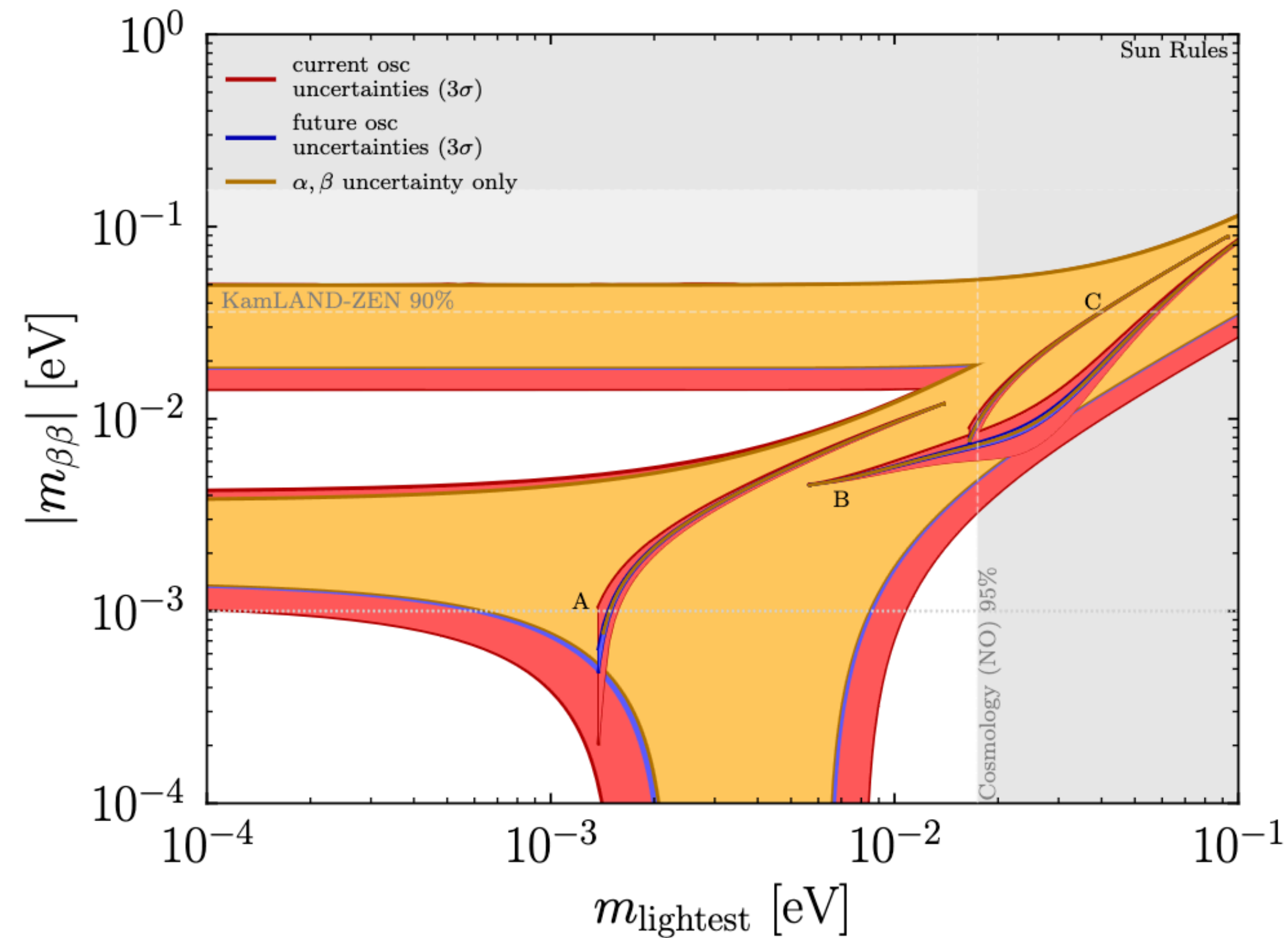


[JG, et al 2203.12169]

# Appendix: Mass sum rules

$$c_1 e^{i\chi_1} (m_1 e^{i\alpha})^d + c_2 e^{i\chi_2} (m_2 e^{i\beta})^d + m_3^d = 0$$

[JG, Denton 2308.09737]



