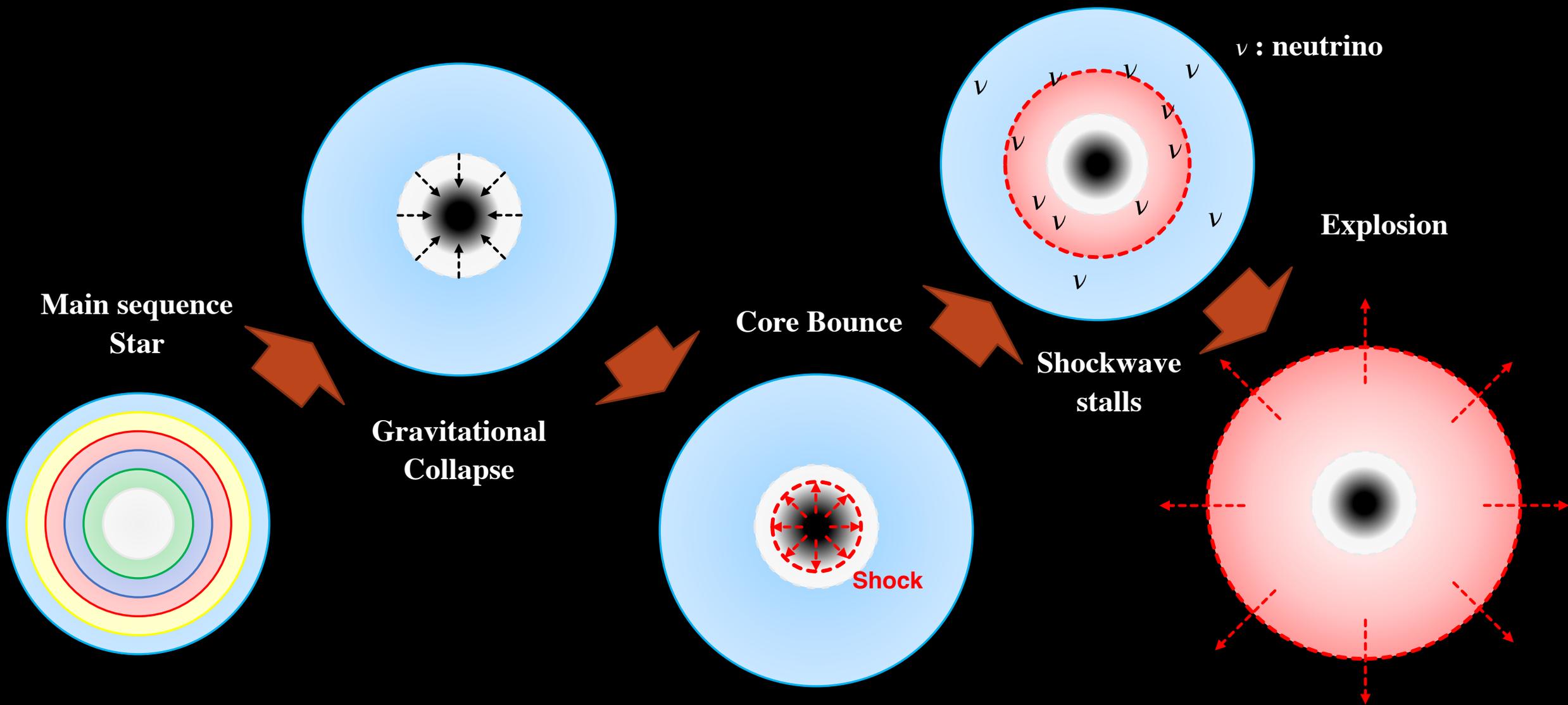


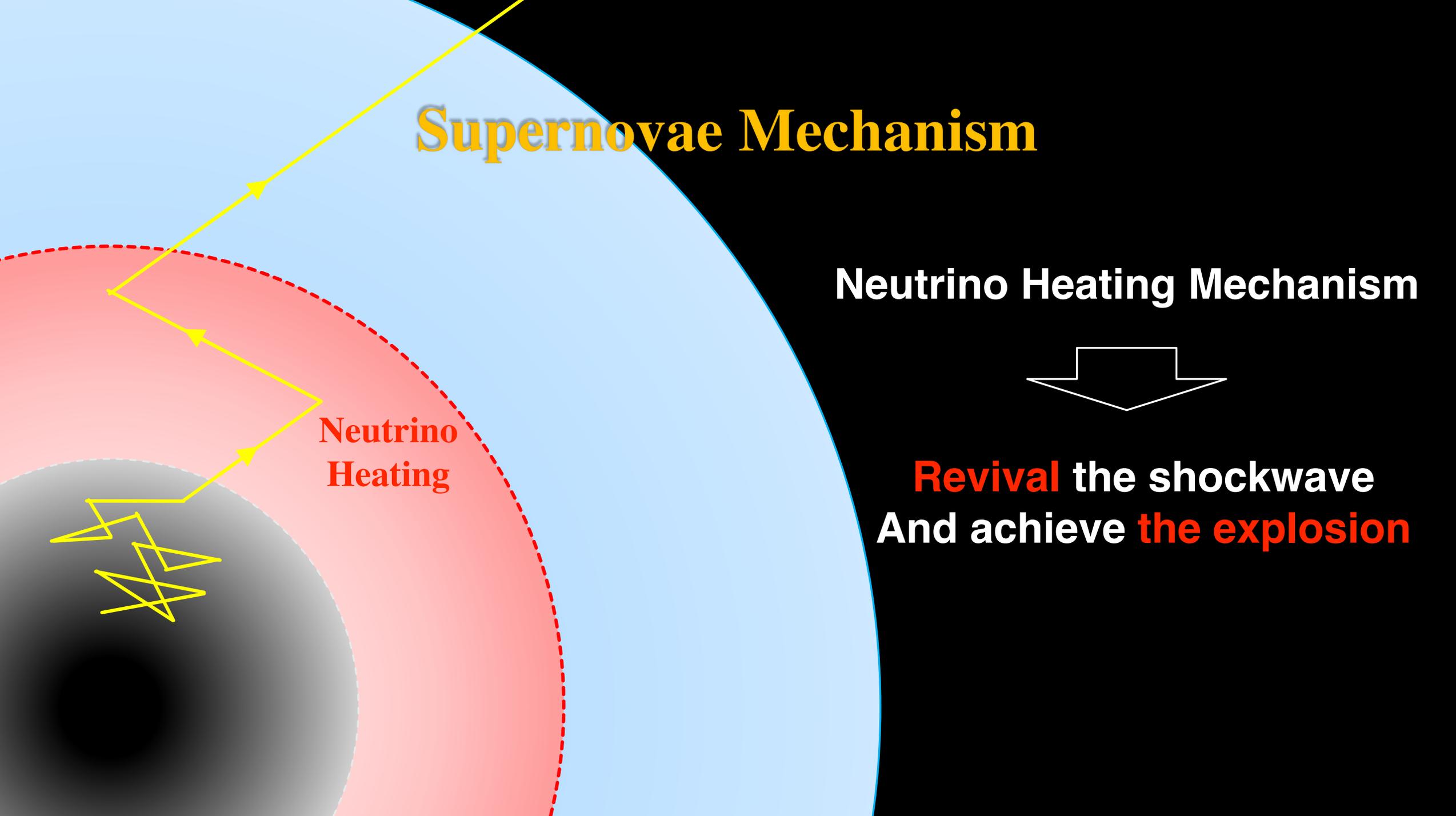
# 高解像度ボルツマン 輻射輸送計算に向けて

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# Supernovae Mechanism



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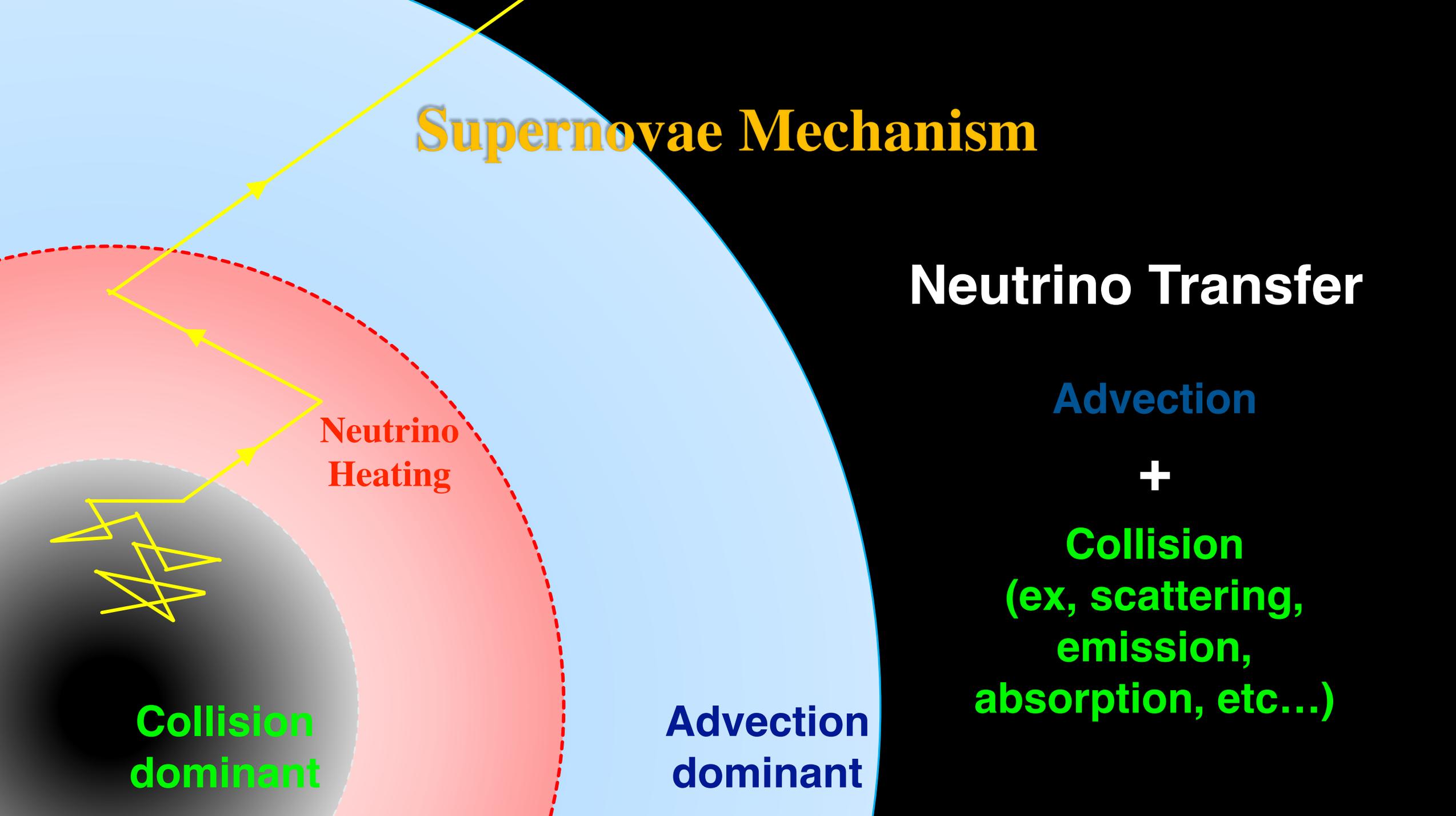


Neutrino Heating Mechanism



**Revival** the shockwave  
And achieve **the explosion**

# Supernovae Mechanism



## Neutrino Transfer

Advection

+

Collision  
(ex, scattering,  
emission,  
absorption, etc...)

Neutrino  
Heating

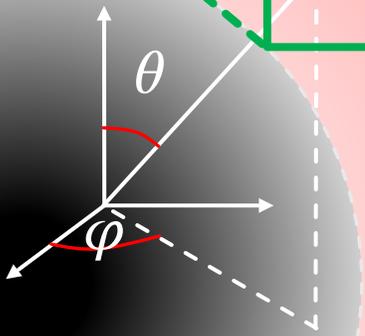
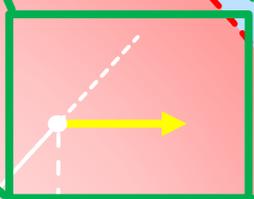
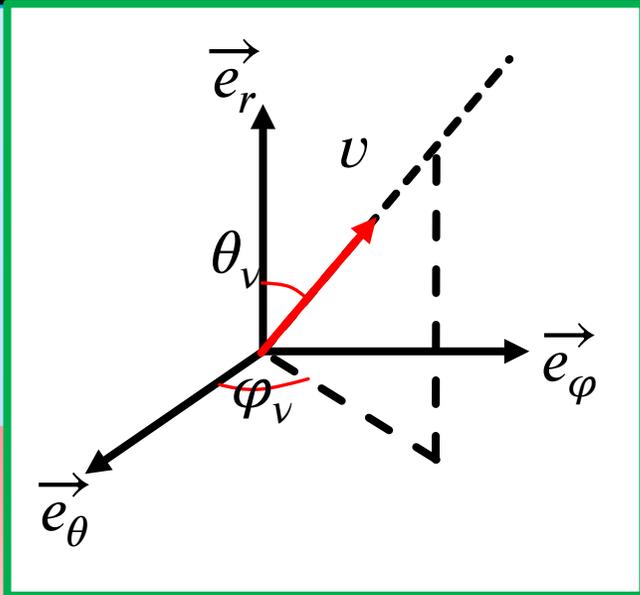
Collision  
dominant

Advection  
dominant

# Boltzmann Equation

$$\Omega = (\theta, \phi)$$

$$\Omega_v = (\theta_v, \phi_v)$$



## Boltzmann Eq.

$$\frac{1}{c} \frac{\partial f(r, \Omega, \epsilon_v, \Omega_v)}{\partial t} + \frac{\partial f(r, \Omega, \epsilon_v, \Omega_v)}{\partial s} = \frac{1}{c} \left[ \frac{\partial f(r, \Omega, \epsilon_v, \Omega_v)}{\partial t} \right]_{coll}$$