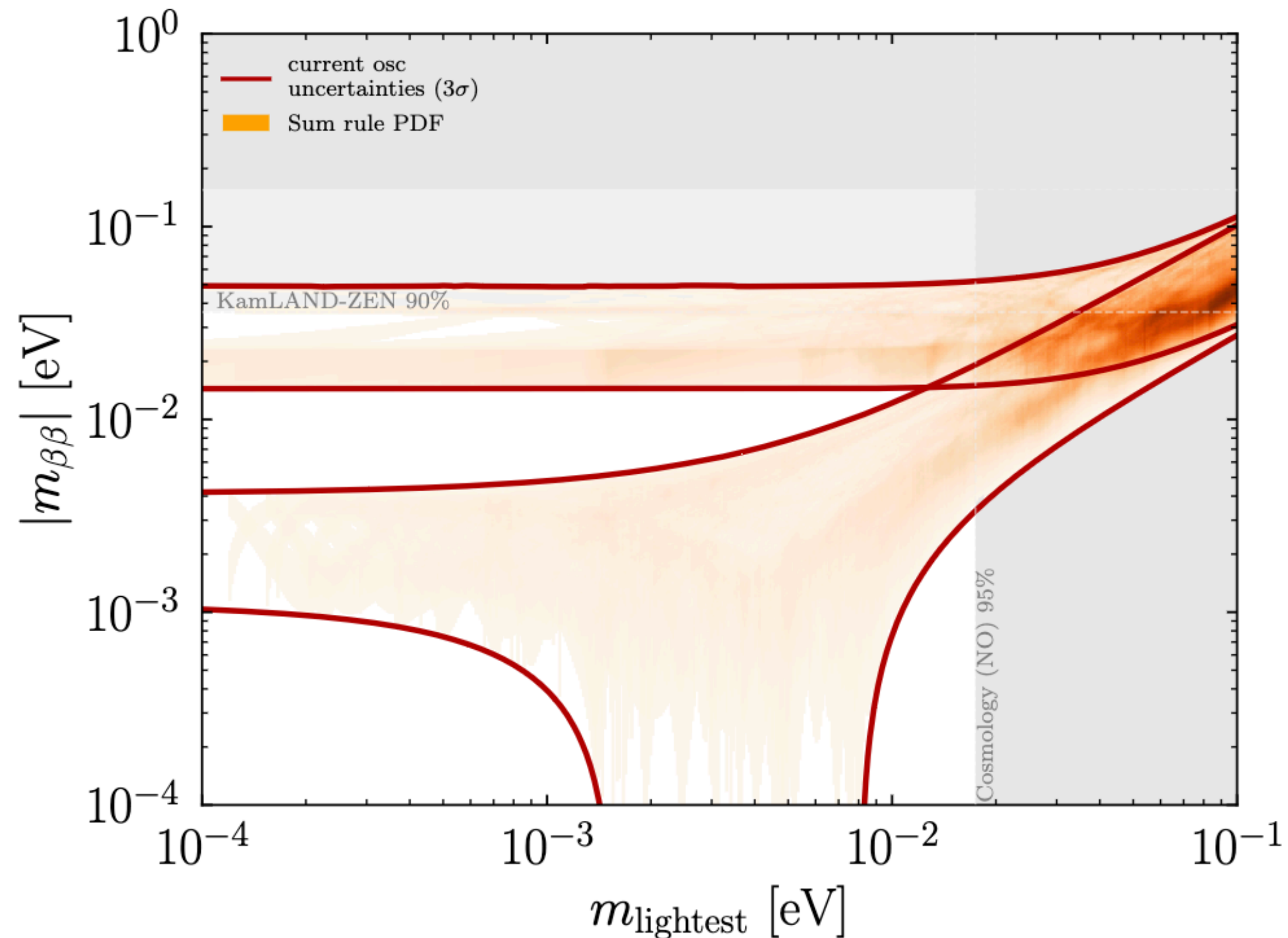
$(c_1, c_2, d, \chi_1, \chi_2) : A : (1, 2, 1/2, \pi, \pi/2), B : (1/2, 1/2, -1/2, \pi, \pi), C : (1, 2, 1, \pi, 0)$

# Appendix: Mass sum rules

3137 models tested, found 1968 viable models

[JG, Denton [2308.09737](#)]

Probability density plot



Predict large neutrino masses

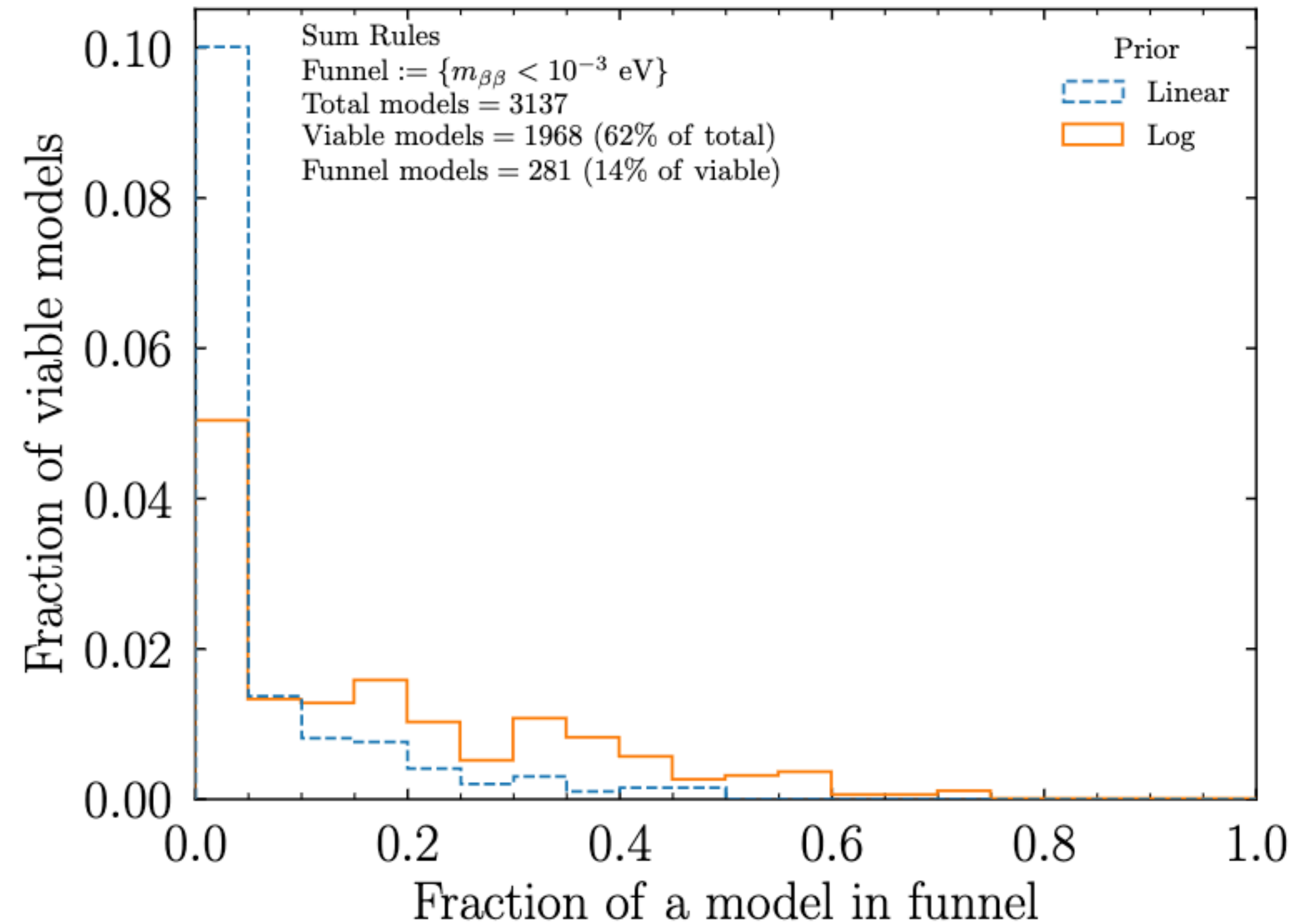
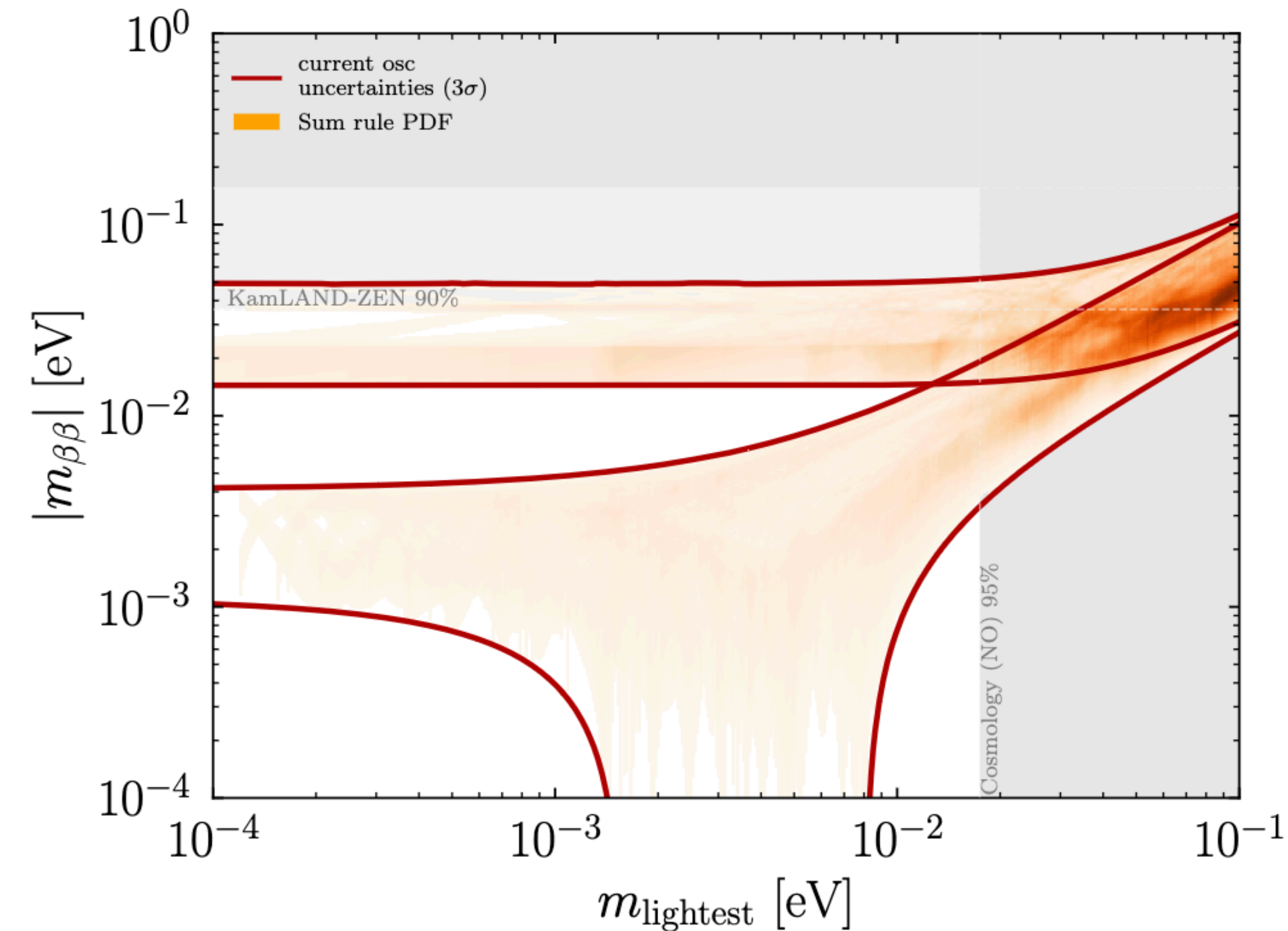
→ tested with cosmology

# Appendix: Mass sum rules

3137 models tested, found 1968 viable models

[JG, Denton [2308.09737](#)]