Boltzmann Equation :

**6D(space + momentum space) time evolution equation**

It costs many computational resources to solve Boltzmann equation

# Boltzmann Equation

**Boltzmann Eq.**  $\frac{1}{c} \frac{\partial f(r, \Omega, \epsilon_\nu, \Omega_\nu)}{\partial t} + \frac{\partial f(r, \Omega, \epsilon_\nu, \Omega_\nu)}{\partial s} = \frac{1}{c} \left[ \frac{\partial f(r, \Omega, \epsilon_\nu, \Omega_\nu)}{\partial t} \right]_{coll}$

**Advection term**                      **Collision term**

**Advection term**

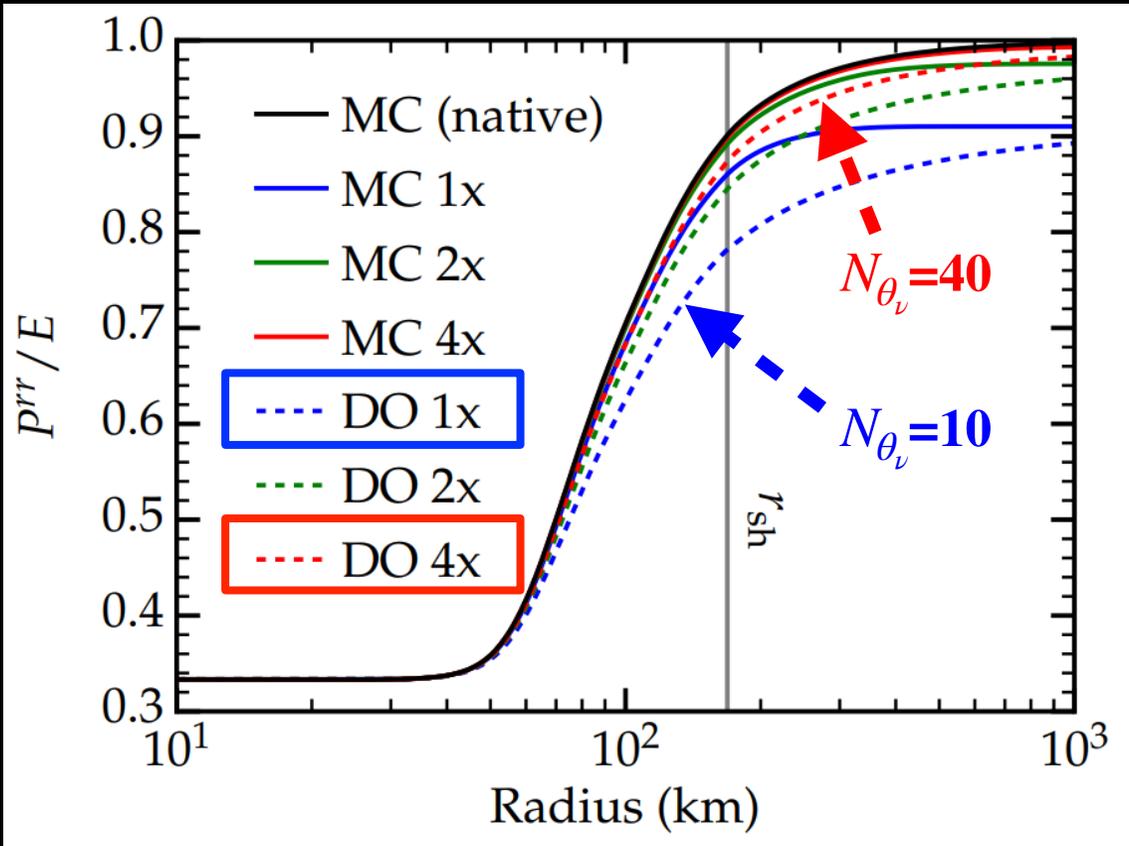
**Neutrino propagation**

**Collision term**

**Neutrino-matter interactions**

# Resolution Dependence

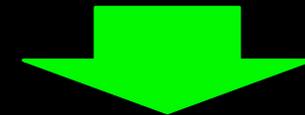
The  $rr$  component of the lab-frame energy-integrated Eddington tensor



S. Richers et. al. 2017

If the  $\theta_\nu$  resolution is 4 times higher than current, Eddington tensor  $P^{rr}$  differs by 10%.