Probability density plot

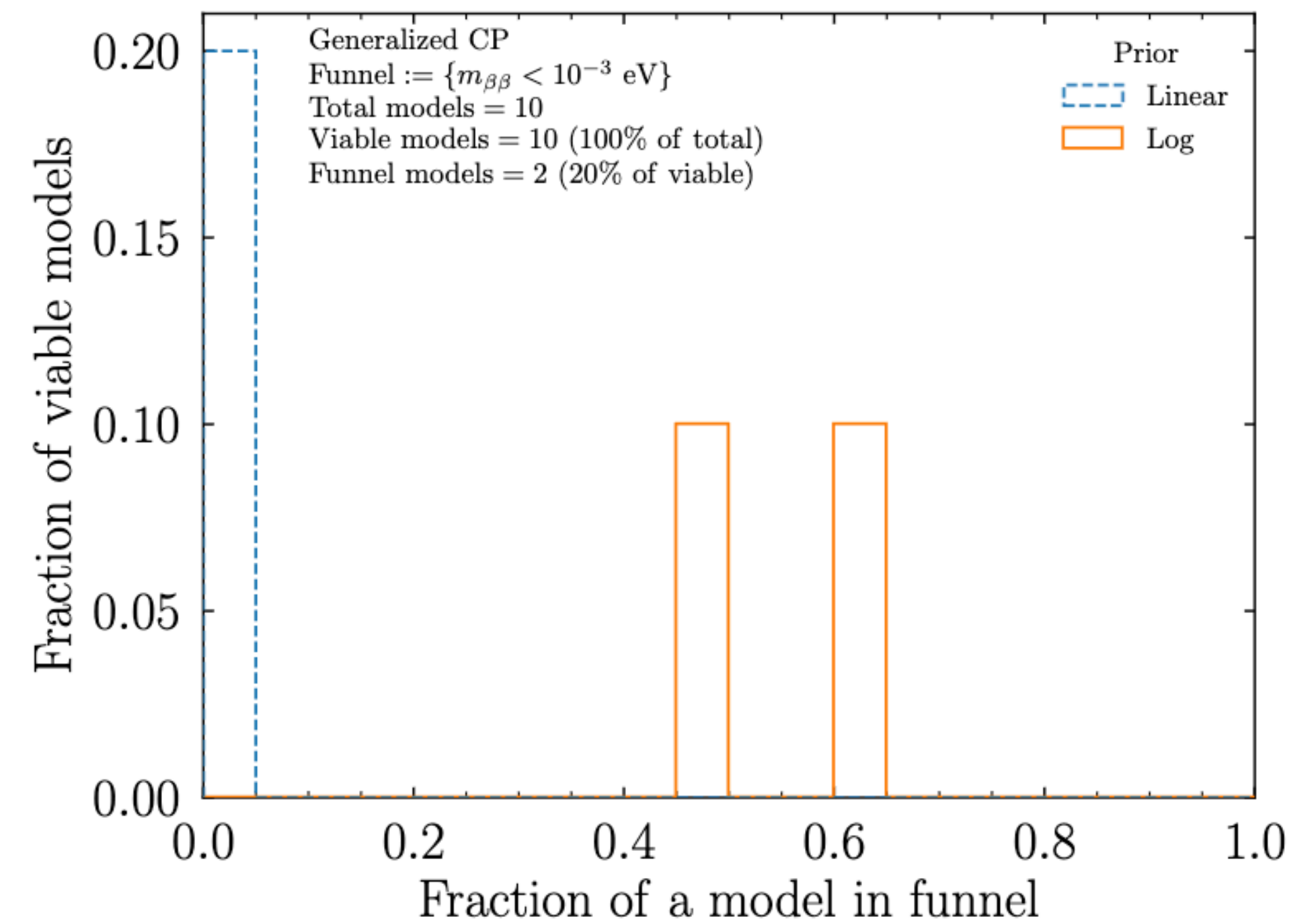
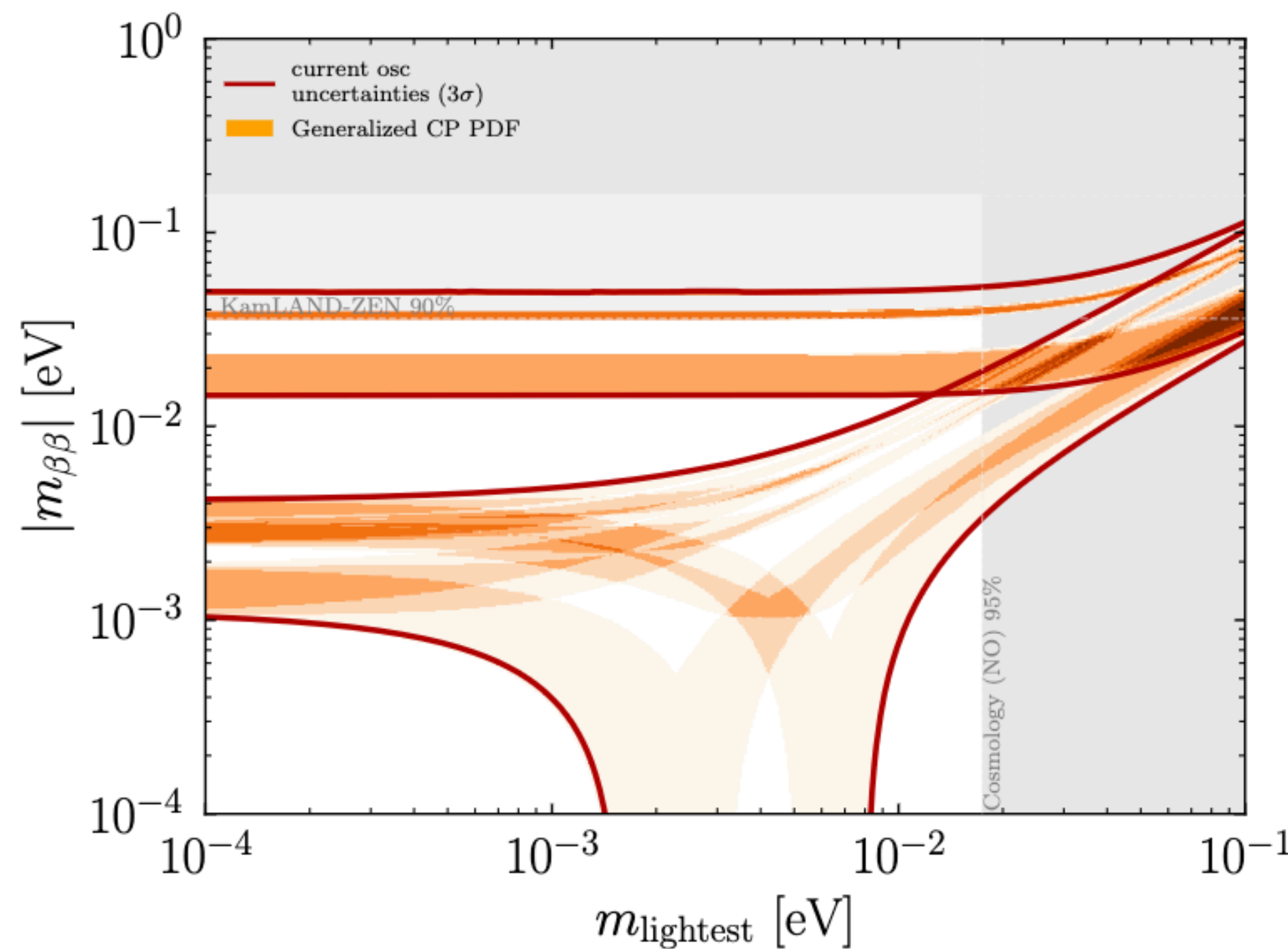