As the radius is bigger, the error due to resolution is bigger



The resolution of the advection is not sufficient

# Solving the Boltzmann Equation

$$\text{Boltzmann Eq. } \frac{1}{c} \frac{\partial f(r, \Omega, \epsilon_v, \Omega_v)}{\partial t} + \frac{\partial f(r, \Omega, \epsilon_v, \Omega_v)}{\partial s} = \frac{1}{c} \left[ \frac{\partial f(r, \Omega, \epsilon_v, \Omega_v)}{\partial t} \right]_{coll}$$

**Advection term**                      **Collision term**

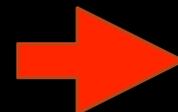
**Explicit**                                      **Implicit**

## Problem

- **Advection term** should be solved with **high angle resolution**
- If **Collision term** is solved with high resolution, a **significant increase** in costs is necessary.

## Idea

**Only the Advection term** is solved with **high angular resolution**  
**The resolution of the collision term is kept.**



**Dual Resolution Prescription**

# Method : Dual Resolution Prescription

$$\text{Advection : } f^*(\epsilon_v, \Omega_v) = f^n(\epsilon_v, \Omega_v) - c\Delta t S(\epsilon_v, \Omega_v, t(r, \Omega))$$

High Resolution

$$\text{Collision : } f^{n+1}(\epsilon_v, \Omega_v) = f^*(\epsilon_v, \Omega_v) - \Delta t \left[ \frac{\partial f^{n+1}(\epsilon_v, \Omega_v)}{\partial t} \right]_{coll}$$

Low Resolution

# Method : Dual Resolution Prescription

$$\text{Advection} : f^*(\epsilon_v, \Omega_v) = f^n(\epsilon_v, \Omega_v) - c\Delta t \mathcal{S}(\epsilon_v, \Omega_v, t(r, \Omega))$$

High Resolution

$$\text{Resolution Conversion} : f_{N_{fine}} \mapsto f_{N_{rough}}$$

$$\text{Collision} : f^{n+1}(\epsilon_v, \Omega_v) = f^*(\epsilon_v, \Omega_v) - \Delta t \left[ \frac{\partial f^{n+1}(\epsilon_v, \Omega_v)}{\partial t} \right]_{coll}$$

Low Resolution

$$\text{Resolution Conversion} : f_{N_{rough}} \mapsto f_{N_{fine}}$$

# Method : Dual Resolution Prescription

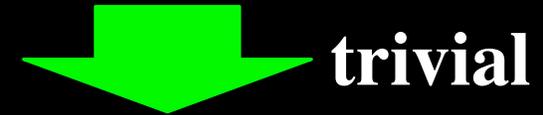
**Resolution Conversion** :  $f_{N_{fine}} \mapsto f_{N_{rough}}$

**Resolution Conversion** :  $f_{N_{rough}} \mapsto f_{N_{fine}}$

# Method : Dual Resolution Prescription

**Resolution Conversion** :  $f_{N_{fine}} \mapsto f_{N_{rough}}$

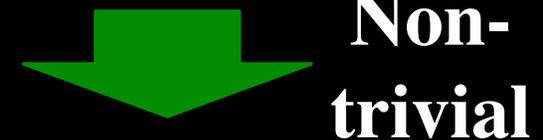
**High Resolution**



**Low Resolution**

**Resolution Conversion** :  $f_{N_{rough}} \mapsto f_{N_{fine}}$

**Low Resolution**



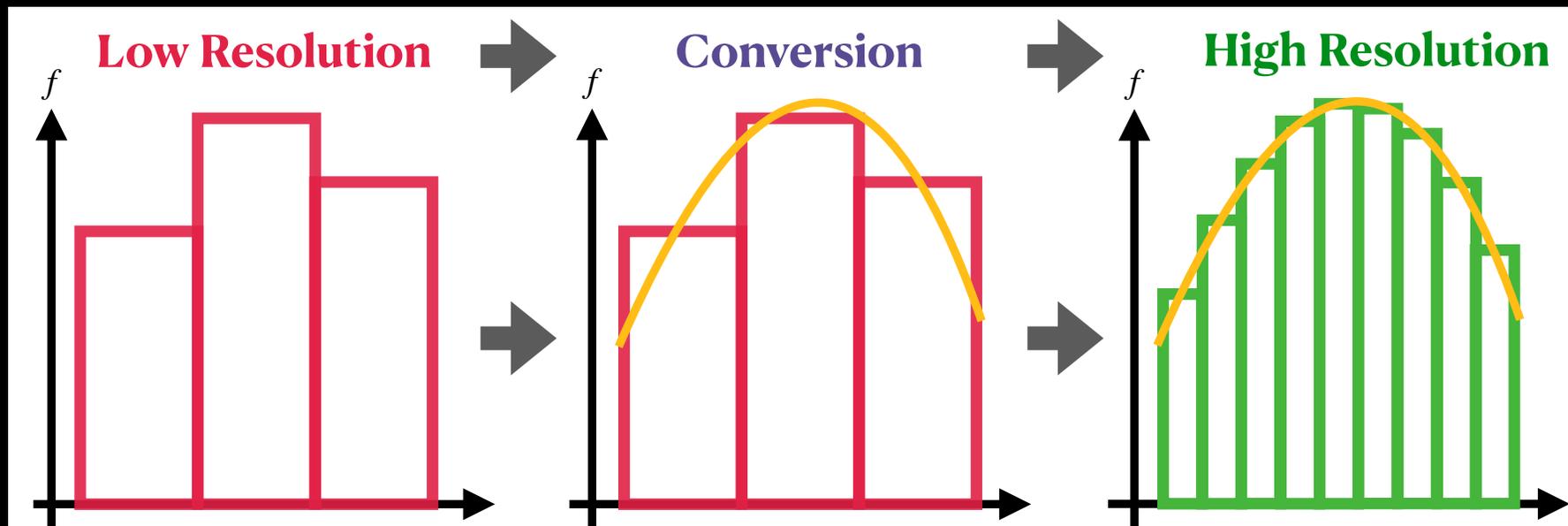
**High Resolution**

# Method : Dual Resolution Prescription

Resolution Conversion :  $f_{N_{rough}} \mapsto f_{N_{fine}}$

Non-trivial

The distribution function is assumed to be expressed by the polynomial



# Method : Dual Resolution Prescription

$$\text{Resolution Conversion} : f_{N_{rough}} \mapsto f_{N_{fine}}$$

The distribution function is assumed to be expressed by the polynomial

$$f_{\text{appr}}(\mu = \cos \theta_\nu) = \sum_{i=0}^{N_{\theta_\nu}^{\text{poly}}} a_n \mu^n \quad N_{\theta_\nu}^{\text{poly}} : \text{The order of polynomial}$$

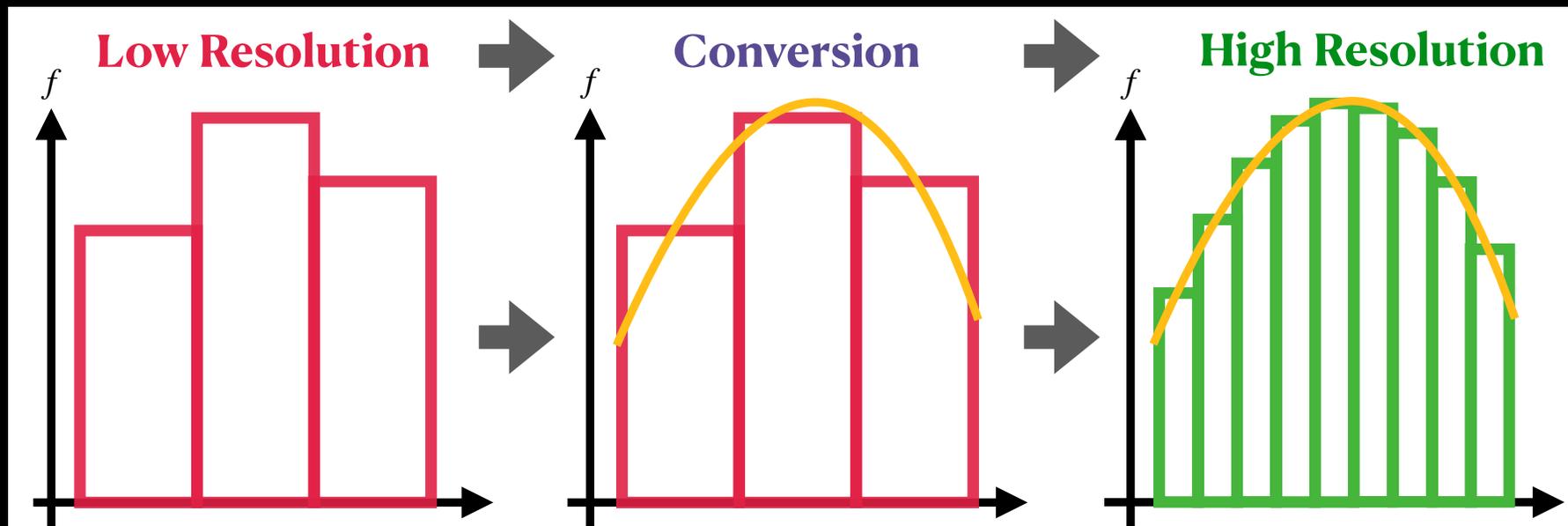
The distribution function in low mesh is

$$\begin{aligned} f_{\text{high}}(\mu_i = \cos \theta_{\nu,i})(\mu_I - \mu_{I-1}) &= \int_{\mu_{I-1}}^{\mu_I} d\mu f_{\text{appr}}(\mu) \\ &= \sum_{n=0}^{N_{\theta_\nu}^{\text{poly}}} \frac{a_n}{n+1} (\mu_I^{n+1} - \mu_{I-1}^{n+1}) \end{aligned}$$

# Method : Dual Resolution Prescription

Resolution Conversion :  $f_{N_{rough}} \mapsto f_{N_{fine}}$

The distribution function is assumed to be expressed by the polynomial



How accurate is this conversion?

# Methods

In the study, we perform **the pilot study**  
in order to **test the effectiveness** of this dual resolution prescription

In dual resolution prescription

the advection term  rigorous expression  
in the equations

we should study

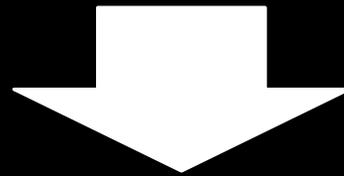
**the resolution dependence of the collision term!**

# Method : Advection term

**Boltzmann Eq.**

**Advection term**

$$\frac{1}{c} \frac{\partial f(r, \Omega, \epsilon_\nu, \Omega_\nu)}{\partial t} + \frac{\partial f(r, \Omega, \epsilon_\nu, \Omega_\nu)}{\partial s} = \frac{1}{c} \left[ \frac{\partial f(r, \Omega, \epsilon_\nu, \Omega_\nu)}{\partial t} \right]_{coll}$$



**the spatial dependence is eliminated**

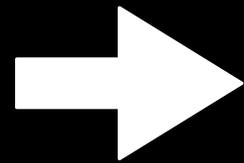
$$\frac{1}{c} \frac{\partial f(\epsilon_\nu, \Omega_\nu)}{\partial t} = \frac{1}{c} \left[ \frac{\partial f(\epsilon_\nu, \Omega_\nu)}{\partial t} \right]_{coll} + S(\epsilon_\nu, \Omega_\nu, t(r, \Omega))$$

**Source term**

# Methods

we should study **the resolution dependence of the collision term!**

**Advection term** is replaced by the artificial source term



Only 3D(momentum) remains

As the artificial source term,

**2 types of the source terms** were carried out:

Steady state Test & Time evolution Test

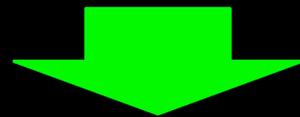
# Method : Advection term

**Steady state Test : resolution dependence**

$$S_{\text{steady}} = -\frac{1}{c} \left[ \frac{\partial f_{\text{ref}}(\epsilon_\nu, \Omega_\nu)}{\partial t} \right]_{\text{coll}}$$

$$\frac{1}{c} \frac{\partial f(\epsilon_\nu, \Omega_\nu)}{\partial t} = \frac{1}{c} \left[ \frac{\partial f(\epsilon_\nu, \Omega_\nu)}{\partial t} \right]_{\text{coll}} + S_{\text{steady}}(\epsilon_\nu, \Omega_\nu, t(r, \Omega))$$

**In this case, the source term and collision term are canceled**



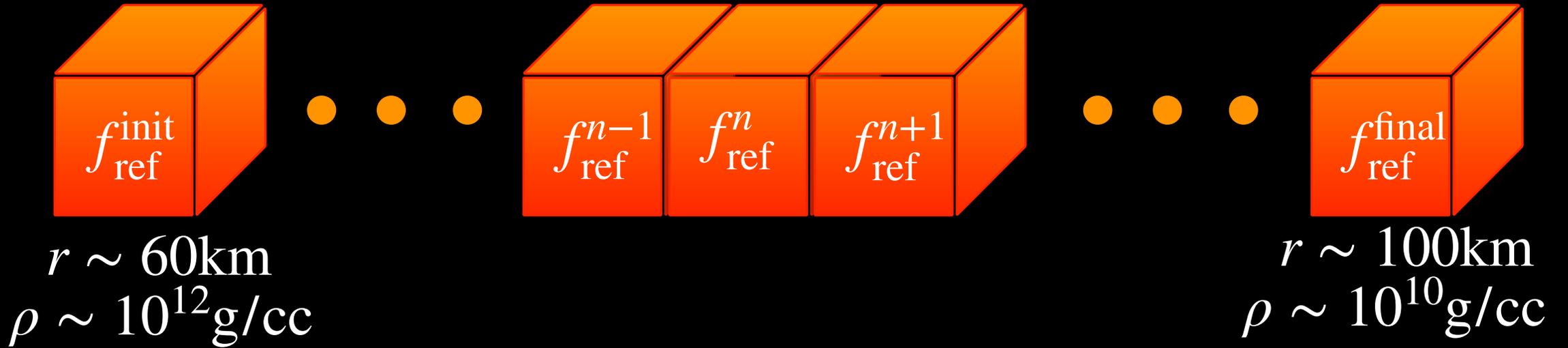
**The source term reproduce the reference distribution function  
if the resolutions are reference ones**

# Method : Advection term

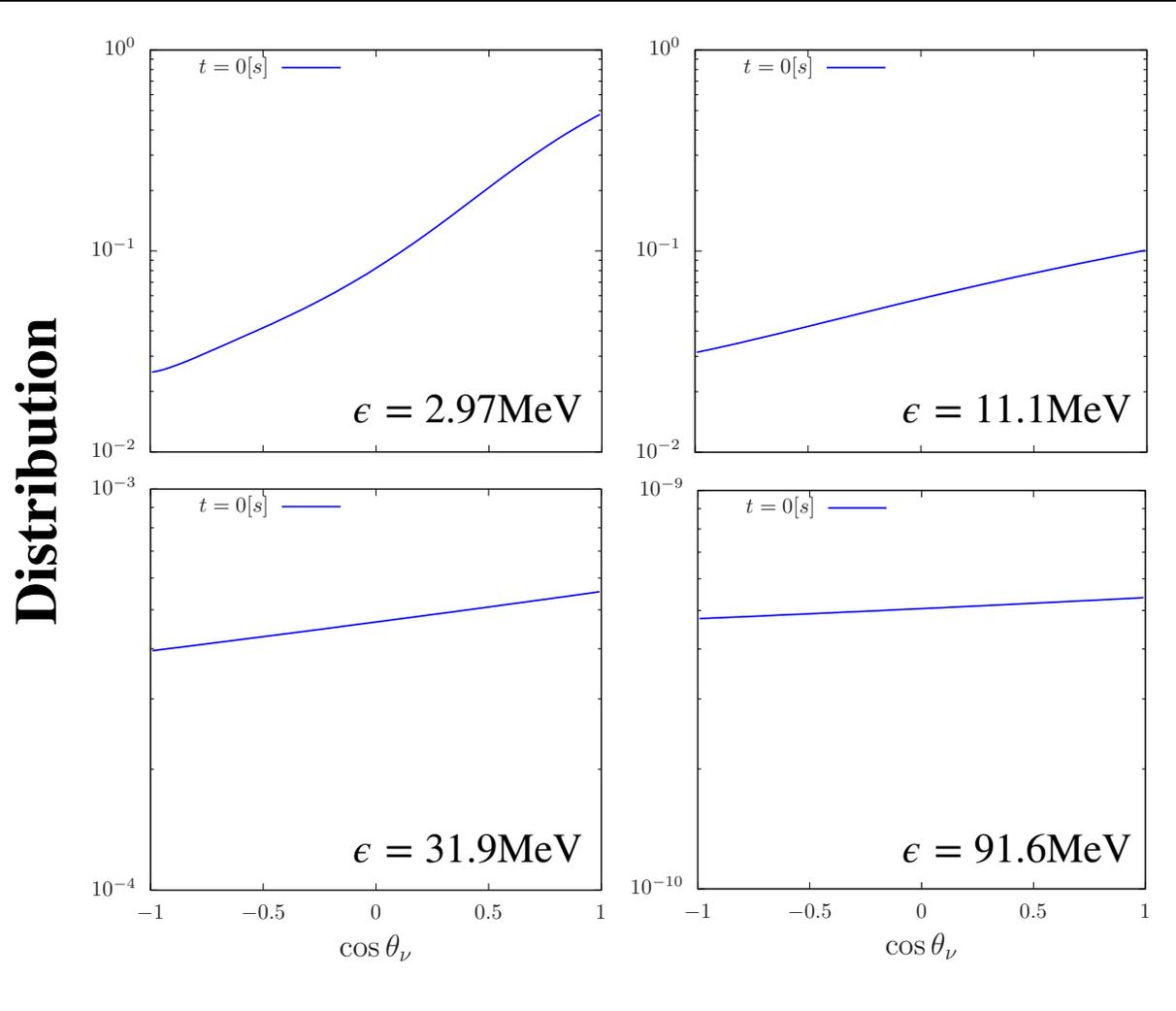
**Time evolution Test : resolution and interpolation dependence**