# Appendix: Results for generalized CP

Phases have specific values

$(\alpha, \beta)$	$(\alpha, \beta)$
$(\mathbf{0}, \boldsymbol{\pi})$	$(0, \pi/2)$ or $(0, 3\pi/2)$
$(\boldsymbol{\pi}, \mathbf{0})$	$(\pi/2, 3\pi/2)$ or $(3\pi/2, \pi/2)$
$(\mathbf{0}, \mathbf{0})$	$(\pi, \pi/2)$ or $(\pi, 3\pi/2)$
$(\boldsymbol{\pi}, \boldsymbol{\pi})$	$(\pi/2, 0)$ or $(3\pi/2, 0)$
	$(\pi/2, \pi/2)$ or $(3\pi/2, 3\pi/2)$
	$(\pi/2, \pi)$ or $(3\pi/2, \pi)$

[JG, Denton 2308.09737]



# Appendix: Results for charged lepton corrections

$$U_{\text{PMNS}} = U_e^\dagger U_\nu$$

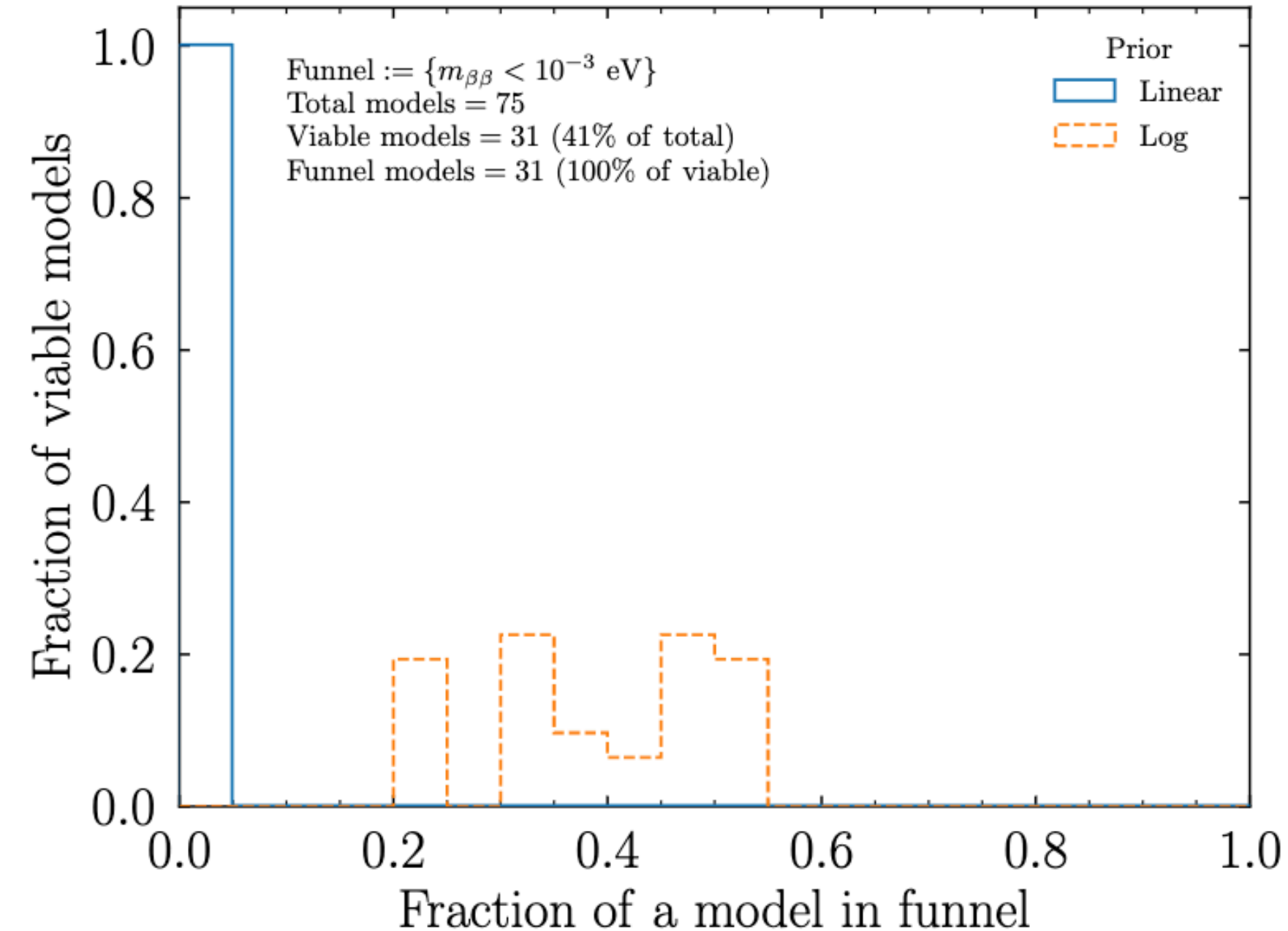
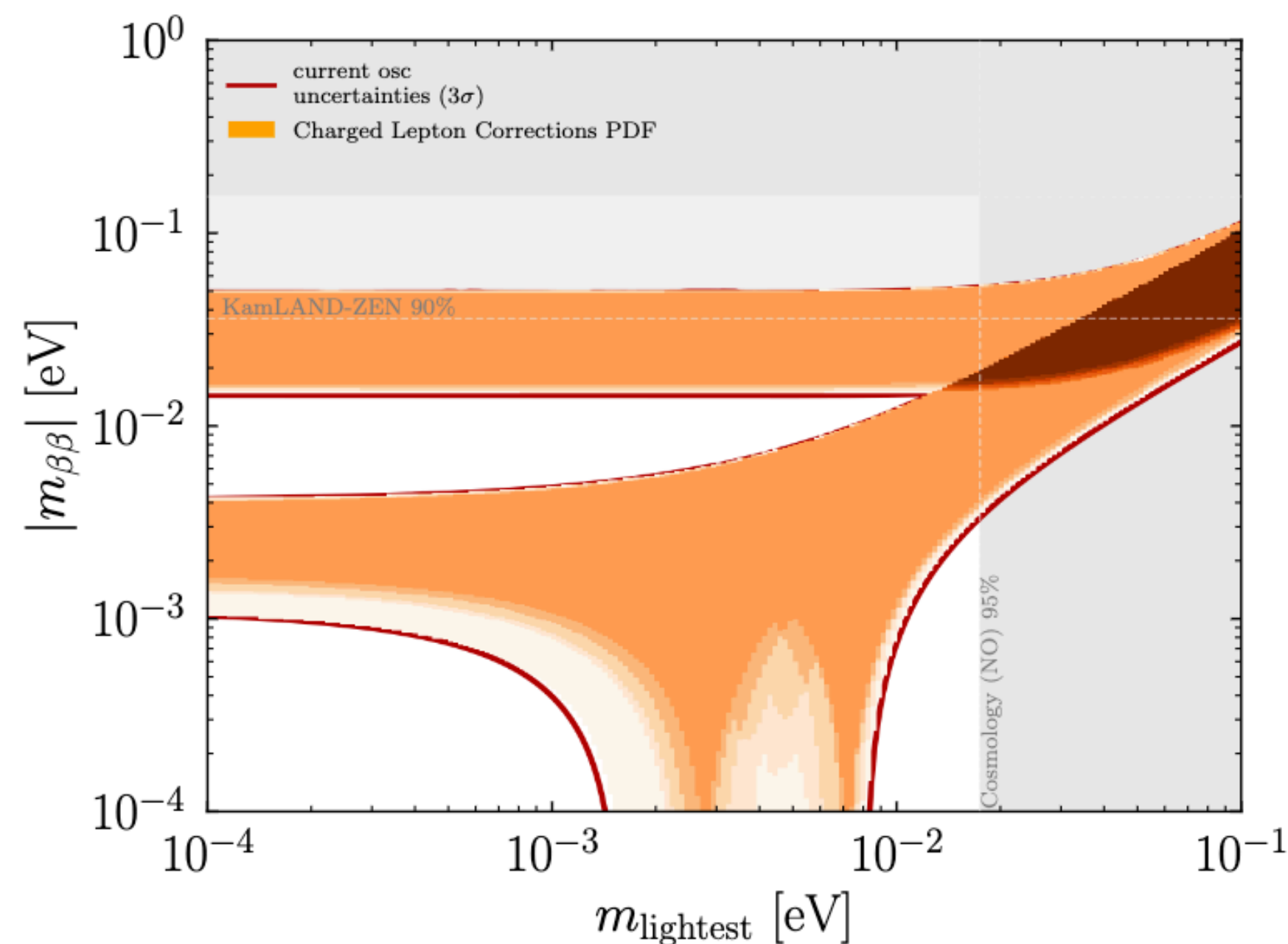
[JG, Denton [2308.09737](#)]