$$S_{\text{time}}^n = \frac{f_{\text{ref}}^{n+1} - f_{\text{ref}}^n}{\delta t} - \frac{1}{c} \left[ \frac{\partial f_{\text{ref}}^{n+1}(\epsilon_\nu, \Omega_\nu)}{\partial t} \right]_{\text{coll}}$$

**The source term reproduce the reference distribution function  
if the resolutions are reference ones**



# Result : Steady state Test in $\theta_\nu$



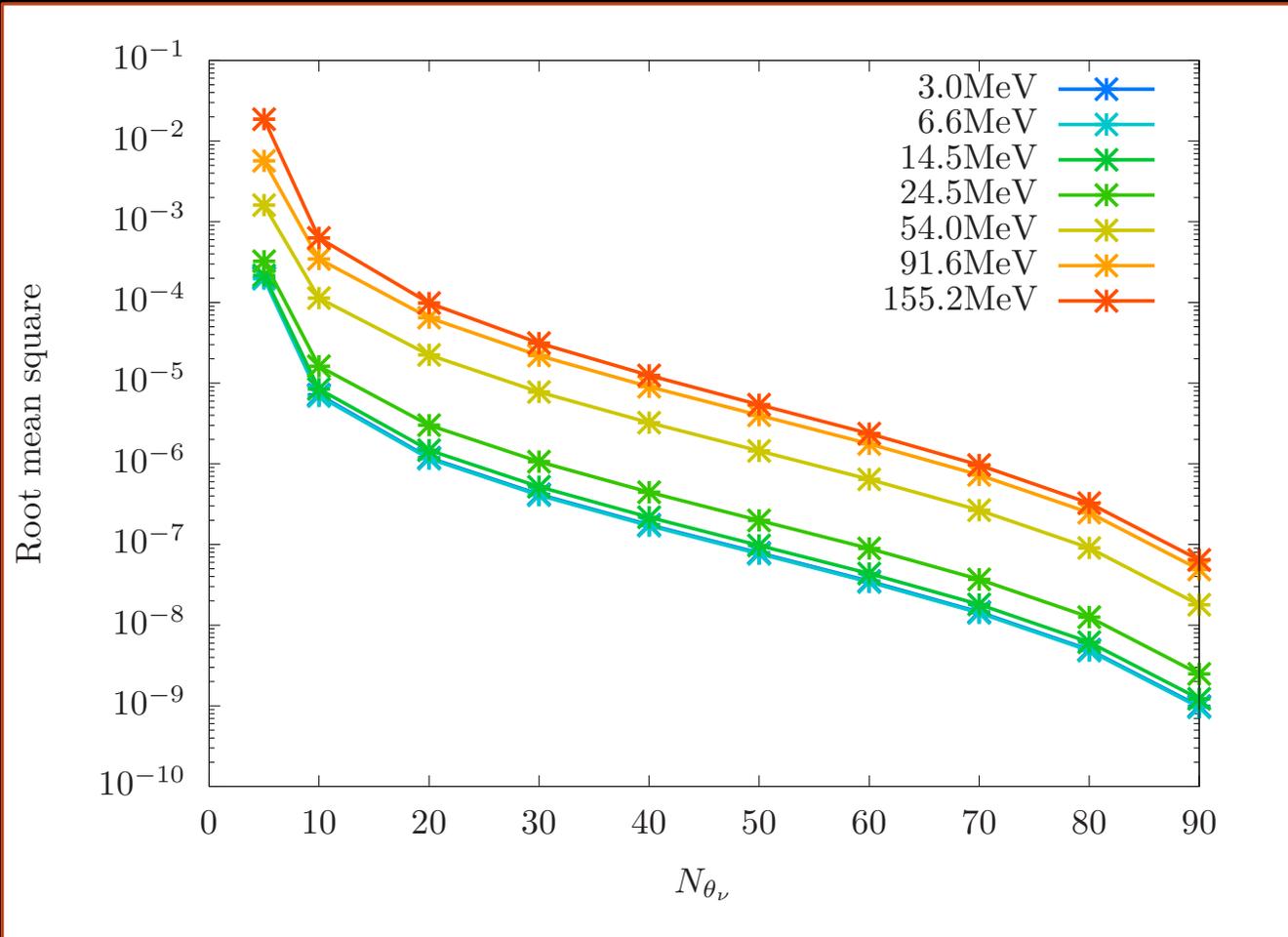
**Reference Resolution :**

$$N_\epsilon = 20$$

$$N_{\theta_\nu} = 100$$

$$N_{\phi_\nu} = 6$$

# Result : Steady state Test in $\theta_\nu$



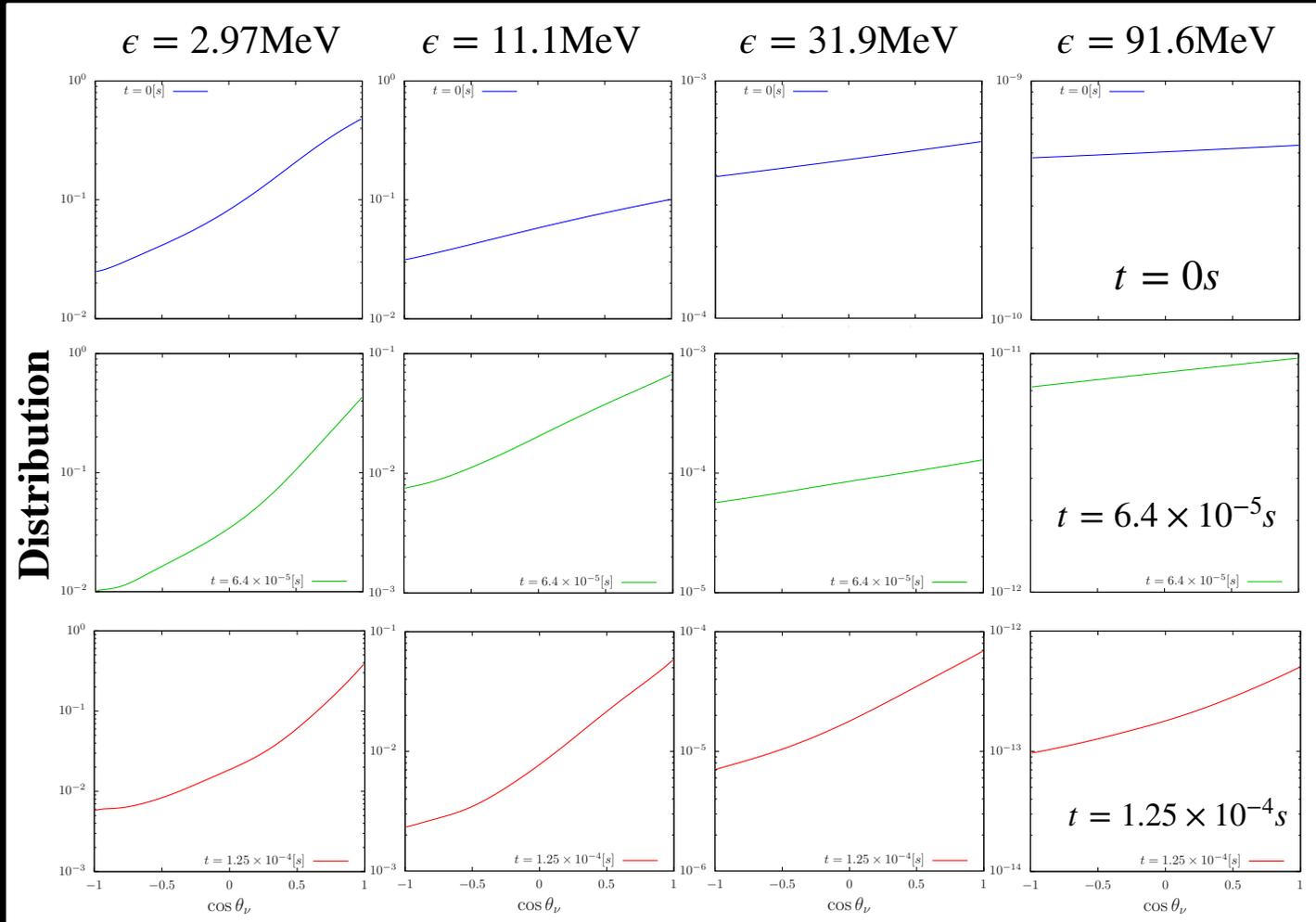
Reference Resolution :  $N_{\theta_\nu} = 100$