Angles in neutrino sector determined by underlying symmetry

Studied two rotations in the neutrino sector, one charged lepton rotation

two rotations in the neutrino sector, two charged lepton rotation

three rotations in the neutrino sector, one charged lepton rotation





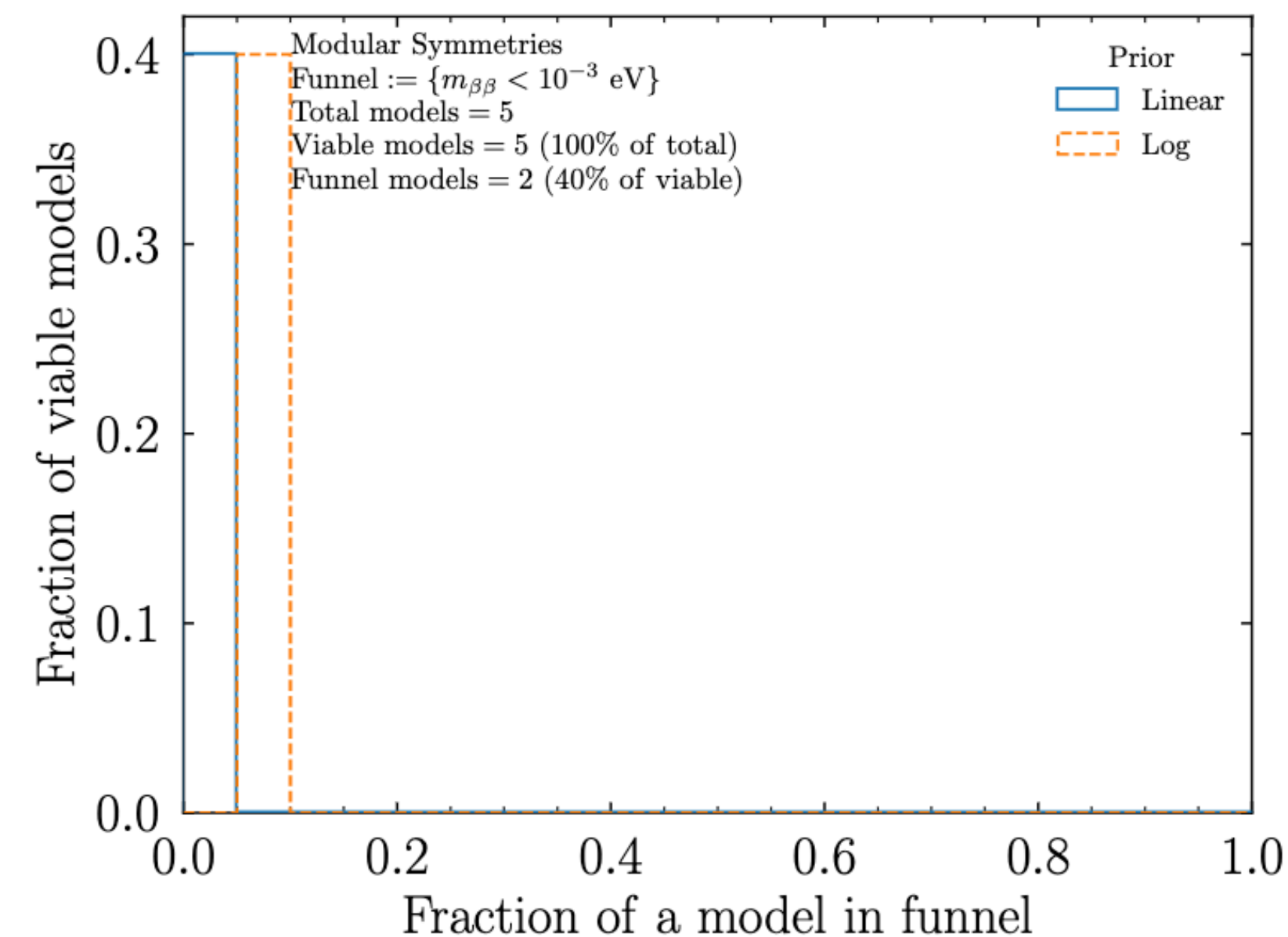
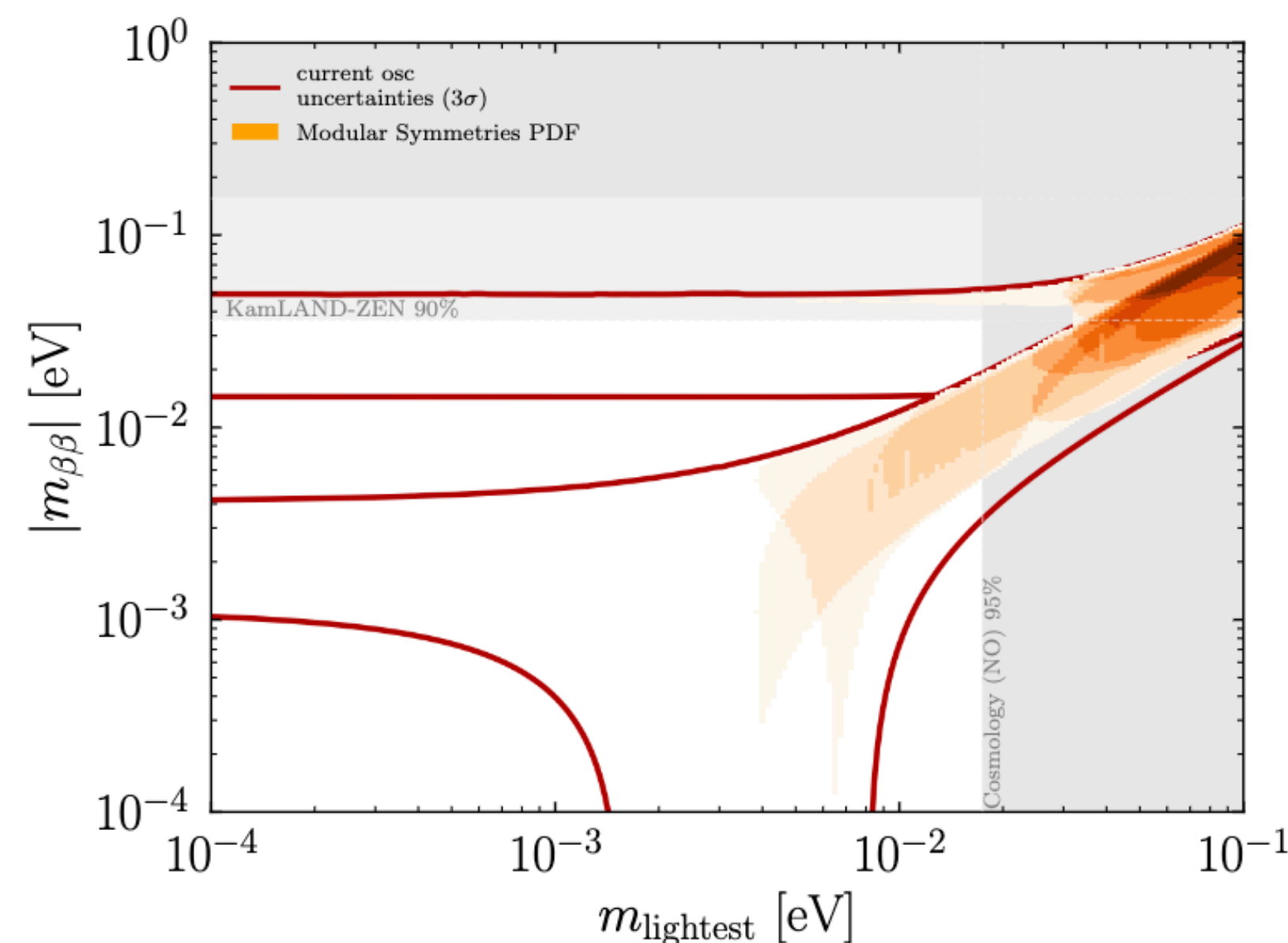
# Appendix: Results for modular symmetries

[JG, Denton [2308.09737](#)]

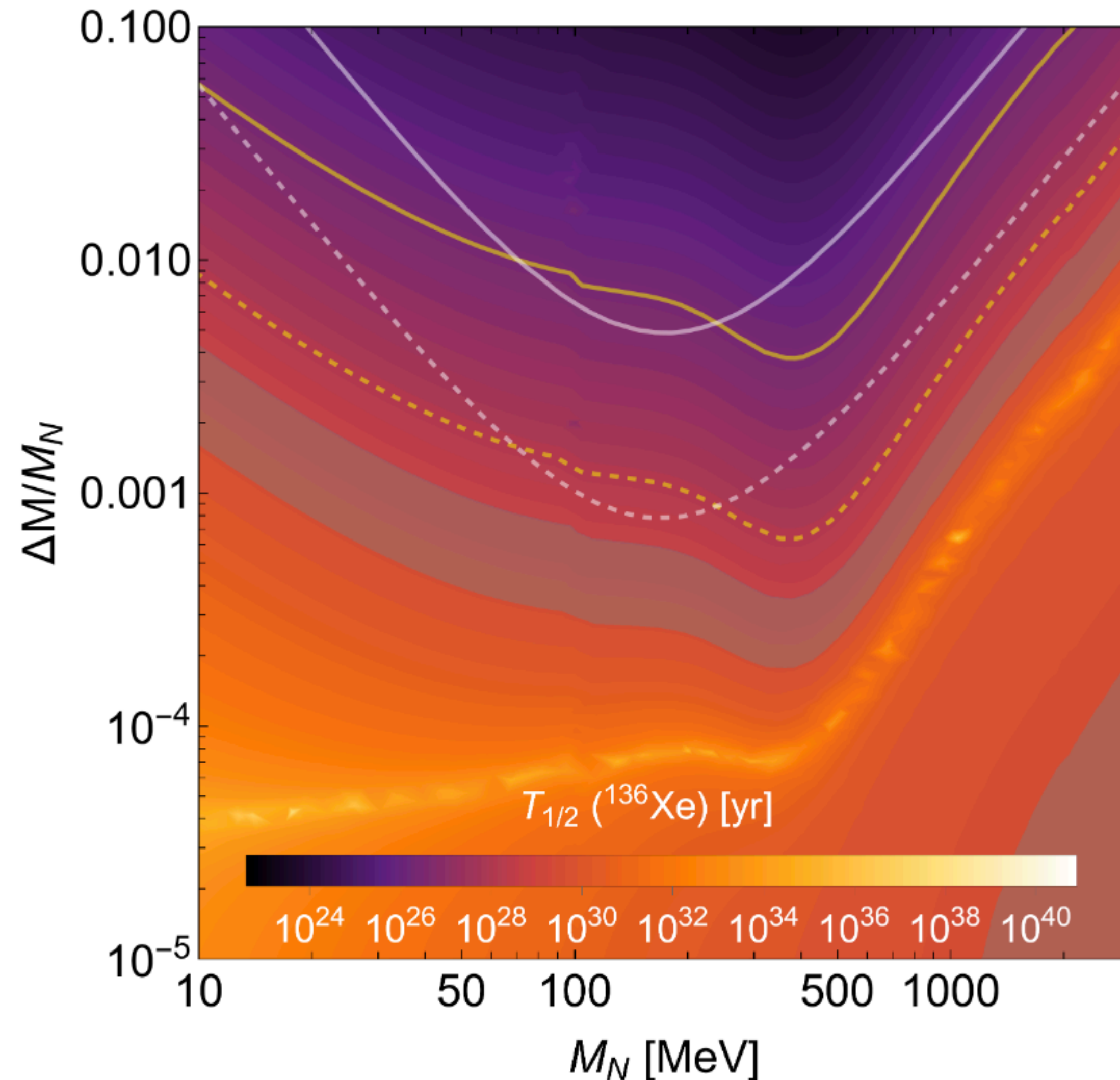
Reduced numbers of fields which break flavor symmetry [F. Feruglio '17]

5 models with maximal number of predictions realized in literature

Coefficients of sum rules depend on mixing parameters [JG, Spinrath [2012.04131](#)]



# Appendix: Sterile neutrinos in $0\nu\beta\beta$



[Dekens et al '24]

3+2 scenario