$$\text{Error} : \left| \frac{f_{N_{\theta_\nu}} - f_{N_{\theta_\nu}=100}}{f_{N_{\theta_\nu}=100}} \right|$$

If the resolution is current, ( $N_{\theta_\nu}=10$ ),  
the accuracy (RMS  $\sim 10^{-4}$ ) is acceptable

# Result : Time Evolution in $\theta_\nu$

## Reference distribution



## Reference Resolution :

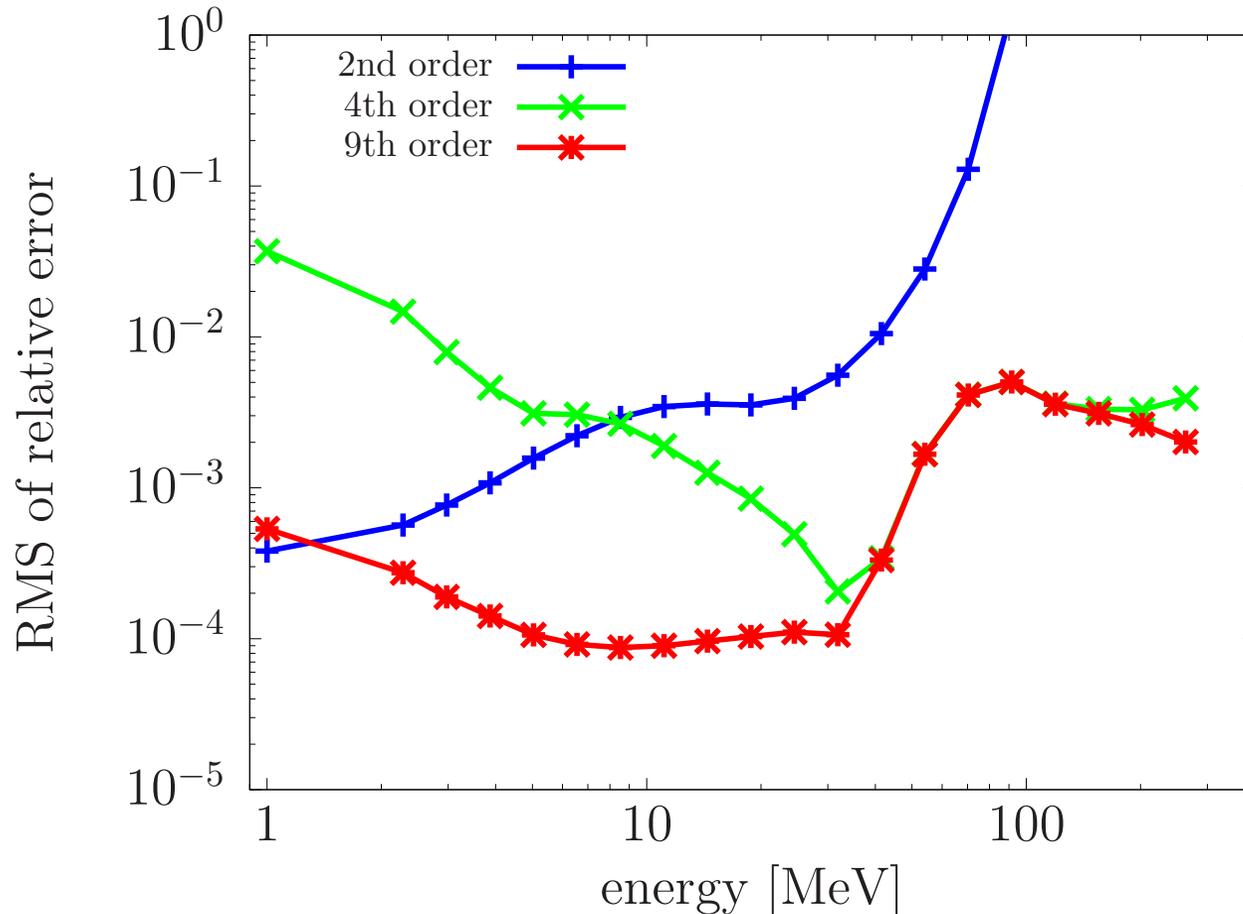
$$N_\epsilon = 20$$

$$N_{\theta_\nu} = 40$$

$$N_{\phi_\nu} = 6$$

# Result : Time Evolution in $\theta_\nu$

## The interpolation dependence

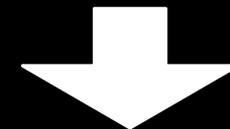


Reference Resolution :  $N_{\theta_\nu} = 40$

Interpolation : Polynomial

$$\text{Error} : \left| \frac{f_{N_{\theta_\nu}} - f_{N_{\theta_\nu}=40}}{f_{N_{\theta_\nu}=40}} \right|$$

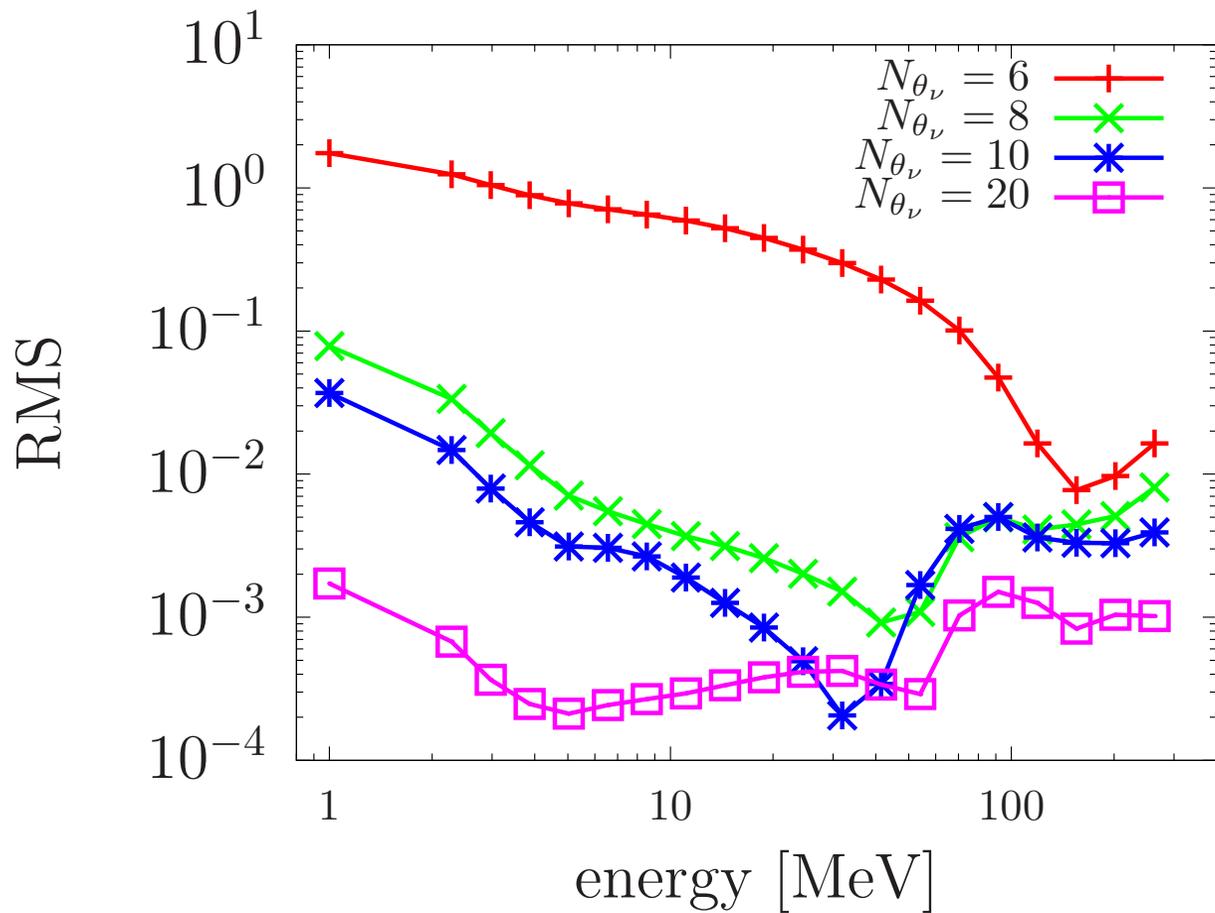
Higher order of polynomial



Smaller relative error

# Result : Time Evolution in $\theta_\nu$

The resolution dependence



Reference Resolution :  $N_{\theta_\nu} = 40$

Interpolation : 4th-order Polynomial

$$\text{Error} : \left| \frac{f_{N_{\theta_\nu}} - f_{N_{\theta_\nu}=40}}{f_{N_{\theta_\nu}=40}} \right|$$

Higher resolution



Smaller relative error

# Result : Steady state Test in $\phi_\mu$

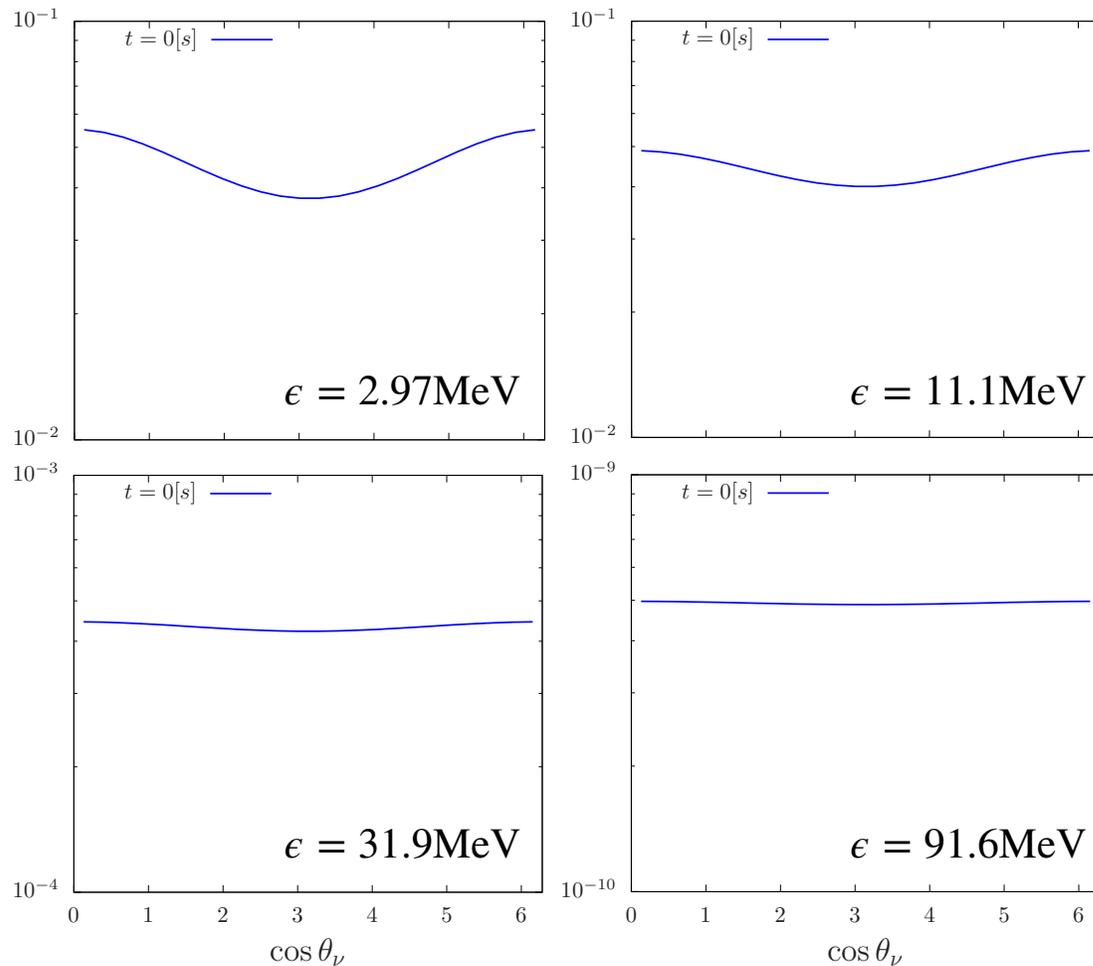
Reference Resolution :

$$N_\epsilon = 20$$

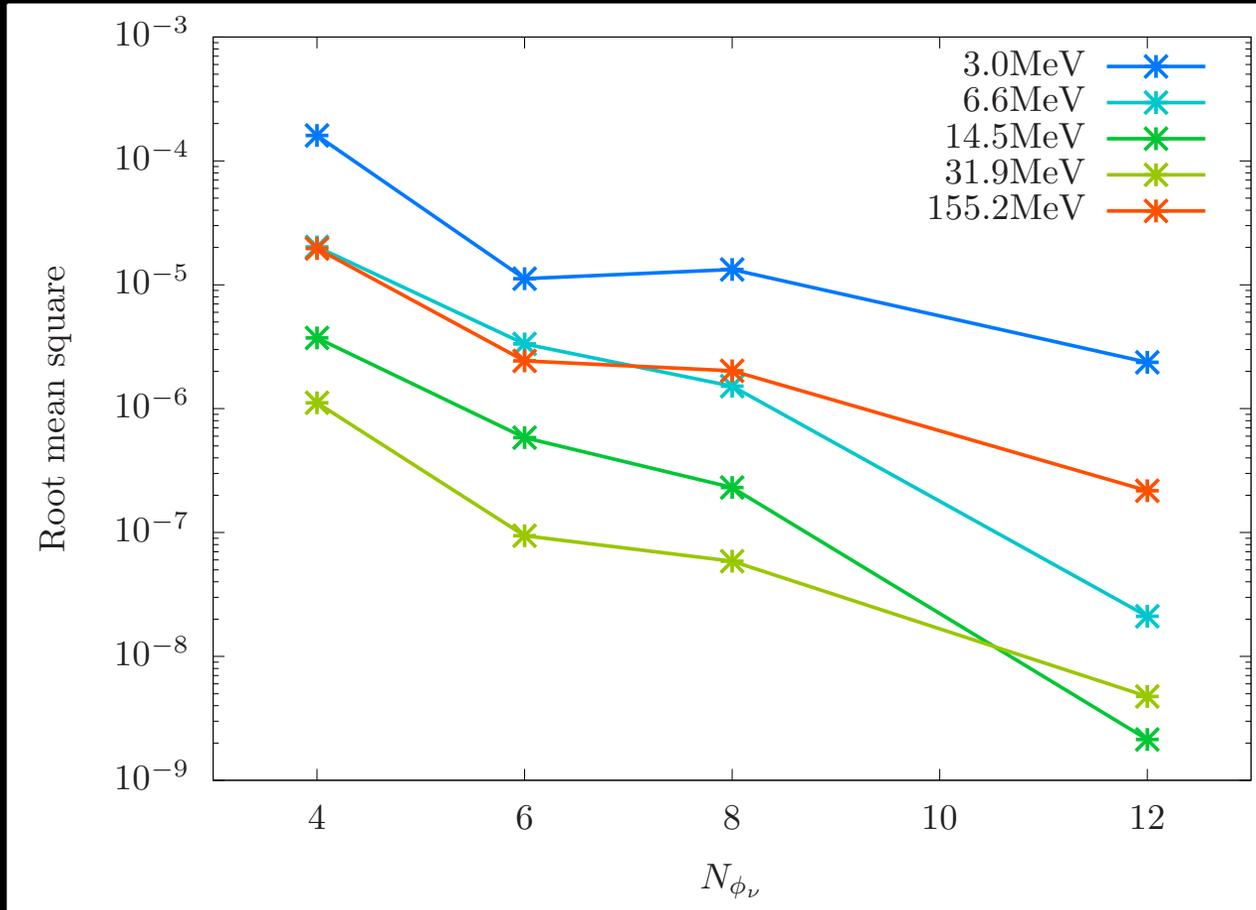
$$N_{\theta_\nu} = 10$$

$$N_{\phi_\nu} = 24$$

Distribution



# Result : Steady state Test in $\phi_\mu$



**Reference Resolution :**

$$N_e = 20$$

$$N_{\theta_\nu} = 10$$

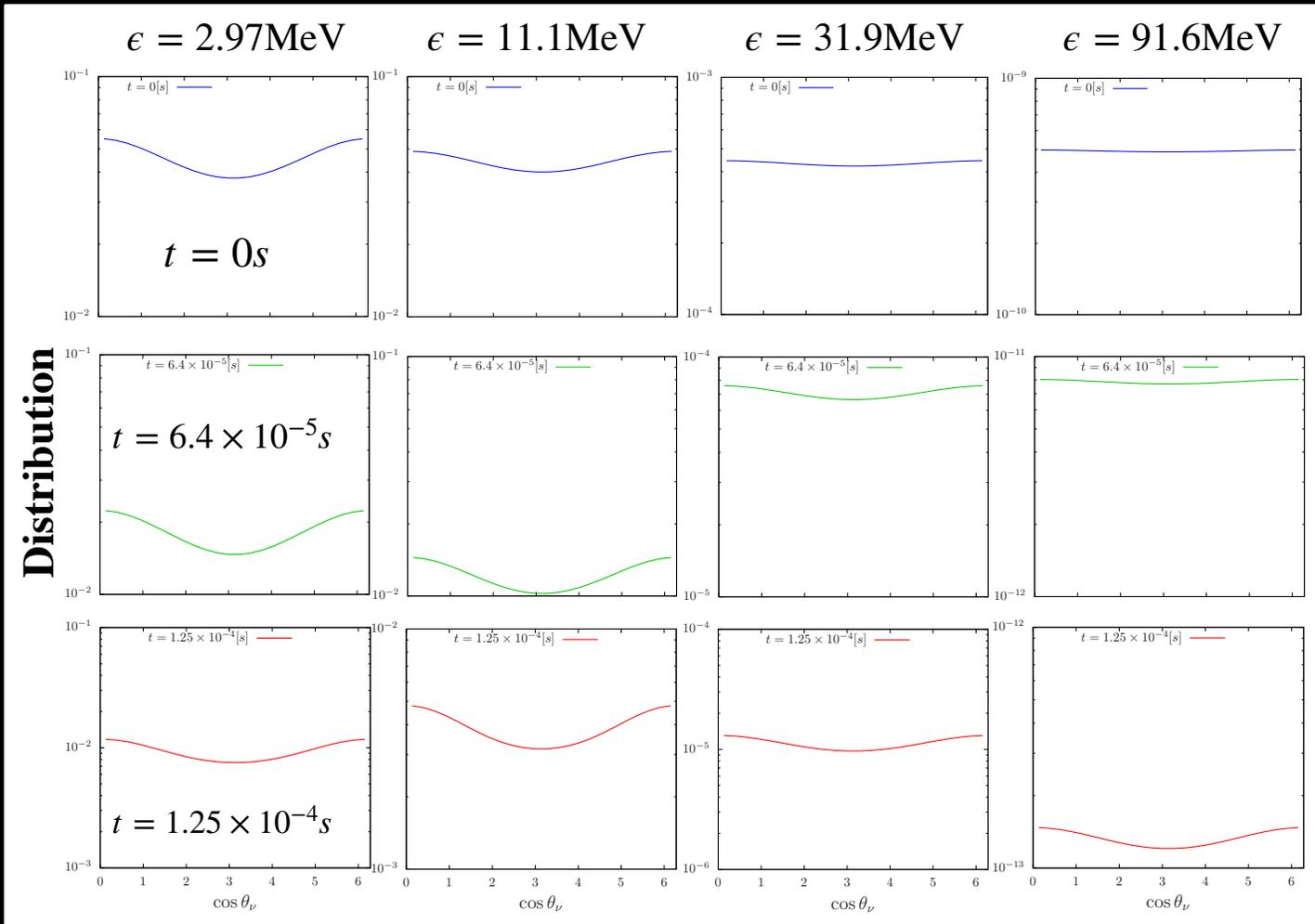
$$N_{\phi_\nu} = 24$$

$$\text{Error : } \left| \frac{f_{N_{\phi_\nu}} - f_{N_{\phi_\nu}=24}}{f_{N_{\phi_\nu}=24}} \right|$$

**If the resolution is current, ( $N_{\phi_\nu}=6$ ),  
the accuracy ( $\text{RMS} \lesssim 10^{-4}$ ) is acceptable**

# Result : Time Evolution in $\phi_\mu$

## Reference Distribution



Reference Resolution :

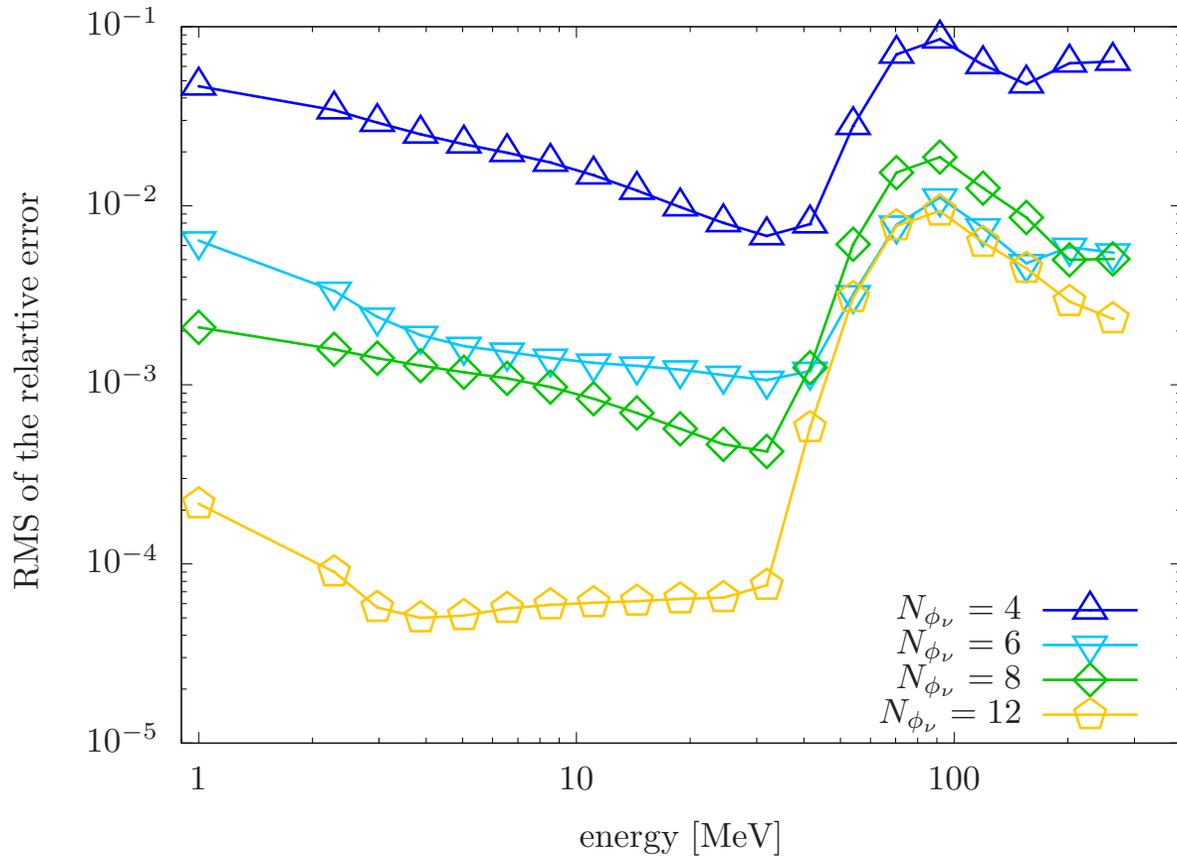
$$N_\epsilon = 20$$

$$N_{\theta_\nu} = 10$$

$$N_{\phi_\nu} = 24$$

# Result : Time Evolution in $\phi_\mu$

## The resolution dependence



Reference Resolution :

$$N_\epsilon = 20$$

$$N_{\theta_\nu} = 10$$

$$N_{\phi_\nu} = 24$$

Interpolation : 4th-order Polynomial

$$\text{Error} : \left| \frac{f_{N_{\phi_\nu}} - f_{N_{\phi_\nu}=24}}{f_{N_{\phi_\nu}=24}} \right|$$

Higher resolution



Smaller relative error

## Summary

- ◆ We perform the pilot study with dual resolution prescription and study the momentum angle resolution dependence.
- ◆ When  $N_{\theta_\nu} = 10$  ( $N_{\phi_\nu} = 6$ ) with 9th (4th) order polynomial, the error ( $\lesssim 10^{-2}$ ) is acceptable and the dual resolution prescription is valid.

## Future Work

- ◆ We will implement the Dual Resolution Prescription into the Boltzmann code and reveal the neutrino behavior with the high resolution calculation.